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A FINITE ELEMENT FORMULATION
FOR NONLINEAR VISCOPLASTIC
MATERIAL USING A Q-MODEL

by

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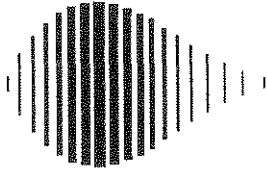
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ABSTRACT

A direct comparison of the rheological behavior of a mass supported by a complex spring and a mass supported by a spring-dashpot arrangement provides a means of establishing a relationship between material damping in the frequency domain and the frequency dependent $Q(\omega)$. For the special case where $Q(\omega)$ is constant over a given frequency range, the expansion of $Q^{-1}(\omega)$ into a Laurent Series yields a set of damping coefficients whose values are determined by minimizing the mean square error of the series over the prescribed frequency range. The resulting damping expression is used in conjunction with an elastoplastic constitutive matrix in finite element discretization to produce a viscoplastic model suitable for a direct step by step time integration. The proposed model is very convenient for use in finite element discretization for the analyses of earthquake, blast, shock, and other soil-structure interaction problems involving cyclic loading.

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SECTION 1

INTRODUCTION

The complex behavior of geologic materials has in recent years generated extensive investigation by researchers to develop material models to predict the path-dependent behavior of such materials within the framework of the theory of classical plasticity. To better predict the hysteresis displayed by soils when loaded and unloaded hydrostatically, the cap model (27) was developed. This model, which is based on a modified classical plasticity theory, was further modified to include the effect of kinematic hardening (11). The bounding surface plasticity theory which utilizes two yield surfaces was proposed to simulate the behavior of sand under cyclic loading (19). While the above mentioned models predict the path-dependent behavior of material fairly well, they fail to address the dependence of material behavior on time. In many geomechanical problems, however, material behavior is also governed by rheological behavior. To account for energy dissipation which is associated with foundation media in soil-structure interaction problems, the material has been idealized to be viscoelastic solid (23,24). The effect of Q on wave attenuation and velocity dispersion in geologic materials has been investigated by researchers in recent years within the framework of viscoelasticity (2,3,4). In Ref. 28, the convolution integral relating stress to strain history in geologic media (assumed viscoelastic) has been transformed into a convergent sequence of constant coefficient differential operators of increasing order, through the use of Pade approximants, with the value of Q assumed over a prescribed frequency range. In this format, the convolution integral is reduced into a form suitable for time stepping in finite difference schemes for wave propagation problems. The above mentioned viscoelastic models, however, fail to account for the path-dependent behavior of materials. A more realistic model would be a viscoplastic model which accounts for the path dependence as well as the energy dissipation characteristics of the material. Several attempts have been made by researchers to predict the time as well as path-dependence of material in recent years (5,6,18,20).

The contribution of this research is to provide explicit representation of the viscous and energy dissipation characteristics of material suitable for time stepping in finite element discretization through the Q -model, $Q^{-1}(\omega)$ being a measure of energy dissipation per cycle in a medium during cyclic loading. A relation between damping in the frequency domain and the frequency dependent $Q^{-1}(\omega)$ is established through a direct comparison of rheological representation of a mass

supported by a complex spring and a mass supported by a spring-dashpot arrangement, in the frequency domain. For the special case where $Q(\omega)$ is constant over a prescribed frequency range, the expansion of $Q^{-1}(\omega)$ into a Laurent Series yields a set of damping coefficients whose values are determined by minimizing the mean square error of the series over the prescribed frequency range. The resulting damping expression is used in conjunction with an elastoplastic constitutive matrix in finite element discretization to produce an elasto-viscoplastic model suitable for step by step time integration. It should be mentioned that the proposed model follows quite a different approach from those proposed in Refs. 5 and 6. Also, earlier finite element models for elastoviscoplastic materials have focused on the quasi-static behavior of the materials (25,26). The merits of the proposed model lie not only in the ease of its application in dynamic problems but also in the readiness in which the parameters involved could be determined from dynamic tests in the laboratory. The proposed model will find very useful application in seismic problems such as design of dams, bridges, buildings, subjected to earthquake excitation. The model could be adapted to suit viscoelastic problems simply by replacing the elastoplastic constitutive matrix with an elastic constitutive matrix. This makes the model suitable for a much wider variety of problems.

SECTION 2 SURVEY OF SOIL MODELS

2.1 Simple Plasticity Models

The elastic ideally-plastic soil model with a fixed yield surface defined by

$$F(\sigma_{ij}) = 0 \quad (2-1)$$

characterizes the earlier plasticity soil models. Here σ_{ij} defines a stress point in the stress space. In terms of stress invariants, the yield surface is given by

$$F(J_1, J_2, J_3) = 0 \quad (2-2)$$

where J_1, J_2, J_3 are the first, second, and third stress invariants, respectively.

The material is elastic when the stress point lies inside the yield surface, in which case changes in stresses result in recoverable deformations. The increment in elastic strain is given by

$$d\epsilon_{ij}^e = \frac{1}{2G} ds_{ij} + \frac{1}{9K} \delta_{ij} d\sigma_{kk} \quad (2-3)$$

where:

ds_{ij} = Increment in deviatoric stress

$d\sigma_{kk}$ = Increment in volumetric stress

δ_{ij} = Kronecker delta

K = Bulk modulus

G = Shear modulus

On the yield surface, total strain increment is given by

$$d\epsilon_{ij} = d\epsilon_{ij}^e + d\epsilon_{ij}^p \quad (2-4)$$

where $d\epsilon_{ij}^e$ is the elastic strain increment defined by equation (2-3) and

$$d\epsilon_{ij}^p = \lambda \frac{\partial q}{\partial \sigma_{ij}} \quad (2-5)$$

where:

q = plastic potential

λ = nonnegative scalar function

$d\epsilon_{ij}^p$ = plastic strain rate

If $F = q$ then equation (2-5) becomes

$$d\epsilon_{ij}^p = \begin{cases} \lambda \frac{\partial F}{\partial \sigma_{ij}} & F=0 \\ 0 & F<0 \end{cases} \quad (2-6)$$

and the flow rule is said to be associated, and hence the plastic strain rate is normal to the yield surface at the current stress point. Furthermore, for a convex yield surface, uniqueness is assured. In general, however, when the yield surface does not coincide with the plastic potential, i.e., $F \neq q$, the corresponding flow rule is said to be non-associated. In this case, uniqueness cannot, in general, be proved.

Stress outside the yield surface, i.e., $F > 0$ is not permitted.

Uniqueness and Stability

Uniqueness, stability, and continuity are basic requirements for continuum models. From Drucker's stability postulate (Ref. 7) nonnegative work must be done by an external agent in any excursion from equilibrium. In particular for any stress cycle, where σ_{ij}^* is the stress at equilibrium state

$$\int (\sigma_{ij} - \sigma_{ij}^*) d\epsilon_{ij} \geq 0 \quad (2-7)$$

The equal sign applies only for elastic or reversible paths. Satisfying Drucker's postulate is sufficient (but not necessary) to insure uniqueness and continuity.

By eliminating the elastic or reversible strains and by choosing σ_{ij}^* (equation (2-7) on the yield surface, the condition for stability in the "small" for elastic-plastic models is obtained, i.e.,

$$d\sigma_{ij} d\epsilon_{ij}^p \geq 0 \quad (2-8)$$

This means that the yield condition can only move outward (or not move) at a stress point, i.e., work softening or strain softening is not permitted.

The condition for stability in the "large" for elastoplastic materials is given by

$$(\sigma_{ij} - \sigma_{ij}^*) d\epsilon_{ij}^p \geq 0 \quad (2-9)$$

It can be noted that equation (2-9) requires the normality of the plastic strain rate vector, and the convexity of the yield surface. (See Figure 2-1.)

2.2 von Mises

The von Mises yield condition is given by

$$\sqrt{J_2'} = k \quad (2-10)$$

where J_2' is the second invariant of the deviatoric stress and is given by

$$J_2' = \frac{1}{2} s_{ij} s_{ij} \quad (2-11)$$

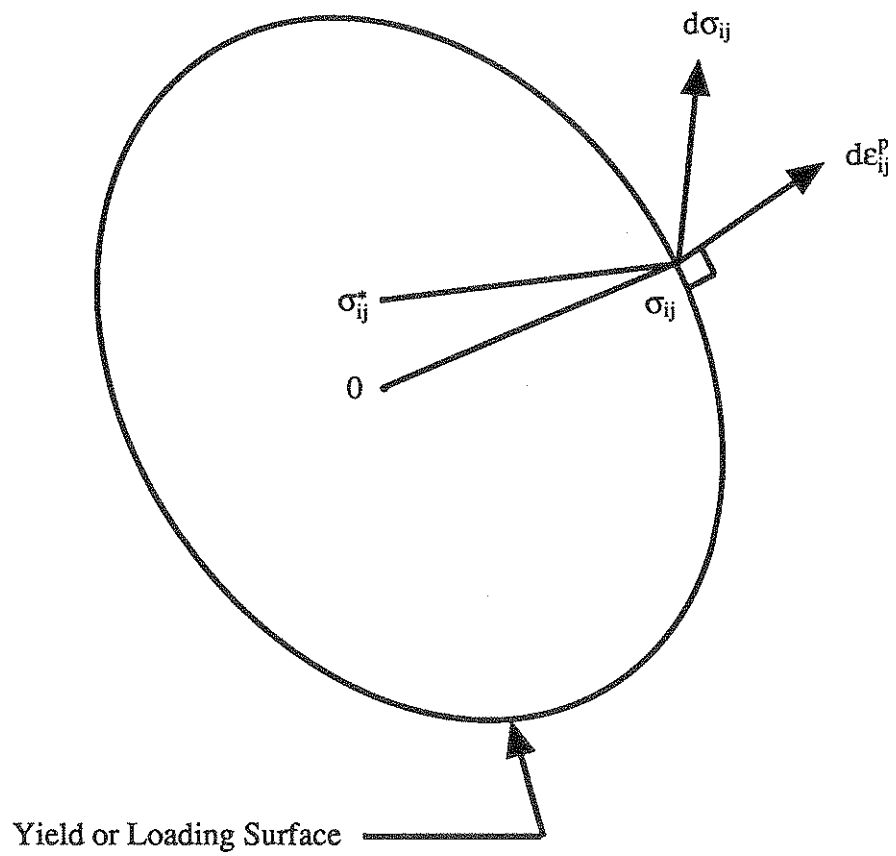


FIGURE 2-1 Geometric Interpretation of Drucker's Stability Postulate

and k is a constant.

The von Mises yield surface is a cylinder in the principal stress space (see Figure 2-2), and is a good representation of the failure surface of many saturated clays.

2.3 Drucker and Prager

Drucker and Prager, Ref. (10), proposed an ideally plastic yield condition for granular materials given by

$$F = \sqrt{J_2} + \alpha J_1 - k = 0 \quad (2-12)$$

This is a modified Mohr-Coulomb failure criterion which reduces to the Mohr-Coulomb equations in a triaxial test when

$$\alpha = \frac{2 \sin \phi}{\sqrt{3} (3 - \sin \phi)} \quad (2-13)$$

and

$$k = \frac{6c \cos \phi}{\sqrt{3} (3 - \sin \phi)} \quad (2-14)$$

In equation (2-12), J_1 and J_2 are, respectively, the first and second invariants of stress and stress deviator. α and k are defined by equations (2-13) and (2-14) respectively, ϕ is the friction angle and c is cohesion of soil. (See Figures 2-3 and 2-4.) This model satisfies Drucker's postulates for stability and uniqueness when used with the associated flow rule as defined earlier. However, the model has the following shortcomings:

1. Due to the normality principle and the concept of the associated flow rule, considerable dilatancy effects are predicted which are much greater than observed experimentally;

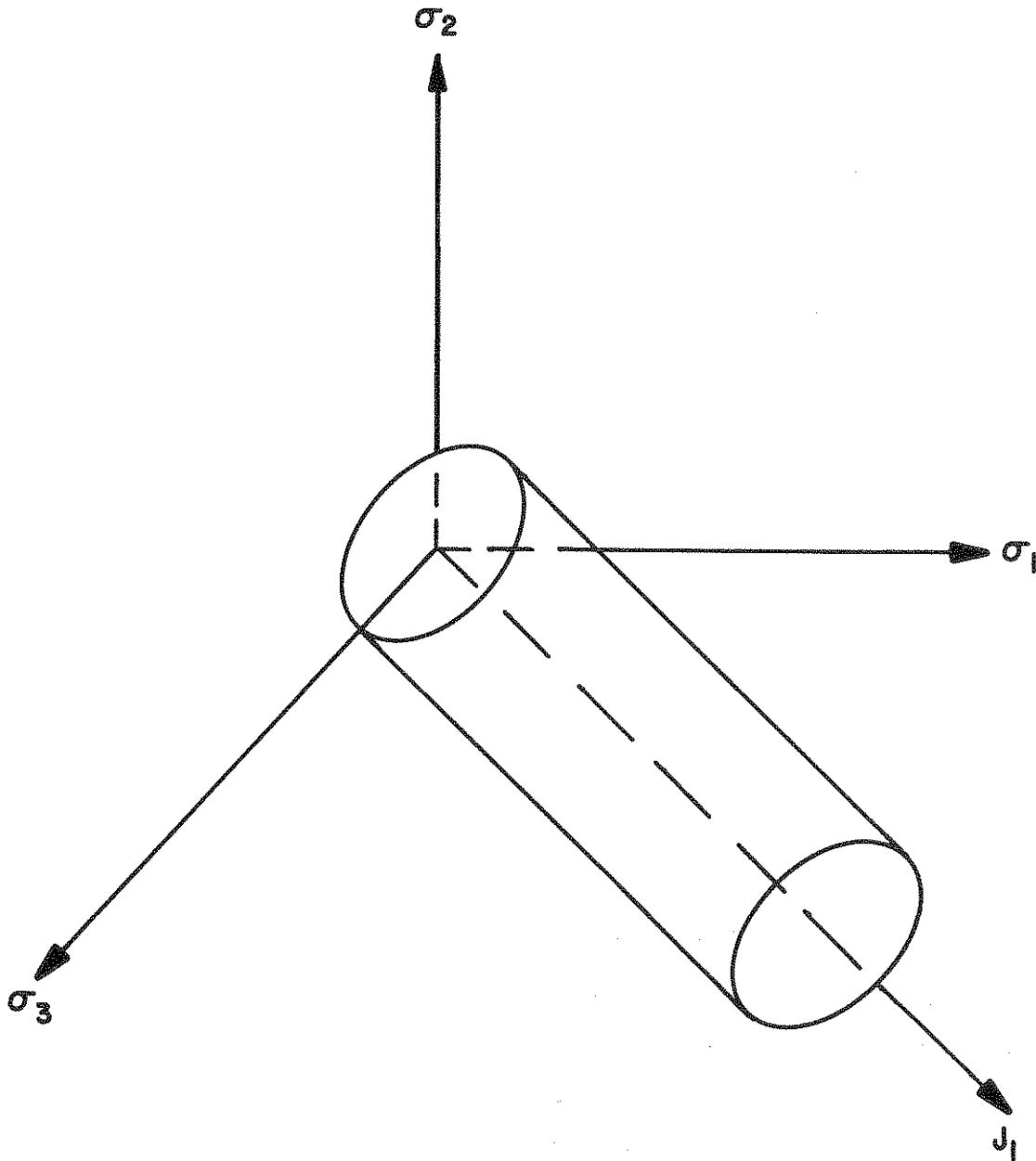


FIGURE 2-2 Von Mises Yield Surface Cylinder

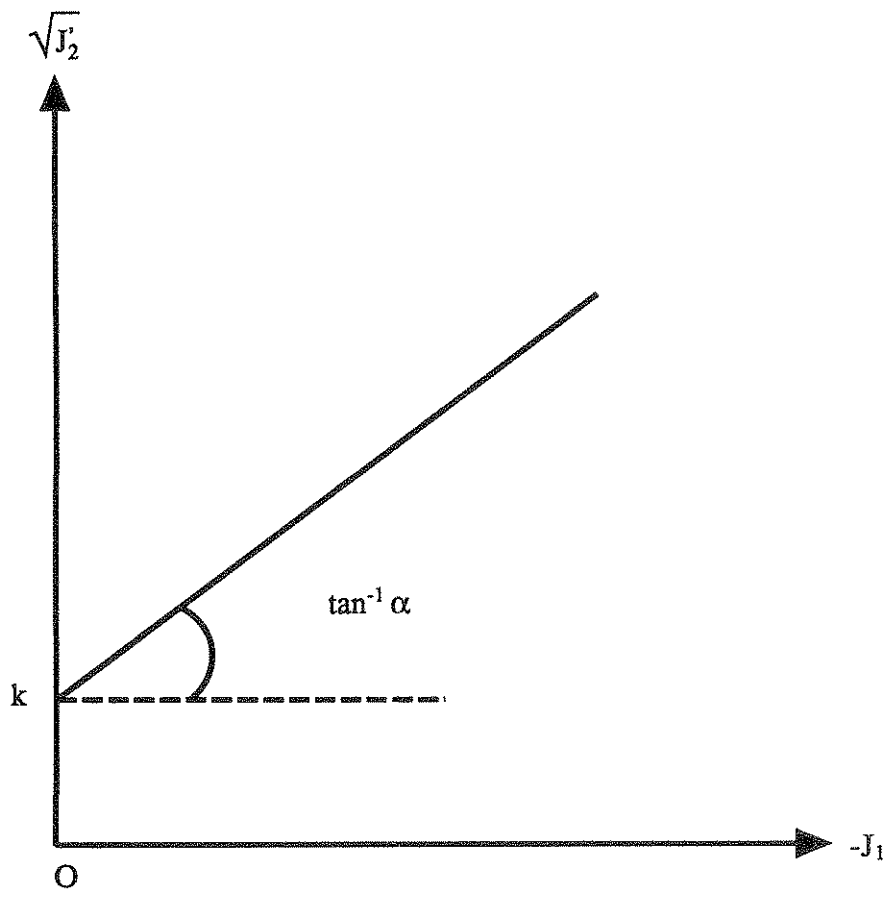


FIGURE 2-3 Drucker - Prager Yield Surface

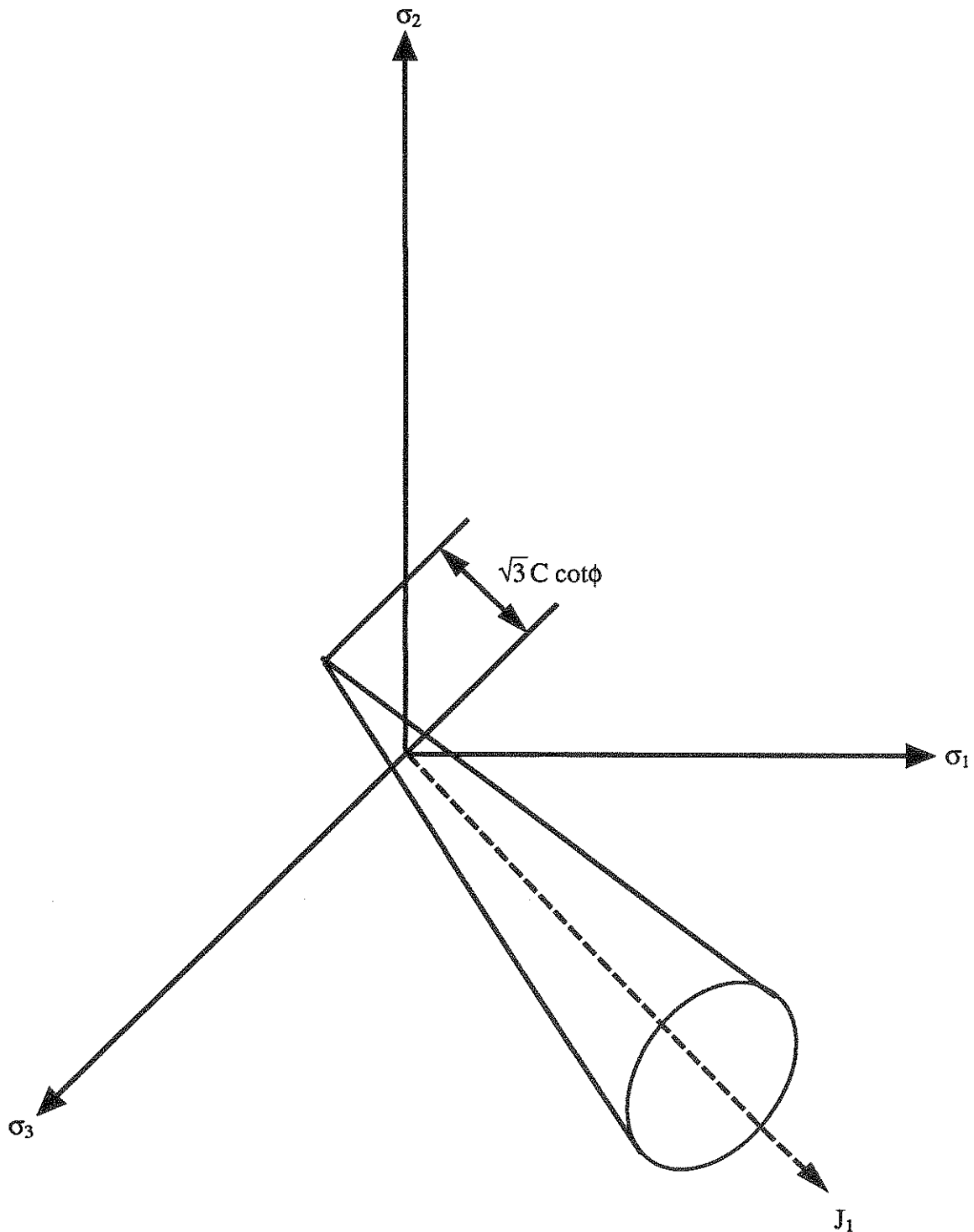


FIGURE 2-4 Drucker - Prager Yield Surface (Cone)

2. Experimental observations have shown that considerable hysteresis in a hydrostatic load-unloading path occurs which cannot be predicted using the same elastic bulk modulus of loading and unloading and a yield surface which does not cross the hydrostatic (J_1) axis; and
3. Soils pass through the fluid state at high pressures where shear strength does not vary with hydrostatic pressure. Therefore, the yield condition should essentially be independent of J_1 for large J_1 .

2.4 The Development of the Cap Model

To control the plastic volumetric change or dilatation of soils, Drucker, Gibson and Henkel (Ref. 8) added a movable cap to the Drucker-Prager model, which crosses the hydrostatic loading axis. This modification reproduces better the hysteresis which soils display when loaded and unloaded hydrostatically.

Several strain hardening plasticity models based on critical-state concepts have since then been developed by a group at Cambridge University.

A group of researchers from MIT also worked on a similar model where the yield curves are ellipses of constant eccentricity. However, Drucker-Prager's postulate of stability in the small is not fully satisfied at all points on the yield surface and, therefore, the irreversibility condition

$$\sigma_{ij} de_{ij}^p \geq 0 \quad (2-15)$$

is not satisfied.

DiMaggio and Sandler (Ref. 27) have proposed a new cap model for granular soil which satisfies continuity, stability, and uniqueness conditions.

The new cap model has an ideally plastic modified Drucker-Prager yield condition denoted by

$$f_1 \left(J_1, \sqrt{J_2'} \right) = 0 \quad (2-16)$$

and a strain-hardening cap, which expands or contracts as the plastic volumetric strain decreases or increases, respectively, denoted by

$$f_2 \left(J_1, \sqrt{J_2'}, \varepsilon_v^p \right) = 0 \quad (2-17)$$

A full discussion of the above model and its application to McCormick Ranch sand is covered in a different section.

Modifications to the cap model to include kinematic hardening has been made for materials whose hysteresis is independent of strain rate (Ref. 11). This modification is achieved by replacing the stress tensor σ_{ij} by the quantity $(\sigma_{ij} - \alpha_{ij})$, where α_{ij} is the tensor whose components are memory parameters defining the translation of the yield surface in stress space. The kinematic hardening is assumed to occur in shear only. Hence

$$\alpha_{ii} = 0 \quad (2-18)$$

where the summation convention of the subscripts apply. The memory parameter α_{ij} is governed by a kinematic hardening rule given by

$$\alpha_{ij} = C_\alpha \dot{\varepsilon}_{ij}^p \quad (2-19)$$

where $\dot{\varepsilon}_{ij}^p$ are the deviatoric components of plastic strain, and C_α is a constant. Extensive coverage of this subject has been made in Ref. 11.

2.5 Bounding Surface Plasticity Theory

The bounding surface (Ref. 19) is represented in the 2-D octahedral, shear-vs-normal stress space by a half-ellipse with a variable axis length (Figure 2-5). The ellipse's horizontal axis coincides with the normal stress axis and has one of its end points fixed to the origin of coordinate system,

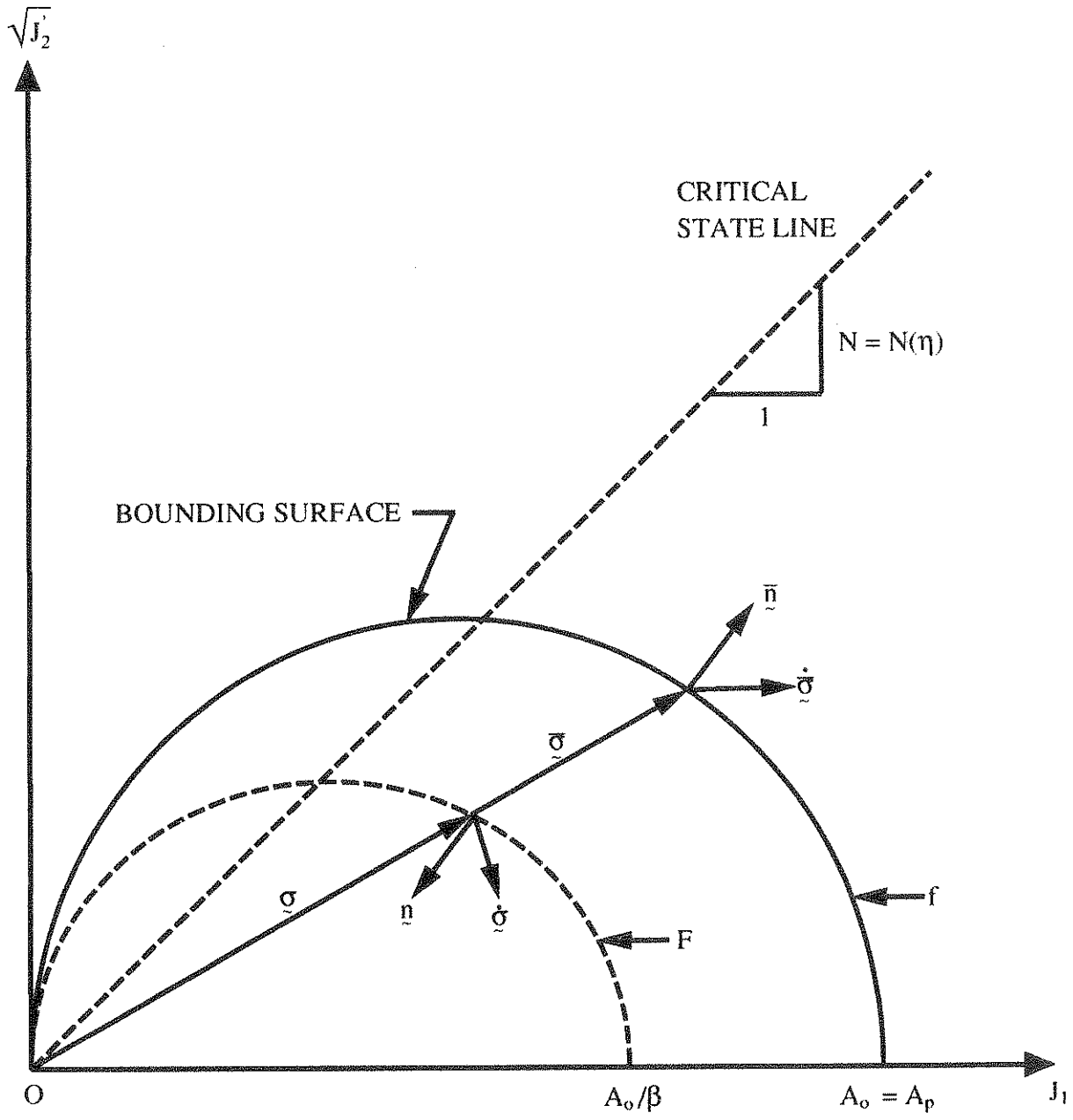


FIGURE 2-5 Bounding Surface

and the other end point, as well as the ratio of the axis lengths, is determined by the hardening law. The hardening law is obtained along the standard triaxial stress path. An ellipse is formulated such that:

1. It contains the stress point representing the state of stress in the triaxial test sample at all times; and
2. The "vector" normal at this stress point is proportional to the plastic strain rate "vector" measured in the triaxial test. The major axis and the ratio between the axes are then expressed as functions of the stress level and accumulated plastic strain. Under this new hardening law, the bounding surface, for any chosen loading path, always expands.

Elasto-Plastic Constitutive Equations

$$\dot{\epsilon}^p = \frac{\underline{n} : \dot{\sigma}}{k_p} \underline{n} \quad (2-20)$$

where:

$\langle \rangle = \text{Operation}$ $\langle L \rangle = LH(L)$

$H = \text{Heaviside's step function}$

$\dot{\epsilon}^p = \text{Plastic strain rate tensor}$

$k_p = \text{Generalized plastic modulus}$

$\underline{n} = \text{Second order tensor such that } \underline{n} : \underline{n} = 1$

$$\dot{\epsilon} = \underline{D}^{-1} : \dot{\sigma} - \left\langle \frac{\underline{n} : \dot{\sigma}}{k_p} \right\rangle \underline{n} \quad (2-21)$$

or inverted

$$\dot{\underline{\sigma}} = \underline{D}:\dot{\underline{\epsilon}} - \left\langle \frac{\underline{n}:\underline{D}:\dot{\underline{\epsilon}}}{k_p + \underline{n}:\underline{D}:\underline{n}} \right\rangle \underline{D}:\underline{n} \quad (2-22)$$

where $\dot{\underline{\epsilon}}$ is the total strain rate tensor.

Radial Mapping

The radial projection of a stress point representing $\underline{\sigma}$ produces $\bar{\underline{\sigma}}$ on the bounding surface, where

$$\bar{\underline{\sigma}} = \beta \underline{\sigma} \quad (2-23)$$

with β , the radial mapping scalar, being unity when $\underline{\sigma}$ lies on the bounding surface itself.

Loading Condition

$$\text{Loading path: } \underline{n} = \bar{\underline{n}} ; k_p = \bar{k}_p \text{ with } \beta = 1 \quad (2-24)$$

$$\text{unloading path: } \underline{n} = -\bar{\underline{n}} ; k_p = H_u / (\beta - 1) \quad (2-25)$$

with $\dot{\underline{\sigma}}$ pointing inwards from F

$$\text{reloading path: } \underline{n} = \bar{\underline{n}} ; k_p = H_R (\beta - 1) \quad (2-26)$$

with $\dot{\underline{\sigma}}$ pointing outwards from F where H_u and H_R are material parameters

$$\bar{\underline{n}} = \frac{\underline{\nabla}f}{\|\underline{\nabla}f\|} \quad (2-27)$$

where the operator $\underline{\nabla}$ is defined in the stress space and $\underline{\nabla}f$ is evaluated at the stress state $\bar{\underline{\sigma}}$.

$$\bar{k}_p = - \frac{1}{\|\nabla f\|} \frac{\partial f}{\partial \eta} \sqrt{\frac{1}{2} - \frac{1}{6} [\text{tr}(\eta)]^2} \quad (2-28)$$

where η is the equivalent shear strain defined by the hardening law

$$\eta = \int \dot{\eta} = \int \frac{1}{2} = \dot{\epsilon}^p \dot{\epsilon}^p \quad (2-29)$$

and $\dot{\epsilon}^p$ is the deviator of $\dot{\epsilon}^p$. The link between the generalized plastic moduli k_p and \bar{k}_p is given by

$$\frac{\underline{n}:\dot{\underline{\sigma}}}{k_p} = \frac{\bar{n}:\dot{\underline{\sigma}}}{\bar{k}_p} \geq 0 \quad (2-30)$$

Bounding Surface

The bounding surface f is defined as

$$f = \frac{\bar{J}_2}{J_1 N^2} + \bar{J}_1 - A_0 = 0 \quad (2-31)$$

where:

$$\bar{J}_2 = \beta^2 J_2' \quad (2-32)$$

and

$$\bar{J}_1 = \beta J_1 \quad (2-33)$$

The strain dependency of the hardening surface f is defined by the slope $N = N(\eta)$, of the critical state line, and the axis, $A_0 = pA$ of the half-ellipse; where $A = A(\eta)$ and p is the atmospheric pressure.

The bounding surface parameters A and N are related to the shear dilatation angle α by

$$\tan \alpha = \frac{\bar{J}_1 \sqrt{\bar{J}_2}}{3(N^2 \bar{J}_1^2 - \bar{J}_2')} \quad (2-34)$$

$$= \frac{\sqrt{\frac{1}{2} \bar{\eta}^d : \bar{\eta}^d}}{\text{tr}(\bar{\eta})}$$

where $\bar{\eta}^d = \bar{\eta} - \frac{1}{3} \text{tr}(\bar{\eta})\delta$. The main feature of this 8-parameter model is its hardening law which is defined along the standard triaxial test path to emphasize the shear dilatation of sands. The model's ability to simulate important characteristics of sand under cyclic loading compares well with laboratory tests.

2.6 The Endochronic Theory

The Endochronic Theory is based on the hypothesis that the current state of stress is a linear function of the entire history of plastic strain, with the history defined with respect to a time scale (intrinsic time) which is itself a property of the material at hand.

The equations governing the Endochronic Theory are as follows:

$$\sigma_{ij} = H_D \int_0^{Z_D} \rho(Z_D - Z') \frac{\partial \theta_{ij}}{\partial Z'} dZ' \quad (2-35)$$

$$+ \delta_{ij} H_H \int_0^{Z_H} \phi(Z_H - Z') \frac{\partial \theta}{\partial Z'} dZ'$$

$$d\sigma_{ij} = [K_o (d\varepsilon - d\theta) - 2/3 G_o d\varepsilon] \delta_{ij} \quad (2-36)$$

$$+ 2G_o (d\varepsilon_{ij} - d\theta_{ij})$$

where the intrinsic time scales, Z_D and Z_H , are positive, monotonically increasing quantities defined by the ff expressions:

$$dZ_D^2 = k_{oo} d\zeta_D^2 + k_{oi} d\zeta_H^2 \quad (2-37)$$

$$dZ_H^2 = k_{io} d\zeta_D^2 + k_{ii} d\zeta_H^2 \quad (2-38)$$

where:

H_D = Hardening function

H_H = Softening function

$d\theta_{ij}$ and $d\theta$ are defined such that

$$d\epsilon_{ij}^p = d\theta_{ij} + \frac{1}{3} d\theta \delta_{ij} \quad (2-39)$$

or

$$d\theta = d\epsilon_{kk}^p \quad (2-40)$$

Also K_o G_o denote the bulk and elastic shear moduli, respectively. $\rho(z)$ and $\phi(z)$ are material functions. The intrinsic time measures $d\zeta_D$ and $d\zeta_H$ are defined as follows.

$$d\zeta_D^2 = d\theta_{ij} d\theta_{ij} \quad (2-41)$$

$$d\zeta_H^2 = |d\theta|^2 \quad (2-42)$$

and the k_{rs} are elements of the material-dependent coupling matrix $[k]$. It is evident that the materials described by the above equations are plastic strain history dependent but strain rate

independent (i.e., are independent of the natural time scale given by a clock). The following kernel functions may be used to represent soils using the Endochronic Theory

$$\rho(z) = \frac{e^{-kz}}{\sqrt{z}} \quad (2-43)$$

$$\phi(z) = \frac{\phi_0}{\sqrt{z}} \quad (2-44)$$

$$H_D(z) = H_0 + (H_\infty - H_0)(1 - e^{-\gamma z}) \quad (2-45)$$

$$H_H(\theta) = (\theta_m - \theta)^{-\beta} \quad (2-46)$$

where k , H_0 , H_∞ , γ , θ_m and β are constants. For numerical purposes, the incremental form of equation (2-35) may be used as follows

$$\begin{aligned} d\sigma_{ij} = & \left[H_D \int_0^{Z_D} \rho(Z_D - Z') \frac{\partial^2 \theta_{ij}}{\partial Z'^2} dZ' \right] dZ_D \\ & + \delta_{ij} \left[H_H \int_0^{Z_H} \phi(Z_H - Z') \frac{\partial^2 \theta}{\partial Z'^2} dZ' \right] dZ_H \end{aligned} \quad (2-47)$$

2.7 The Hyperbolic Stress-Strain Law

Considered as a hybrid of plasticity theory and endochronic theory, the hyperbolic stress-strain law was proposed by Kondner and his co-workers to model nonlinear stress-strain relations in soil (Ref. 22). The nonlinear stress-strain curve is represented by a hyperbola of the form

$$(\sigma_1 - \sigma_3) = \frac{\epsilon_a}{a + b\epsilon_a} \quad (2-48)$$

in which $(\sigma_1 - \sigma_3)$ is the principal stress difference, ϵ_a is the axial strain, and a and b are

parameters whose values are determined experimentally. As shown in Figure 2-6, these parameters are the reciprocals of the initial slope (initial tangent modulus) and the asymptote to the $\sigma - \epsilon$ curve. To determine the parameters a and b , equation (2-48) may be transformed into the form

$$\frac{\epsilon_a}{(\sigma_1 - \sigma_3)} = a + b\epsilon_a \quad (2-49)$$

As shown in Figure 2-7, the parameters a and b are, respectively, the intercept and the slope of the straight line. The ratio

$$\frac{(\sigma_1 - \sigma_3)_{\text{failure}}}{(\sigma_1 - \sigma_3)_{\text{ultimate}}} = R_f \quad (2-50)$$

where $R_f < 1$, and is a correlation factor called "failure ratio." Values of R_f for a variety of different soils have been found to range from 0.5 to 1.0 and to be essentially independent of confining pressure.

2.8 Earlier Viscoplastic Models

Viscoplasticity is a term which has been used by researchers to describe material behavior whereby all plastic strains in the material are developed with time; that is, there is a delayed plasticity, in contrast with elasto-plastic strains, which are produced instantaneously, i.e. are independent of time. This section outlines some of the earlier viscoplastic models used for time-dependent materials.

Rheology

Consider the rheological model of an elasto-viscoplastic material shown in Figure 2-8. It consists of a spring which is in series with a slider and dashpot system in parallel. The spring gives the elastic response while the dashpot and slider allow viscous deformation only for stresses greater than a certain limit.

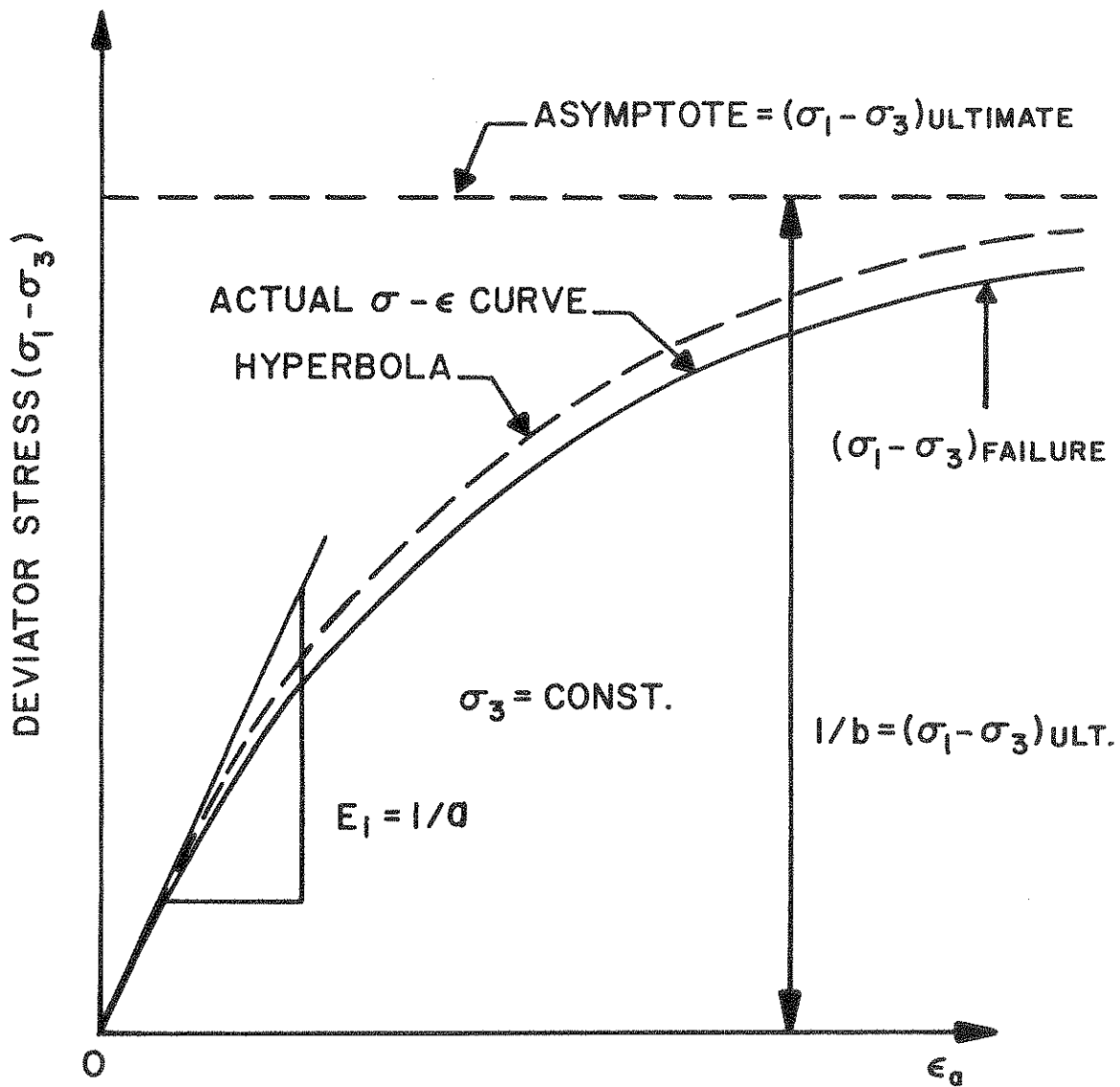


FIGURE 2-6 Hyperbolic Representation Of Stress-Strain Curve

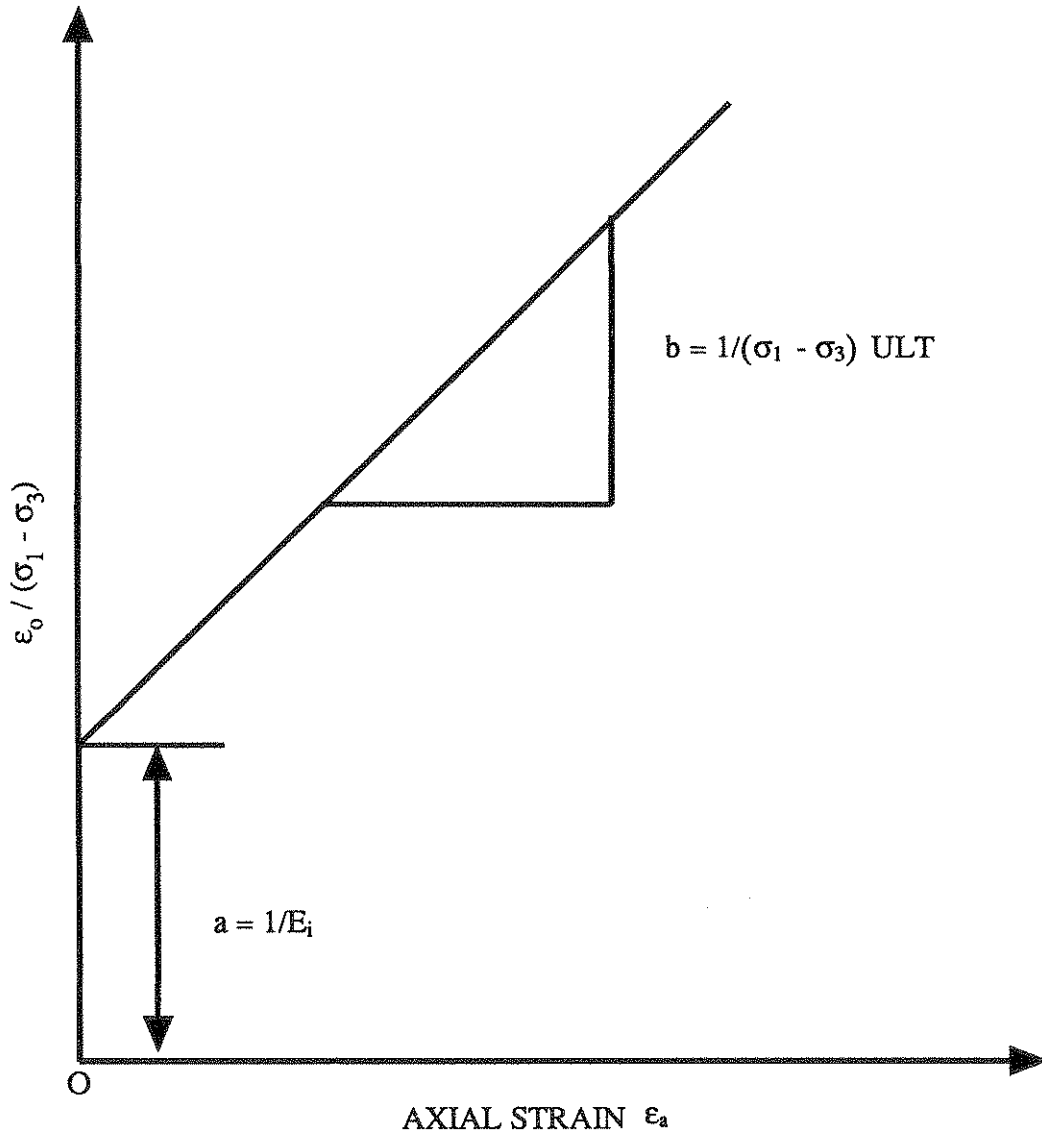


FIGURE 2-7 Transformed Hyperbolic Representation of Stress-Strain Curve

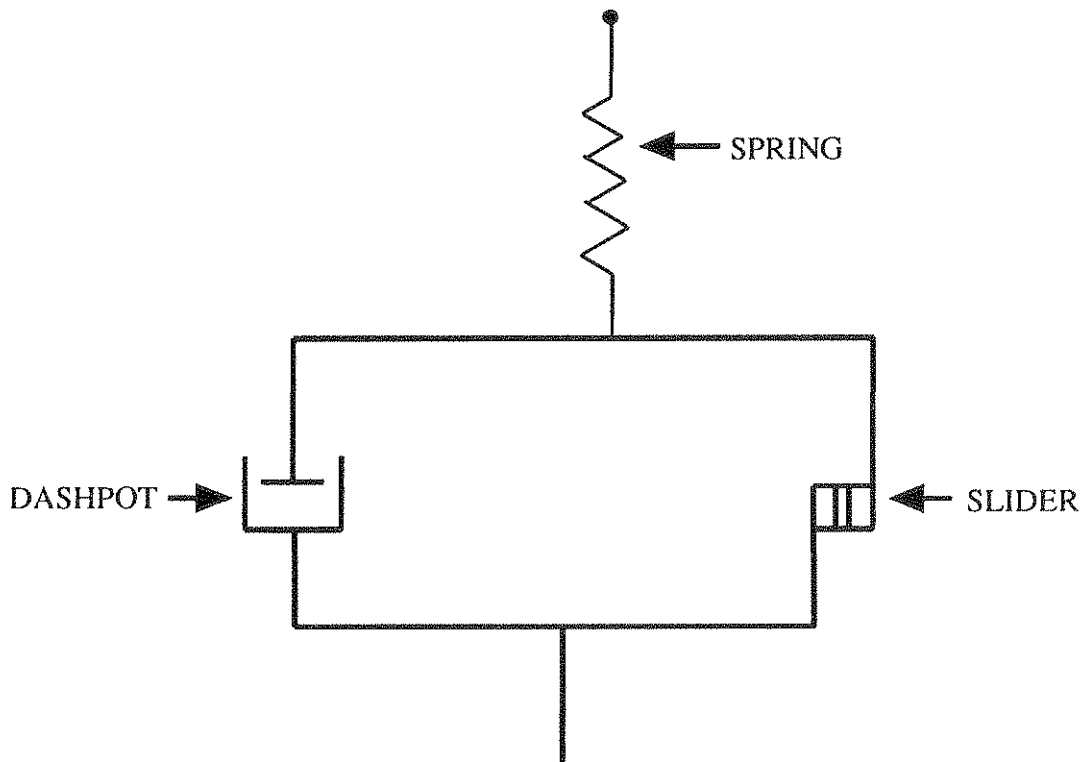


FIGURE 2-8 Rheological Model Of Elasto-Viscoplasticity

Figure 2-9 shows a stress-strain curve in 1-D to illustrate the behavior of an elasto-viscoplastic model. If $\sigma \leq \sigma_y$ where σ_y is the yield stress of the material, no viscoplastic strains are developed. However, if $\sigma > \sigma_y$ viscoplastic strains are developed at a finite rate which depends on the excess of $\sigma - \sigma_y$. A steady state is reached at time $t = T$ when the stress is on the yield surface and no further increase in viscoplastic strains occur. A full discussion on the development of viscoplastic strains is covered in later sections in this report.

2.8.1. The Rigid-Viscoplastic Models

In rigid-viscoplastic models, materials show rigid behavior (no deformation) when applied stress is below a certain limit and show viscous response when the limit is exceeded. These are referred to as Bingham materials.

For such materials, the constitutive equation can be expressed in the form

$$s_{ij} = \frac{\dot{\epsilon}_{ij}}{2\lambda} + 2 \eta \dot{\epsilon}_{ij} \quad (2-51)$$

where s_{ij} is the deviatoric stress given by

$$s_{ij} = \sigma_{ij} - 1/3 \sigma_{kk} \delta_{ij} \quad (2-52)$$

$\dot{\epsilon}_{ij}$ is the deviatoric strain rate given by

$$\dot{\epsilon}_{ij} = \dot{\epsilon}_{ij} - 1/3 \dot{\epsilon}_{kk} \delta_{ij} \quad (2-53)$$

where:

η = Viscosity coefficient

λ = Scaler multiplier

Consider the von Mises yield condition

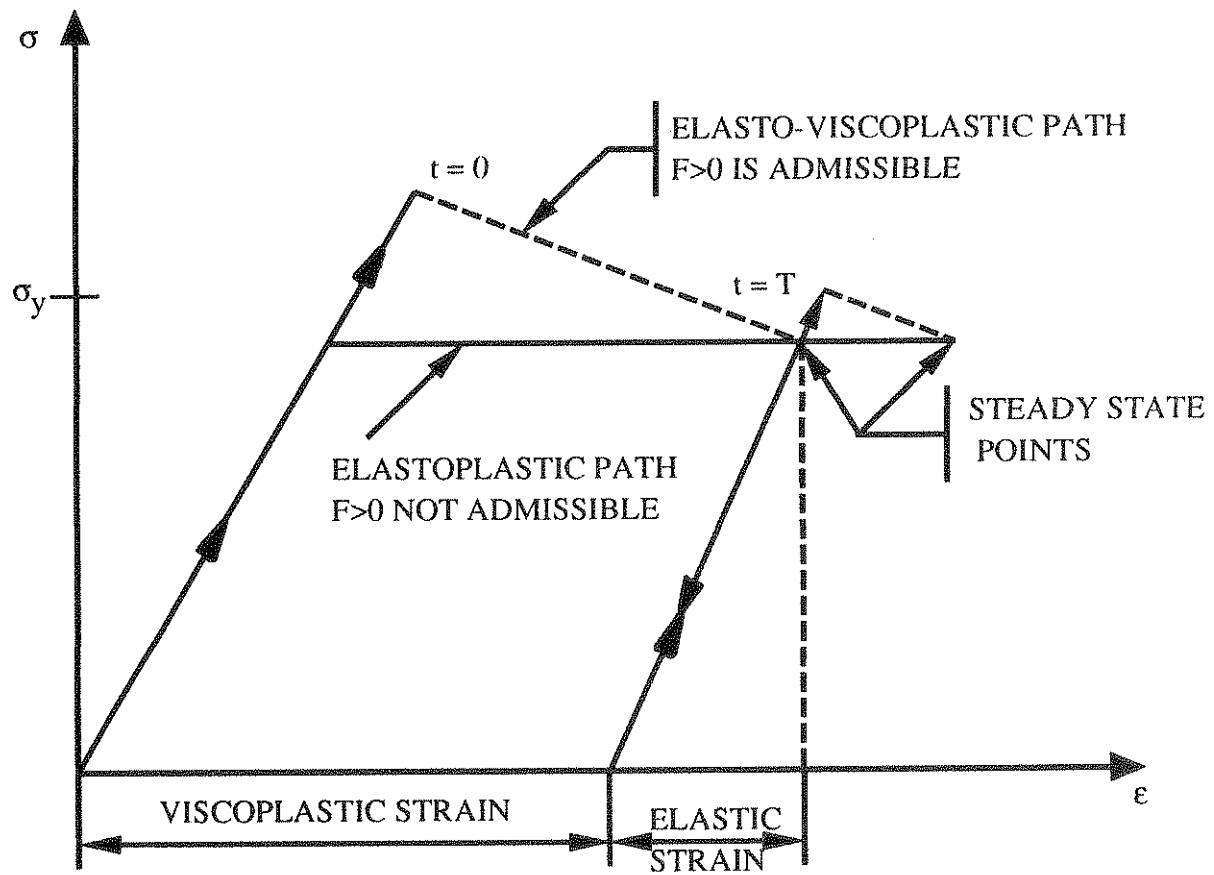


FIGURE 2-9 Typical Stress-Strain Curve For Elasto-Viscoplastic Model

$$J_2' = k^2 \quad (2-54)$$

where J_2' is the second invariant of deviatoric stress given by

$$J_2' = \frac{1}{2} s_{ij} s_{ij} \quad (2-55)$$

and k is the yield stress in pure shear, and rewrite equation (2-51) in the form

$$s_{ij} = \left(\frac{1}{2\lambda} + 2\eta \right) \dot{e}_{ij} = 2\eta' \dot{e}_{ij} \quad (2-56)$$

where η' is a variable viscosity coefficient.

From equation (2-54) we can obtain

$$2\eta' = 2\eta \left(1 + \frac{k}{2\eta\sqrt{I_e'}} \right) = \frac{\sqrt{J_2'}}{\sqrt{I_e'}} = \frac{2\eta}{1 - k/\sqrt{J_2'}} \quad (2-57)$$

where:

$$I_e' = \frac{1}{2} \dot{e}_{ij} \dot{e}_{ij}$$

Generally $\eta' > \eta$ and $\eta' = \eta$ only when $\sqrt{J_2'} \rightarrow \infty$. From equation (2-57) we can rewrite equation (2-51) in the form

$$s_{ij} = 2\eta \left(1 - \frac{k}{\sqrt{2J_2'}} \right)^{-1} \dot{e}_{ij} \text{ if } J_2' > k^2 \quad (2-58)$$

For the rate of work to be positive during plastic deformation

$$s_{ij} \dot{e}_{ij} = 2k\sqrt{I_e} + 4\eta I_e' > 0 \quad (2-59)$$

This, however, becomes zero as $J_2' \rightarrow k^2$.

The volume remains incompressible during deformation, i.e., $\dot{e}_{ii} = 0$. For $J_2' \leq k^2$ the body is rigid.

2.8.2 Elasto-Viscoplastic Models

In elasto-viscoplastic models, the material behavior is elastic when applied stress is below a certain limit, and shows viscous response under the action of higher stress.

The constitutive equation for such a material behavior is composed of elastic and viscoplastic components.

We have

$$\dot{e}_{ij} = \dot{e}_{ij}^E + \dot{e}_{ij}^{VP} \quad (2-60)$$

or

$$\dot{e}_{ij} = \frac{\dot{s}_{ij}}{2G} + \frac{2\lambda}{1+4\lambda\eta} s_{ij} \quad (2-61)$$

or from equation (2-58)

$$\dot{e}_{ij} = \frac{\dot{s}_{ij}}{2G} + \frac{1}{2\eta} \left(1 - \frac{k}{\sqrt{J_2'}} \right) s_{ij} \text{ if } J_2' > k^2 \quad (2-62)$$

and $\sigma_{kk} = 3K\varepsilon_{kk}$ (elastic volumetric behavior). In the above equations, $\dot{\epsilon}_{ij}$, $\dot{\epsilon}_{ij}^E$, $\dot{\epsilon}_{ij}^{VP}$ are respectively the total, elastic, and viscoplastic strain rates, and G is the shear modulus.

Equation (2-62) can be written more generally as

$$\dot{\epsilon}_{ij} = \frac{\dot{s}_{ij}}{2G} + \frac{1}{2\eta} \left\langle 1 - \frac{k}{\sqrt{J_2}} \right\rangle s_{ij} \quad (2-63)$$

where the notation $\langle \rangle$ is defined such that

$$\begin{aligned} \langle F \rangle &= 0 \text{ if } F \leq 0 \\ \langle F \rangle &= F \text{ if } F > 0 \end{aligned} \quad (2-64)$$

where F is an arbitrary function. The above model can be viewed as a modification of the Prandtl-Reuss model to include the viscous effect of materials.

Loading Condition

The loading condition may be defined locally using the rate of working. Now

$$\dot{W} = \dot{W}^E + \dot{W}^{VP} \quad (2-65)$$

where \dot{W} , \dot{W}^E , \dot{W}^{VP} denote, respectively, the total rate of working, the rate of work due to elastic and viscoplastic deformations. Hence

$$s_{ij} \dot{\epsilon}_{ij} = s_{ij} \dot{\epsilon}_{ij}^E + s_{ij} \dot{\epsilon}_{ij}^{VP} \quad (2-66)$$

or

$$s_{ij} \dot{\epsilon}_{ij} = \frac{1}{2G} \left(\dot{J}_2 + 2\frac{G}{\eta} \left(1 - \frac{k}{\sqrt{J_2}} \right) J_2 \right) \quad (2-67)$$

Notice that $\dot{W}^{VP} > 0$ if $J'_2 > k^2$. If

$$\dot{J}'_2 > k^2 \text{ and } s_{ij} \dot{e}_{ij} > 0 \quad (2-68)$$

then we have loading. There are three types of loading:

1. If in addition to equation (2-68)

$$\dot{J}'_2 > 0, \text{ and therefore } \dot{W}^{VP} > 0, \dot{W}^E > 0 \quad (2-69)$$

the process is called total loading.

2. If in addition to equation (2-68)

$$\dot{J}'_2 = 0, \text{ and therefore } \dot{W}^{VP} > 0, \dot{W}^E = 0 \quad (2-70)$$

then we have neutral loading.

3. If in addition to equation (2-68)

$$\dot{J}'_2 < 0 \text{ but } \dot{W} > 0 \quad (2-71)$$

we have partial loading. If

$$\begin{aligned} J'_2 > k^2 \text{ and } s_{ij} \dot{e}_{ij} &= 0 \\ \text{i.e. when } -\dot{W}^E = \dot{W}^{VP} &> 0 \end{aligned} \quad (2-72)$$

then we have a state of pure relaxation. If

$$J_2' > k^2 \text{ and } \dot{s}_{ij} \dot{e}_{ij} < 0 \quad (2-73)$$

the process is called quasi-unloading because the decrease of the stress is more rapid than in a relaxation process, and therefore corresponds to stress decrease at the boundary of the body, while still $\dot{W}^{VP} > 0$. A state of pure unloading (instantaneous) corresponds to $\dot{J}_2' = \infty$.

During stress relaxation at constant strain, equation (2-67) becomes

$$\dot{J}_2' + \frac{2G}{\eta} \left(1 - \frac{k}{\sqrt{J_2'}} \right) J_2' = 0 \quad (2-74)$$

which has a solution of the form

$$\sqrt{J_2'}(t) = k + \left(\sqrt{J_2'^0} - k \right) \exp \left[- \frac{G}{\eta} (t-t_0) \right] \quad (2-75)$$

where:

$$\sqrt{J_2'^0} = \sqrt{J_2'}(t_0)$$

and t_0 is the reference time at which relaxation process begins. Now when $t \rightarrow \infty \sqrt{J_2'} \rightarrow k^+$ (from above)

$$\lim_{t \rightarrow \infty} \sqrt{J_2'}(t) = k^+ \quad (2-76)$$

In other words, as the time approaches infinity, the stress deviator relaxes towards a point on the yield surface $J_2' = k^2$.

Perzyna (Ref. 6) has introduced more parameters in the visco-plastic model including hardening properties. He postulated a constitutive equation of the form

$$\dot{\epsilon}_{ij} = \frac{\dot{s}_{ij}}{2G} + \gamma \langle \phi(F) \rangle \frac{\partial f}{\partial \sigma_{ij}} \quad (2-77)$$

$$\dot{\epsilon}_{ii} = \frac{\dot{\sigma}_{ii}}{3K} \quad (2-78)$$

where:

$$F = F(\sigma_{ij}, \epsilon_{ij}^p) = f\left(\frac{\sigma_{ij}, \epsilon_{ij}^p}{\kappa}\right) - 1 \quad (2-79)$$

is the statical yield function. G and K are, respectively, the elastic shear and bulk moduli, and γ is a material constant. The symbol $\langle \phi(F) \rangle$ is defined as follows

$$\langle \phi(F) \rangle = \begin{cases} 0 & \text{for } F \leq 0 \\ \phi(F) & \text{for } F > 0 \end{cases} \quad (2-80)$$

The function $f(\sigma_{ij}, \epsilon_{ij}^p)$ depends on the state of stress σ_{ij} and on the state of anelastic strain ϵ_{ij}^p and

$$\kappa = \kappa(W_p) = \kappa\left(\int_{\sigma}^{\epsilon_{ij}^p} \sigma_{ij} d\epsilon_{ij}^p\right) \quad (2-81)$$

is the work-hardening parameter. Consider the anelastic part of the constitutive equation (2-77)

$$\dot{\epsilon}_{ij}^p = \gamma \phi(F) \frac{\partial f}{\partial \sigma_{ij}} \quad (2-82)$$

squaring both sides of equation (2-82) we obtain

$$\dot{\epsilon}_{ij}^p \dot{\epsilon}_{ij}^p = \gamma^2 \left\{ \phi(F) \frac{\partial f}{\partial \sigma_{ij}} \right\}^2 \quad (2-83)$$

Now let

$$I_2^p = \frac{1}{2} \dot{\epsilon}_{ij}^p \dot{\epsilon}_{ij}^p \quad (2-84)$$

where I_2^p is the second invariant of the plastic strain rate. Then equation (2-83) becomes

$$2I_2^p = \gamma^2 \left\{ \phi(F) \frac{\partial f}{\partial \sigma_{ij}} \right\}^2 \quad (2-85)$$

or

$$F - \phi^{-1} \left(\frac{I_2^p}{\frac{1}{2} \gamma^2 \frac{\partial f}{\partial \sigma_{ij}} \frac{\partial f}{\partial \sigma_{ij}}} \right)^{\frac{1}{2}} = 0 \quad (2-86)$$

where ϕ^{-1} denotes the functional inverse of ϕ or

$$f(\sigma_{ij}, \epsilon_{ij}^p) = \kappa(W_p) \left\{ 1 + \phi^{-1} \left[\frac{I_2^p}{\frac{1}{2} \gamma^2 \frac{\partial f}{\partial \sigma_{kl}} \frac{\partial f}{\partial \sigma_{kl}}} \right]^{\frac{1}{2}} \right\} \quad (2-87)$$

This expression implicitly represents dynamical yield condition for elasto-viscoplastic work hardening materials and also describes the dependence of the yield criterion on the strain rate. From equation (2-82) it is evident that the plastic strain rate vector $\dot{\epsilon}_{ij}^p$ in the 9-D stress hyper-space is always along the normal to the subsequent loading surface. (See Figure 2-10.) The yield theory described above reduces to classical rate independent (inviscid) theory for vanishingly small strain rates.

Different forms of the function $\phi(F)$ in equation (2-77) have been proposed by Perzyna in Ref. 6 as

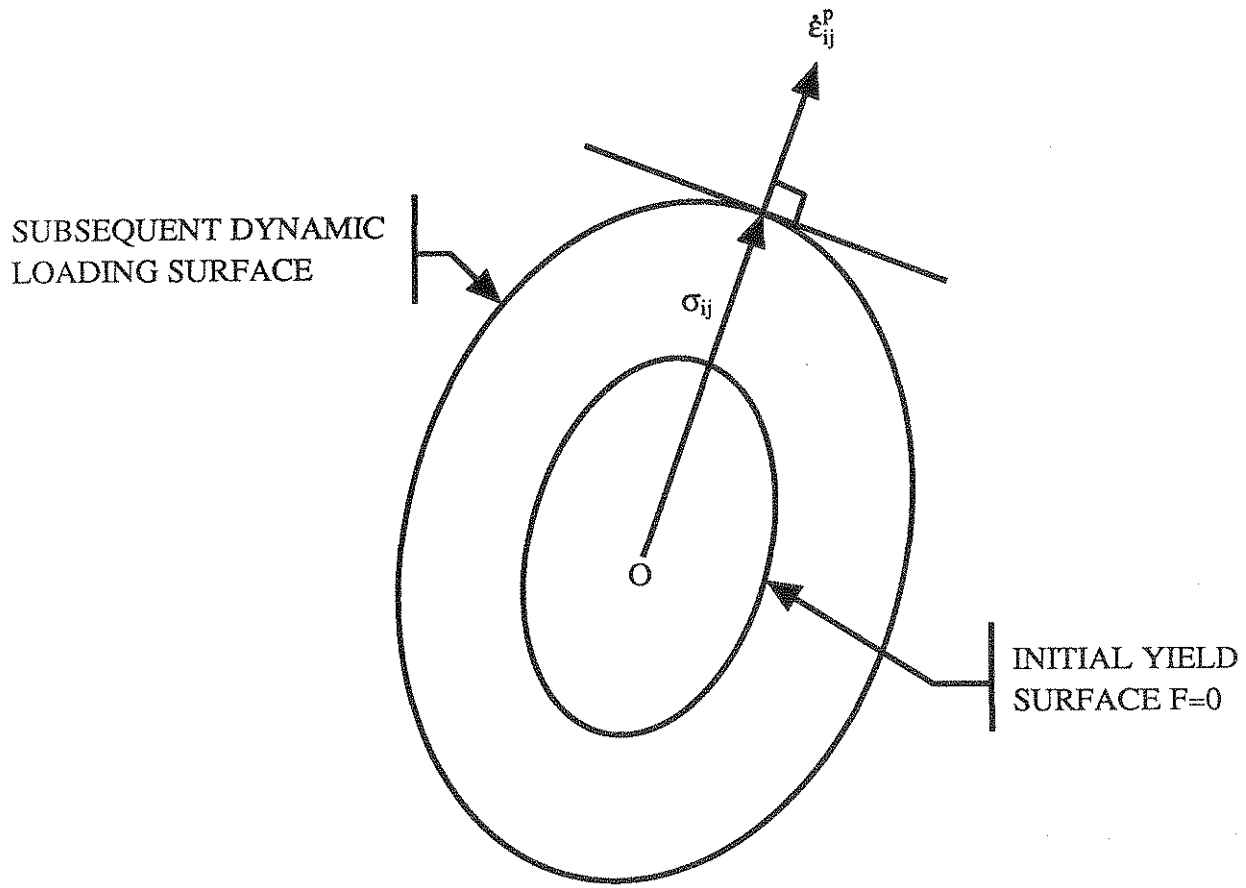


FIGURE 2-10 Typical Yield Surface For An Elasto-Viscoplastic Material

$$\phi(F) = F^\delta$$

$$\phi(F) = F$$

$$\phi(F) = \exp F - 1 \quad (2-88)$$

$$\phi(F) = \sum_{\alpha=1}^N A_\alpha [\exp F^\alpha - 1]$$

$$\phi(F) = \sum_{\alpha=1}^N B_\alpha F^\alpha$$

The above functions, however, describe perfectly plastic rate sensitive material under dynamic loading.

Hohenemser and Prager have proposed an elasto-viscoplastic constitutive equation of the form

$$\dot{\epsilon}_{ij} = \gamma [F_M] \frac{s_{ij}}{\sqrt{J_2}} \quad (2-89)$$

and

$$\sqrt{J_2'} = k_0 \left(1 + \frac{\dot{I}_2}{\gamma} \right) \quad (2-90)$$

where k_0 is the yield stress in simple shear and

$$\dot{I}_2 = \frac{1}{2} \dot{\epsilon}_{ij} \dot{\epsilon}_{ij} \quad (2-91)$$

γ is a material constant and

$$F_M = \frac{\sqrt{J_2}}{k_0} - 1 \quad (2-92)$$

To generalize equation (2-89) to include the effect of work hardening, Rosenblatt (Ref. 12) proposed a constitutive equation of the form

$$\dot{\epsilon}_{ij} = \gamma [\phi(F)] \frac{s_{ij}}{\sqrt{J_2}} \quad (2-93)$$

This relation, however, violates the normality requirement sufficient for uniqueness, in the limit of vanishing inelastic strain rates except for the special case where the expression is governed by equation (2-92).

2.8.3 Elastic-Viscoplastic Models

In these models, the material behavior is elastic if stresses are below a certain limit. If stresses exceed this limit, material exhibits instantaneous plastic deformation in addition to a delayed (viscous) deformation. The total strain rate is given by

$$\dot{\epsilon}_{ij} = \dot{\epsilon}_{ij}^E + \dot{\epsilon}_{ij}^P + \dot{\epsilon}_{ij}^{VP} \quad (2-94)$$

where $\dot{\epsilon}_{ij}^E$, $\dot{\epsilon}_{ij}^P$, $\dot{\epsilon}_{ij}^{VP}$ denote, respectively, the elastic, plastic, viscoplastic strain rates.

Yannis F. Dafalias (Ref. 20) has used the above postulate to model the behavior of cohesive soils. Dafalias used constitutive equations of the form

$$\begin{aligned} \dot{\epsilon}_{ij} = & C_{ijkl} \dot{\sigma}_{kl} \\ & + \left\langle \frac{1}{\kappa_p} \frac{\partial f}{\partial \sigma_{ij}} \dot{\sigma}_{ij} \right\rangle \frac{\partial f}{\partial \sigma_{ij}} + \langle \phi(\Delta \hat{\sigma}) \rangle R_{ij}^v \end{aligned} \quad (2-95)$$

where:

C_{ijkl} = Elastic compliance

$f(\sigma_{ij}, q_n^p) = 0$ = Yield surface

$$\dot{q}_n^p = \left\langle \frac{1}{\kappa_p} \frac{\partial f}{\partial \sigma_{ij}} \dot{\sigma}_{ij} \right\rangle r_n^p \quad (2-96)$$

(the plastic modulus)

$$\kappa_p = - \frac{\partial f}{\partial q_n^p} r_n^p \quad (2-97)$$

and r_n^p is a function of the state only, $\kappa_p > 0$ denotes stable response and $\kappa_p \leq 0$ denotes unstable response, $\Delta \hat{\sigma}$ is overstress. The notation $\langle \rangle$ is defined such that

$$\langle F \rangle = 0 \text{ if } F \leq 0 \text{ and}$$

$$\langle F \rangle = F \text{ if } F > 0$$

From normality principle

$$R_{ij}^v = \frac{\partial F}{\partial \bar{\sigma}_{ij}} \quad (2-98)$$

where F defines the bounding surface given by

$$F(\bar{\sigma}_{ij}, q_n^p) = 0 \quad (2-99)$$

The functional form of $F = 0$ can be similar to the form of $f = 0$. Bars over stress quantities indicate points on $F = 0$, and the actual stress lies in or on $F = 0$.

2.9 Viscoelastic-Plastic Models

In the viscoelastic-plastic model (Ref. 5), the total strain is the sum of a viscoelastic component and a plastic component, i.e.

$$e_{ij} = e_{ij}^{VE} + e_{ij}^p \quad (2-100)$$

where e_{ij} , e_{ij}^{VE} , e_{ij}^p denote, respectively, the total, viscoelastic, and plastic strains. The viscoelastic component of strain follows a creep integral law of classical linear viscoelastic theory of the form

$$e_{ij}^{VE}(t) = s_{ij}^0(x)J_1(t) + \int_0^t J_1(t-\tau) \frac{\partial s_{ij}(x,\tau)}{\partial \tau} d\tau \quad (2-101)$$

$$\epsilon_{kk}^{VE}(t) = \sigma_{kk}^0(x)J_2(t) + \int_0^t J_2(t-\tau) \frac{\partial \sigma_{kk}(x,\tau)}{\partial \tau} d\tau \quad (2-102)$$

where J_1 and J_2 denote, respectively, the creep functions in shear and isotropic compression (or dilatation) which may have finite jump discontinuities at $t = 0$, $s_{ij}^0(x)$ stands for $s_{ij}(x,0^+)$, and so on. The initial response of the viscoelastic solid of the type represented by equations (2-101) and (2-102) is assumed to be elastic. That is

$$e_{ij}^E = e_{ij}^{VE}(0) = \frac{s_{ij}}{2G} \quad (2-103)$$

Similarly

$$\epsilon_{kk}^E = \epsilon_{kk}^{VE}(0) = \frac{\sigma_{kk}}{3K} \quad (2-104)$$

Consider the arbitrary 9-dimensional yield surface in stress space

$$f = f(\sigma_{ij}, \epsilon_{ij}^p, \chi_{ij}, \kappa_{ij}) \quad (2-105)$$

$$i = 1,2,3$$

$$j = 1,2,3$$

where:

σ_{ij} = State of stress at a generic point in the stress space

ϵ_{ij}^p = Plastic strain

χ_{ij} = Work-hardening effect due to time history

κ_{ij} = Effect of work-hardening due to path history alone

Define the functional

$$\chi_{ij} = \chi_{ij}(\epsilon_{kl}^V - \epsilon_{kl}^E) \quad (2-106)$$

and

$$\epsilon_{ij}^p = \epsilon_{ij}^p(\sigma_{kl}, \chi_{mn}, \kappa_{pq}) \quad (2-107)$$

where ϵ_{kl}^V denotes time dependent strain tensor and ϵ_{kl}^E denotes elastic strain tensor. Now consider the time rate of f in equation (2-105)

$$\dot{f} = \frac{\partial f}{\partial \sigma_{ij}} \dot{\sigma}_{ij} + \frac{\partial f}{\partial \epsilon_{ij}^p} \dot{\epsilon}_{ij}^p + \frac{\partial f}{\partial \chi_{ij}} \dot{\chi}_{ij} + \frac{\partial f}{\partial \kappa_{ij}} \dot{\kappa}_{ij} \quad (2-108)$$

The loading criterion is as follows: If

$$\frac{\partial f}{\partial \sigma_{ij}} \dot{\sigma}_{ij} + \frac{\partial f}{\partial \chi_{ij}} \dot{\chi}_{ij} < 0 \text{ and } f < 0; \text{ (unloading)} \quad (2-109)$$

If

$$\frac{\partial f}{\partial \sigma_{ij}} \dot{\sigma}_{ij} + \frac{\partial f}{\partial \chi_{ij}} \dot{\chi}_{ij} = 0 \text{ and } f = 0; \text{ (neutral loading)} \quad (2-110)$$

If

$$\frac{\partial f}{\partial \sigma_{ij}} \dot{\sigma}_{ij} + \frac{\partial f}{\partial \chi_{ij}} \dot{\chi}_{ij} > 0 \text{ and } f = 0; \text{ (loading)} \quad (2-111)$$

Constitutive Equations

Consider the yield surface given by

$$f = f(\sigma_{ij}, \kappa_{ij}, \chi) \quad (2-112)$$

where f is the instantaneous yield surface similar to equation (2-105). However, here $\kappa_{ij} = \epsilon_{ij}^p$ and χ_{ij} is replaced by χ .

$$\dot{f} = \frac{\partial f}{\partial \sigma_{ij}} \dot{\sigma}_{ij} + \frac{\partial f}{\partial \epsilon_{ij}^p} \dot{\epsilon}_{ij}^p + \frac{\partial f}{\partial \chi} \dot{\chi} \quad (2-113)$$

Since $\dot{\epsilon}_{ij}^p$ is directed to the normal of the instantaneous loading surface f

$$\dot{\epsilon}_{ij}^p = \begin{cases} \Lambda \frac{\partial f}{\partial \sigma_{ij}} & \text{if } f = 0 \\ 0 & \text{if } f < 0 \end{cases} \quad (2-114)$$

Substitute equation (2-114) into equation (2-113)

$$\frac{\partial f}{\partial \sigma_{ij}} \dot{\sigma}_{ij} + \frac{\partial f}{\partial \epsilon_{ij}^p} \Lambda \frac{\partial f}{\partial \sigma_{ij}} + \frac{\partial f}{\partial \chi} \dot{\chi} = 0 \quad (2-115)$$

Therefore

$$\Lambda = - \frac{\frac{\partial f}{\partial \sigma_{kl}} \dot{\sigma}_{kl} + \frac{\partial f}{\partial \chi} \dot{\chi}}{\frac{\partial f}{\partial \epsilon_{kl}^p} \frac{\partial f}{\partial \sigma_{kl}}} \quad (2-116)$$

Substitute equation (2-116) into equation (2-114)

$$\dot{\epsilon}_{ij}^p = \frac{\frac{\partial f}{\partial \sigma_{kl}} \dot{\sigma}_{ij} + \frac{\partial f}{\partial \chi} \dot{\chi}}{\frac{\partial f}{\partial \epsilon_{kl}^p} \frac{\partial f}{\partial \sigma_{kl}}} \frac{\partial f}{\partial \sigma_{ij}} \text{ if } f = 0 \quad (2-117)$$

$$\dot{\epsilon}_{ij}^p = 0 \text{ if } f < 0$$

SECTION 3

THE CAP MODEL AND DEVELOPMENT OF AN EXPLICIT FORM OF ELASTO-PLASTIC CONSTITUTIVE MATRIX OF MATERIAL

3.1 The Cap Model

The cap model (Ref. 1) is a classical incremental plasticity model defined by a yield surface and a plastic strain rate vector (Figure 3-1). The model exhibits three different modes of behavior: elastic, failure, and cap.

Elastic Mode: The elastic mode of behavior occurs when the stress point is within the failure envelope, and stress changes result in recoverable deformations.

Failure Mode: During the failure mode of behavior, the stress point lies on the failure envelope with a stress-strain relation given by equation (3-1). As shown in Figure 3-1, the associated flow rule requires that the plastic strain rate vector be directed upward and to the left. Therefore, the plastic strain during failure is composed of a deviatoric or shear component together with a volumetric, or dilatant component.

Cap Mode: The cap mode of behavior occurs when the stress point lies on the movable cap and pushes it outward. The stress-strain relation is given by equation (3-2). As shown in Figure 3-1, the associated flow rule requires that during cap action the plastic strain rate vector be directed upward and to the right. This implies that the plastic strain rate produces an irreversible decrease in volume in conjunction with the irreversible shearstrain. This reduction in volume is referred to as compaction.

This compaction leads to an increase in the cap parameter ϵ_V^P which, in turn, through equation (3-5), leads to an increase in $X(\kappa)$ and hence the cap moves to the right. Either J_1 or $\sqrt{J_2}$ or both must increase to maintain the cap mode of behavior.

In soils, the dilatancy associated with failure leads to a decrease in $\bar{\epsilon}_V^P$, resulting in a leftward movement of the cap. This cap movement is limited if and when the cap reaches the stress point

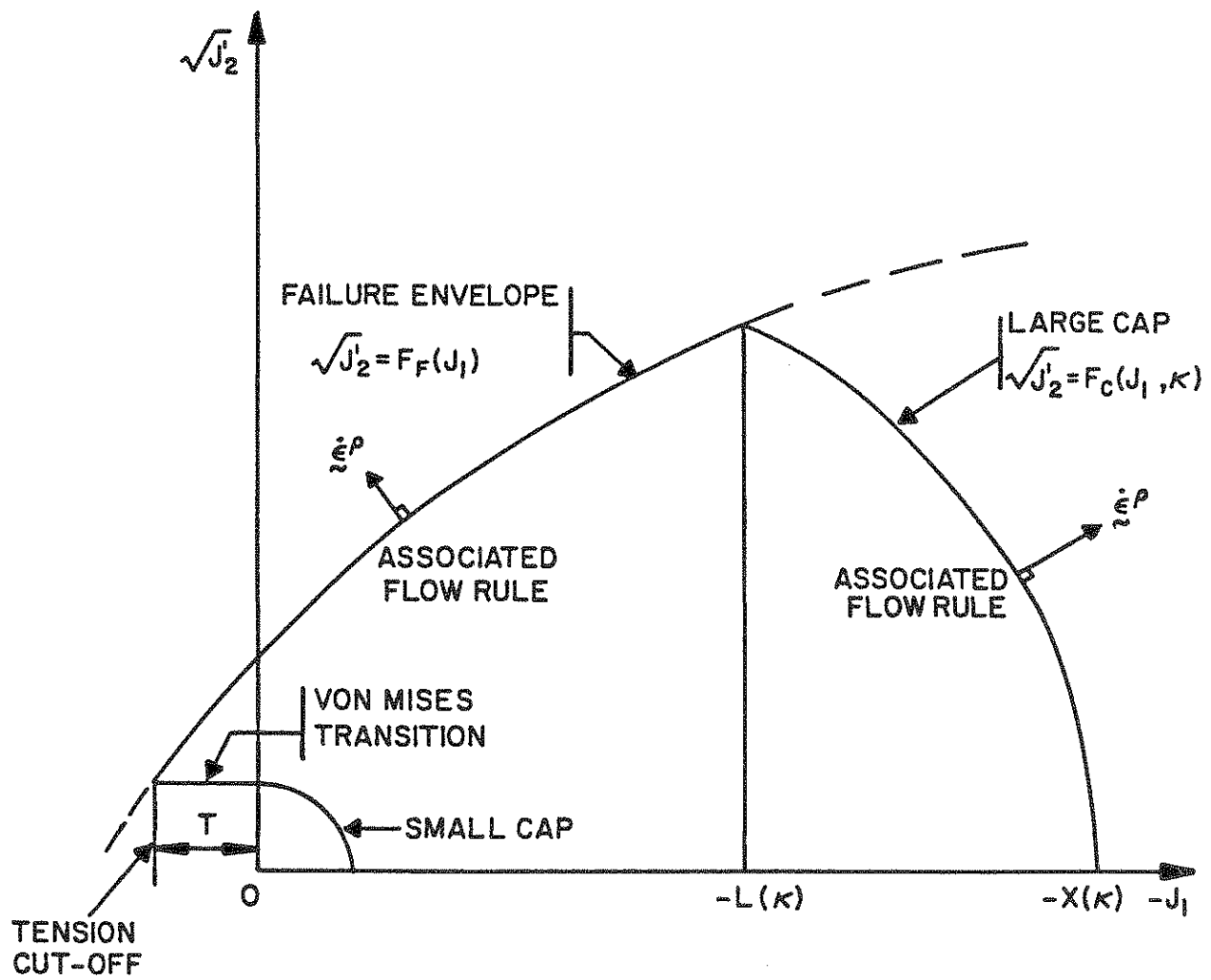


FIGURE 3-1 Typical Yield Surface In Cap Model

(so that the stress point is at a corner of the yield surface). When this occurs, the associated flow rule requires that the plastic strain rate vector lie between the outward drawn normals at the corner. At this point, the subsequent plastic straining is purely in shear.

Constitutive Equations (see Figure 3-1)

The following constitutive relations were developed for the McCormick Ranch Sand (Ref. 1).

Failure Mode:

$$F_F(J_1) = A-C \exp(BJ_1) \quad J_1 < L \quad (3-1)$$

Cap Mode:

$$F_c(J_1, \kappa) = \frac{1}{R} \sqrt{\{[X(\kappa) - L(\kappa)]^2 - [J_1 - L(\kappa)]^2\}} \quad L < J_1 < X \quad (3-2)$$

where:

$$L(\kappa) = \begin{cases} \kappa & \text{if } \kappa < 0 \\ 0 & \text{if } \kappa \geq 0 \end{cases} \quad (3-3)$$

$$X(\kappa) = \kappa - RF_F(\kappa) \quad (3-4)$$

and A,B,C, and R are material parameters and κ is a hardening parameter.

$$\bar{\epsilon}_V^p = W \{ \exp [DX(\kappa)] - 1 \} \quad (3-5)$$

$$\dot{\bar{\epsilon}}_V^p = \begin{cases} \dot{\epsilon}_V^p & \text{if } \dot{\epsilon}_V^p \leq 0 \text{ or } \kappa < J_1 \text{ and } \kappa < 0 \\ 0 & \text{otherwise} \end{cases} \quad (3-6)$$

Tension cutoff is assumed to be represented by

$$J_1 = T \quad (3-7)$$

where T denotes the maximum allowable hydrostatic tension.

3.1.1 Derivatives of Cap Functions

Consider the yield function in terms of stress invariants and a hardening parameter given by

$$F(J_1, \sqrt{J_2'}, \kappa) = 0 \quad (3-8)$$

where:

$$J_1 = \sigma_{ij} \delta_{ij} = \sigma_{kk} \quad (3-9)$$

in which δ_{ij} is the Kronecker delta and

$$J_2' = \frac{1}{2} s_{ij} s_{ij} = \frac{1}{2} \left(\sigma_{ij} - \frac{1}{3} \sigma_{kk} \delta_{ij} \right) \left(\sigma_{ij} - \frac{1}{3} \sigma_{kk} \delta_{ij} \right) \quad (3-10)$$

and κ is the hardening parameter. Now differentiating equation (3-8) we get

$$\frac{\partial F}{\partial J_1} \partial J_1 + \frac{1}{2\sqrt{J_2'}} \frac{\partial F}{\partial \sqrt{J_2'}} \partial J_2' + \frac{\partial F}{\partial \kappa} \partial \kappa \quad (3-11)$$

Divide equation (3-11) by σ_{ij} and get

$$\frac{\partial F}{\partial \sigma_{ij}} = \frac{\partial F}{\partial J_1} \frac{\partial J_1}{\partial \sigma_{ij}} + \frac{1}{2\sqrt{J_2'}} \frac{\partial F}{\partial \sqrt{J_2'}} \frac{\partial J_2'}{\partial \sigma_{ij}} \quad (3-12)$$

or

$$\frac{\partial F}{\partial \sigma_{ij}} = \frac{\partial F}{\partial J_1} \delta_{ij} + \frac{s_{ij}}{2\sqrt{J_2}} \frac{\partial F}{\partial \sqrt{J_2}} \quad (3-13)$$

Failure Mode:

From equation (3-1)

$$F_F = \sqrt{J_2} - A + C \exp(BJ_1) \quad (3-14)$$

Hence

$$\frac{\partial F_F}{\partial J_1} = CB \exp(BJ_1) \quad (3-15)$$

$$\frac{\partial F_F}{\partial \sqrt{J_2}} = 1 \quad (3-16)$$

Hence from equation (3-13)

$$\frac{\partial F_F}{\partial \sigma_{ij}} = CB \exp(BJ_1) \delta_{ij} + \frac{s_{ij}}{2\sqrt{J_2}} \quad (3-17)$$

Cap Mode:

From equation (3-2)

$$F_c = \sqrt{J_2} - \frac{1}{R} \sqrt{\{[X(\kappa) - L(\kappa)]^2 - [J_1 - L(\kappa)]^2\}} \quad (3-18)$$

Differentiating equation (3-18) we get

$$-R^2 \sqrt{J_2'} \partial F_c + 2R^2 \partial J_2' = (X-L)\partial X + (J_1 - X) \partial L + (L - J_1) \partial J_1 \quad (3-19)$$

Divide equation (3-19) by $\partial \sigma_{ij}$ and rearrange terms

$$\frac{\partial F_c}{\partial \sigma_{ij}} = \frac{2s_{ij}}{\sqrt{J_2'}} - \frac{(L - J_1)}{R^2 \sqrt{J_2'}} \delta_{ij} \quad (3-20)$$

$$\frac{\partial F_c}{\partial X} = \frac{(L - X)}{R^2 \sqrt{J_2'}} \quad (3-21)$$

$$\frac{\partial F}{\partial L} = \frac{(X - J_1)}{R^2 \sqrt{J_2'}} \quad (3-22)$$

From equation (3-4)

$$\frac{\partial X}{\partial \kappa} = 1 + RCB \exp(B\kappa) \quad (3-23)$$

From equation (3-3)

$$\frac{\partial L}{\partial \kappa} = \begin{cases} 1 & \text{if } \kappa < 0 \\ 0 & \text{if } \kappa \geq 0 \end{cases} \quad (3-24)$$

Now

$$\frac{\partial F_c}{\partial \kappa} = \frac{\partial F_c}{\partial X} \cdot \frac{\partial X}{\partial \kappa} + \frac{\partial F_c}{\partial L} \cdot \frac{\partial L}{\partial \kappa} \quad (3-25)$$

$$= \frac{(L - X)}{R^2 \sqrt{J_2}} [1 + RCB \exp (B\kappa)] + \frac{(X - J_1)}{R^2 \sqrt{J_2}} \quad (3-26)$$

if $\kappa < 0$

$$\frac{\partial F_c}{\partial \kappa} = \frac{(L - x)}{R^2 \sqrt{J_2}} [1 + RCB \exp (B\kappa)] \quad \text{if } \kappa \geq 0 \quad (3-27)$$

Elasto-Plastic Constitutive Matrix

Consider the yield surface defined by

$$F(\sigma, \kappa) = 0 \quad (3-28)$$

where:

κ = Hardening parameter

σ = General state of stress

From incremental theory of plasticity

$$d\xi = d\xi^e + d\xi^p \quad (3-29)$$

where:

$d\xi$ = Increment of total strain

$d\xi^e$ = Increment of elastic strain

$d\xi^p$ = Increment of plastic strain

Now

$$d\xi^e = [D]^{-1} d\sigma \quad (3-30)$$

$$\begin{aligned} d\varepsilon^p &= \lambda \frac{\partial F}{\partial \sigma} \quad \text{if } F = 0 \\ &= 0 \quad \text{if } F < 0 \end{aligned} \quad (3-31)$$

This is known as the normality principle and since F is assumed to coincide with the plastic potential, this phenomenon is also referred to as associated flow rule. Equation (3-29) becomes

$$d\varepsilon = [D]^{-1} d\sigma + \lambda \frac{\partial F}{\partial \sigma} \quad (3-32)$$

Now differentiating equation (3-28) we get

$$\partial F = \frac{\partial F}{\partial \sigma} d\sigma + \frac{\partial F}{\partial \kappa} d\kappa = 0 \quad (3-33)$$

or

$$\left\{ \frac{\partial F}{\partial \sigma} \right\}^T d\sigma - A\lambda = 0 \quad (3-34)$$

where:

$$A = - \frac{1}{\lambda} \frac{\partial F}{\partial \kappa} d\kappa \quad (3-35)$$

and

$$d\sigma = [d\sigma_1 \ d\sigma_2 \ d\sigma_3 \ d\sigma_4 \ d\sigma_5 \ d\sigma_6]^T \quad (3-36)$$

is the vector form of the increment of the stress tensor. Hence

$$d\sigma_1 = d\sigma_{11}$$

$$d\sigma_2 = d\sigma_{22}$$

$$d\sigma_3 = d\sigma_{33}$$

$$d\sigma_4 = d\sigma_{23} = d\sigma_{32}$$

$$d\sigma_5 = d\sigma_{13} = d\sigma_{31}$$

$$d\sigma_6 = d\sigma_{12} = d\sigma_{21} \quad (3-37)$$

Multiply equation (3-32) by $\left\{ \frac{\partial F}{\partial \sigma} \right\}^T [D]$. We get

$$\left\{ \frac{\partial F}{\partial \sigma} \right\}^T [D] d\xi = \left\{ \frac{\partial F}{\partial \sigma} \right\}^T d\sigma + \left\{ \frac{\partial F}{\partial \sigma} \right\}^T [D] \left\{ \frac{\partial F}{\partial \sigma} \right\} \lambda \quad (3-38)$$

From equation (3-34)

$$\left\{ \frac{\partial F}{\partial \sigma} \right\}^T d\sigma = A\lambda \quad (3-39)$$

Substituting equation (3-39) into equation (3-38) we get

$$\left\{ \frac{\partial F}{\partial \sigma} \right\}^T [D] d\xi = \left[A + \left\{ \frac{\partial F}{\partial \sigma} \right\}^T [D] \left\{ \frac{\partial F}{\partial \sigma} \right\} \right] \lambda \quad (3-40)$$

$$\lambda = d\xi \cdot \left\{ \frac{\partial F}{\partial \sigma} \right\}^T [D] \cdot \left[A + \left\{ \frac{\partial F}{\partial \sigma} \right\}^T [D] \left\{ \frac{\partial F}{\partial \sigma} \right\} \right]^{-1} \quad (3-41)$$

From equation (3-32)

$$d\sigma = [D] d\varepsilon - \lambda [D] \frac{\partial F}{\partial \sigma} \quad (3-42)$$

Substituting value of λ into equation (3-42) we get

$$d\sigma = \left([D] - [D] \left\{ \frac{\partial F}{\partial \sigma} \right\} \left\{ \frac{\partial F}{\partial \sigma} \right\}^T [D] \cdot \left(A + \left\{ \frac{\partial F}{\partial \sigma} \right\}^T [D] \left\{ \frac{\partial F}{\partial \sigma} \right\} \right)^{-1} \right) d\varepsilon \quad (3-43)$$

The elasto-plastic stress strain relation can be expressed in incremental form as

$$d\sigma = [D^{ep}] d\varepsilon \quad (3-44)$$

where $[D^{ep}]$, the elasto-plastic constitutive matrix, is given by

$$[D^{ep}] = [D] - [D] \left\{ \frac{\partial F}{\partial \sigma} \right\} \left\{ \frac{\partial F}{\partial \sigma} \right\}^T [D] \cdot \left(A + \left\{ \frac{\partial F}{\partial \sigma} \right\}^T [D] \left\{ \frac{\partial F}{\partial \sigma} \right\} \right)^{-1} \quad (3-45)$$

and $[D]$ is the elastic constitutive matrix. The work hardening parameter κ is taken as the amount of plastic work done during plastic deformation. Thus

$$d\kappa = \sigma_1 d\varepsilon_1^p + \sigma_2 d\varepsilon_2^p + \dots + \sigma_6 d\varepsilon_6^p = \sigma^T d\varepsilon^p \quad (3-46)$$

Substituting equation (3-46) into equation (3-31) we get

$$d\kappa = \lambda \sigma^T \frac{\partial F}{\partial \sigma} \quad (3-47)$$

Eliminate λ by substituting equation (3-47) into equation (3-35) and get

$$A = - \frac{\partial F}{\partial \kappa} \mathcal{G}^T \frac{\partial F}{\partial \mathcal{G}} \quad (3-48)$$

SECTION 4 THE PROPOSED Q-MODEL

Wave energy dissipation in a medium results not only from dispersion of wave energy from the source but also from damping by energy losses within the medium. In this report, attention is focused on the latter.

Material damping has been modeled in diverse ways in recent years. The Rayleigh damping of the form

$$C = \alpha M + \beta K \tag{4-1}$$

has been used with some success in time stepping schemes. Here C is the damping matrix, α , β are constants, M and K are mass and stiffness matrices, respectively.

The proposed Q-model emphasizes the viscous and energy dissipation characteristics of the material. This model is of great advantage because all the parameters involved are readily determined from tests, and the resulting damping expression is very easy to apply in time stepping schemes, and thus lends itself readily for use in computer codes. The Q-model is used in conjunction with elastic or elastoplastic constitutive matrix to produce a viscoelastic or viscoplastic material model, respectively.

4.1 Theoretical Development

The specific dissipation factor, sometimes referred to as the coefficient of internal friction is defined as

$$Q^{-1} = \frac{\Delta W}{2\pi W} \tag{4-2}$$

where W is the elastic strain energy stored per unit cycle per unit volume, and ΔW is the energy dissipated per unit cycle per unit volume.

In terms of phase angle

$$Q^{-1} = \tan \phi \quad (4-3)$$

where ϕ is the phase angle between stress and strain.

In terms of logarithmic decrement

$$Q^{-1} = \frac{\delta}{\pi} \quad (4-4)$$

where δ is the logarithmic decrement.

Now consider the complex spring-mass system as shown in Figure 4-1. The equation of motion of the system can be expressed as

$$F(t) - k^* x = m\ddot{x} \quad (4-5)$$

where k^* is the complex spring constant given by

$$k^* = k(1 + i \tan \phi) \quad (4-6)$$

where ϕ is the coefficient of internal friction, and m is the mass of the system. From equation (4-5)

$$m\ddot{x} + k(1 + i \tan \phi) x = F(t) \quad (4-7)$$

By Fourier Transformation of equation (4-7) we get

$$-\omega^2 mx + k(1 + i \tan \phi) x = \bar{F}(\omega) \quad (4-8)$$

Substitute equation (4-3) into equation (4-8) and get

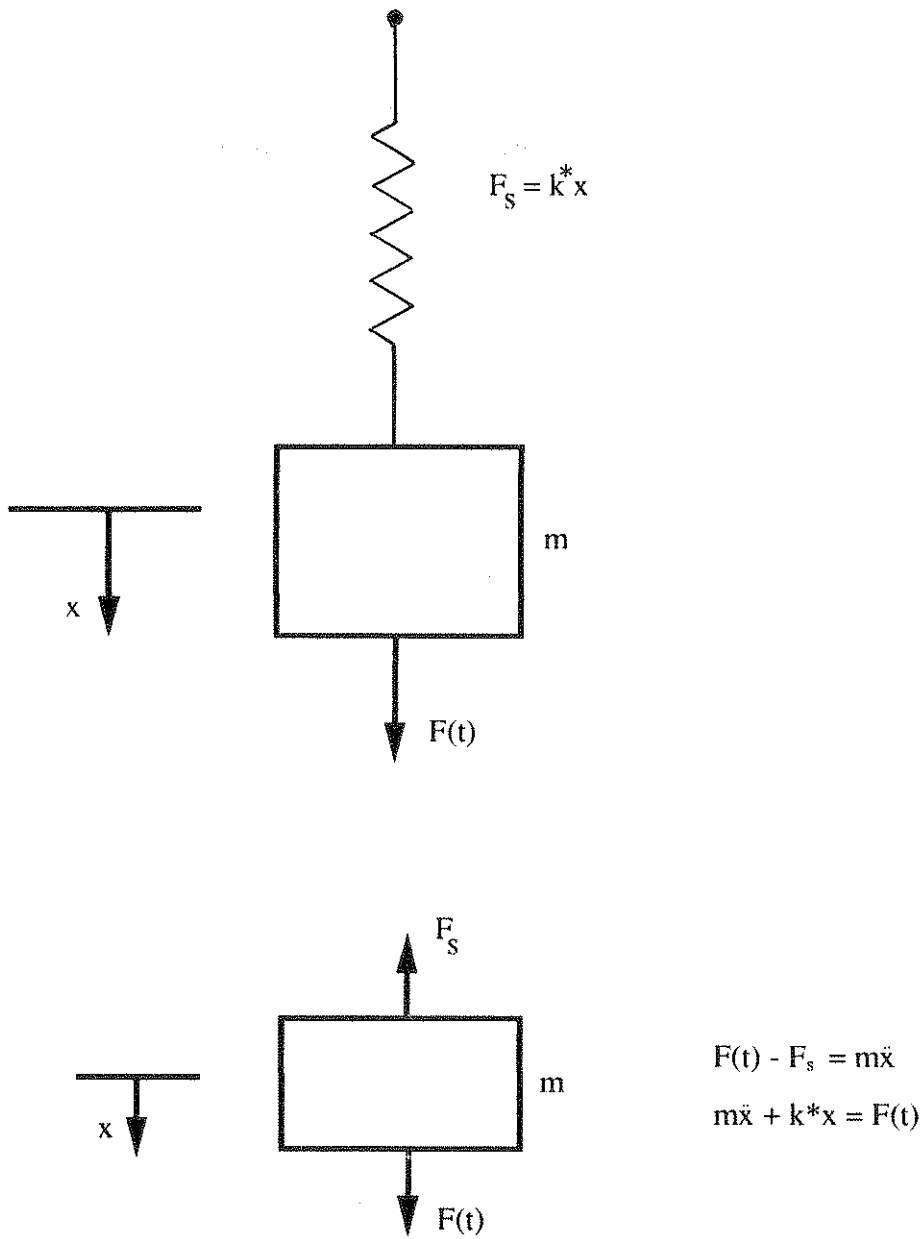


FIGURE 4-1 Mechanical Model For Single Degree Of Freedom Anelastic Spring Mass System

$$-\omega^2 mx + k(1 + iQ^{-1}) x = \bar{F}(\omega) \quad (4-9)$$

Alternatively, for a single degree of freedom system, the equation of motion of the system can be expressed as

$$m\ddot{x} + c(t)\dot{x} + kx = F(t) \quad (4-10)$$

where:

m = mass constant

c = damping constant

k = spring constant

Taking Fourier Transform of equation (4-10) we get

$$-\omega^2 mx + (k + i\omega\bar{c}(\omega)) x = \bar{F}(\omega) \quad (4-11)$$

or

$$-\omega^2 mx + k \left(1 + \frac{i\omega\bar{c}(\omega)}{k} \right) x = \bar{F}(\omega) \quad (4-12)$$

Comparing the complex terms of equations (4-9) and (4-12) we get

$$Q^{-1} = \frac{\omega\bar{c}(\omega)}{k} \quad (4-13)$$

or

$$\bar{c}(\omega) = \frac{kQ^{-1}}{\omega} \quad (4-14)$$

In general, $\bar{c}(\omega)$ could be expanded in the form of a Laurent Series with "even powers of ω ." For

simplicity we can write

$$\bar{c}(\omega) = a_0 + \frac{a_1}{\omega^2} + \frac{a_2}{\omega^4} + \dots \quad (4-15)$$

in order to avoid the involvement of higher derivatives of $x(t)$ associated with terms like b 's in the general form:

$$\bar{c}(\omega) = (b_1 \omega^2 + b_2 \omega^4 + \dots) + a_0 + \frac{a_1}{\omega^2} + \frac{a_2}{\omega^4} + \dots \quad (4-16)$$

Consider a constant Q -model (i.e., frequency independent over a prescribed frequency range).

Let

$$Q(\omega) \approx Q_0 \text{ between } \Omega_1 < \omega < \Omega_2, \omega > 0 \quad (4-17)$$

then from equation (4-14)

$$\bar{c}(\omega) = \frac{kQ_0^{-1}}{\omega} \approx a_0 + \frac{a_1}{\omega^2} + \frac{a_2}{\omega^4} \quad (4-18)$$

or

$$kQ_0^{-1} \approx a_0 \omega + \frac{a_1}{\omega} + \frac{a_2}{\omega^3} \quad \omega \geq 0 \quad (4-19)$$

The mean square error $e(\Omega_1, \Omega_2)$ between Ω_1 and Ω_2 is given by

$$e = \int_{\Omega_1}^{\Omega_2} \left(kQ_0^{-1} - a_0 \omega - \frac{a_1}{\omega} - \frac{a_2}{\omega^3} \right)^2 d\omega \quad (4-20)$$

Minimization of $e(\Omega_1, \Omega_2)$ to obtain the constant a 's gives

$$\frac{\partial e}{\partial a_0} = -2 \int_{\Omega_1}^{\Omega_2} \left(kQ_0^{-1} - a_0 \omega - \frac{a_1}{\omega} - \frac{a_2}{\omega^3} \right) \omega d\omega = 0 \quad (4-21)$$

or

$$\frac{kQ_0^{-1}}{2} (\Omega_2^2 - \Omega_1^2) - \frac{a_0}{3} (\Omega_2^3 - \Omega_1^3) - a_1 (\Omega_2 - \Omega_1) + a_2 \left(\frac{1}{\Omega_2} - \frac{1}{\Omega_1} \right) = 0 \quad (4-22)$$

$$\frac{\partial e}{\partial a_1} = -2 \int_{\Omega_1}^{\Omega_2} \left(kQ_0^{-1} - a_0 \omega - \frac{a_1}{\omega} - \frac{a_2}{\omega^3} \right) \frac{1}{\omega} d\omega = 0 \quad (4-23)$$

or

$$\begin{aligned} & -kQ_0^{-1} \ln \left(\frac{\Omega_2}{\Omega_1} \right) - a_0 (\Omega_2 - \Omega_1) \\ & + a_1 \left(\frac{1}{\Omega_2} - \frac{1}{\Omega_1} \right) + \frac{a_2}{3} \left(\frac{1}{\Omega_2^3} - \frac{1}{\Omega_1^3} \right) = 0 \end{aligned} \quad (4-24)$$

$$\frac{\partial e}{\partial a_2} = -2 \int_{\Omega_1}^{\Omega_2} \left(kQ_0^{-1} - a_0 \omega - \frac{a_1}{\omega} - \frac{a_2}{\omega^3} \right) \frac{1}{\omega^3} d\omega = 0 \quad (4-25)$$

or

$$\begin{aligned} & -\frac{kQ_0^{-1}}{2} \left(\frac{1}{\Omega_2^2} - \frac{1}{\Omega_1^2} \right) + a_0 \left(\frac{1}{\Omega_2} - \frac{1}{\Omega_1} \right) \\ & + \frac{a_1}{3} \left(\frac{1}{\Omega_2^3} - \frac{1}{\Omega_1^3} \right) + \frac{a_2}{5} \left(\frac{1}{\Omega_2^5} - \frac{1}{\Omega_1^5} \right) = 0 \end{aligned} \quad (4-26)$$

Equations (4-22), (4-24) and (4-26) can be expressed in matrix form as

$$\begin{bmatrix} \frac{1}{3} (\Omega_2^3 - \Omega_1^3) & (\Omega_2 - \Omega_1) & -\left(\frac{1}{\Omega_2} - \frac{1}{\Omega_1}\right) \\ (\Omega_2 - \Omega_1) & -\left(\frac{1}{\Omega_2} - \frac{1}{\Omega_1}\right) & -\frac{1}{3} \left(\frac{1}{\Omega_2^3} - \frac{1}{\Omega_1^3}\right) \\ -\left(\frac{1}{\Omega_2} - \frac{1}{\Omega_1}\right) & -\frac{1}{3} \left(\frac{1}{\Omega_2^3} - \frac{1}{\Omega_1^3}\right) & -\frac{1}{5} \left(\frac{1}{\Omega_2^5} - \frac{1}{\Omega_1^5}\right) \end{bmatrix} \begin{bmatrix} a_0 \\ a_1 \\ a_2 \end{bmatrix} = \begin{bmatrix} \frac{kQ_o^{-1}}{2} (\Omega_2^2 - \Omega_1^2) \\ -kQ_o^{-1} \ln \left(\frac{\Omega_2}{\Omega_1}\right) \\ -\frac{kQ_o^{-1}}{2} \left(\frac{1}{\Omega_2^2} - \frac{1}{\Omega_1^2}\right) \end{bmatrix}$$

The above system of equations can be solved for the values of a_0 , a_1 and a_2 . Multiply both sides of equation (4-15) by $i\omega\bar{x}$, and get

$$\bar{c}(\omega) i\omega\bar{x} = a_0 i\omega\bar{x} + \frac{a_1}{\omega^2} i\omega\bar{x} + \frac{a_2 i\omega\bar{x}}{\omega^4} \quad (4-27)$$

or

$$\bar{c}(\omega) i\omega\bar{x} = a_0 i\omega\bar{x} - \frac{a_1 \bar{x}}{i\omega} + \frac{(i)^4 a_2 \bar{x}}{(i\omega)^3} \quad (4-28)$$

Define the operator

$$D = \frac{d}{dt} = i\omega \quad (4-29)$$

then the inverse operator

$$D^{-1} = \int_0^t d\tau = \frac{1}{i\omega} \quad (4-30)$$

From the relation in equations (4-29) and (4-30), equation (4-28) can be rewritten in the time domain as

$$c(t) D^2 x(t) = a_0 D^2 x(t) - a_1 D^{-1} x(t) + a_2 D^{-3} x(t) \quad (4-31)$$

or

$$c(t) \ddot{x} = a_0 \ddot{x} - a_1 \int_0^t x(\tau) d\tau + a_2 \int_0^t \int_0^\tau \int_0^\tau x(\tau) d\tau d\tau d\tau \quad (4-32)$$

Substituting equation (4-32) into equation (4-10), the equation of motion of the system becomes

$$F(t) = m\ddot{x} + a_0 \dot{x} + kx - a_1 \int_0^t x(\tau) d\tau + a_2 \int_0^t \int_0^\tau \int_0^\tau x(\tau) d\tau d\tau d\tau \quad (4-33)$$

4.2 Numerical Procedure

The equation of motion of the system as presented in equation (4-33) can be solved numerically. Two integration schemes -- the explicit time integration scheme using the Central Difference Method and the implicit time integration scheme using the Newmark Method -- are considered here.

$$\text{Let } z_1 = \int_0^t \int_0^\tau \int_0^\tau x(\tau) d\tau d\tau d\tau \quad (4-34)$$

$$z_3 = \int_0^t x(\tau) d\tau \quad (4-35)$$

$$z_4 = x \quad (4-36)$$

$$z_5 = \dot{x} \quad (4-37)$$

$$z_6 = \ddot{x} \quad (4-38)$$

Then

$$\frac{d}{dt} \begin{Bmatrix} z_1 \\ z_2 \\ z_3 \\ z_4 \\ z_5 \end{Bmatrix} = \begin{Bmatrix} z_2 \\ z_3 \\ z_4 \\ z_5 \\ z_6 \end{Bmatrix} \quad (4-39)$$

and equation (4-33) becomes

$$F(t) = mz_6 + a_0 z_5 + kz_4 - a_1 z_3 + a_2 z_1 \quad (4-40)$$

from which

$$z_6 = \frac{1}{m} [F(t) - a_0 z_5 - kz_4 + a_1 z_3 - a_2 z_1] \quad (4-41)$$

4.2.1 Explicit Time Integration Scheme Using Central Difference Method

From Central Difference Method the following expressions are obtained

$${}^t \dot{z}_1 = \frac{1}{2\Delta t} ({}^{t+\Delta t} z_1 - {}^{t-\Delta t} z_1) = {}^t z_2 \quad (4-42)$$

$${}^t \dot{z}_2 = \frac{1}{2\Delta t} ({}^{t+\Delta t} z_2 - {}^{t-\Delta t} z_2) = {}^t z_3 \quad (4-43)$$

$${}^t \dot{z}_3 = \frac{1}{2\Delta t} ({}^{t+\Delta t} z_3 - {}^{t-\Delta t} z_3) = {}^t z_4 \quad (4-44)$$

$${}^t \dot{z}_4 = \frac{1}{2\Delta t} ({}^{t+\Delta t} z_4 - {}^{t-\Delta t} z_4) = {}^t z_5 \quad (4-45)$$

$${}^t \dot{z}_5 = \frac{1}{2\Delta t} ({}^{t+\Delta t} z_5 - {}^{t-\Delta t} z_5) = {}^t z_6 \quad (4-46)$$

From Taylor's expansion of z_4 we get

$${}^{t+\Delta t} z_4 = {}^t z_4 + \Delta t {}^t \dot{z}_4 + \frac{\Delta t^2}{2} {}^t \ddot{z}_4 + \dots \quad (4-47)$$

$${}^t \ddot{z}_4 = \frac{2}{\Delta t^2} ({}^{t+\Delta t} z_4 - {}^t z_4 - \Delta t {}^t \dot{z}_4) \quad (4-48)$$

Substituting value of ${}^t \dot{z}_4$ from equation (4-45) into equation (4-48) and rearranging terms, we get

$${}^t \ddot{z}_4 = \frac{1}{\Delta t^2} ({}^{t+\Delta t} z_4 - 2{}^t z_4 + {}^{t-\Delta t} z_4) = {}^t z_6 \quad (4-49)$$

Similar derivations result in the following expressions

$${}^t \ddot{z}_1 = \frac{1}{\Delta t^2} ({}^{t+\Delta t} z_1 - 2{}^t z_1 + {}^{t-\Delta t} z_1) = {}^t z_3 \quad (4-50)$$

$${}^t \ddot{z}_2 = \frac{1}{\Delta t^2} ({}^{t+\Delta t} z_2 - 2{}^t z_2 + {}^{t-\Delta t} z_2) = {}^t z_4 \quad (4-51)$$

$${}^t \ddot{z}_3 = \frac{1}{\Delta t^2} ({}^{t+\Delta t} z_3 - 2{}^t z_3 + {}^{t-\Delta t} z_3) = {}^t z_5 \quad (4-52)$$

The equilibrium equation for a multi-degree of freedom system is given at time t by

$$[M]\ddot{\underline{X}}_t + [C]\dot{\underline{X}}_t + [K]\underline{X}_t = \underline{F}(t) \quad (4-53)$$

Substituting values of $\ddot{\underline{X}}_t$, $\dot{\underline{X}}_t$, \underline{X}_t and $[C]$ into equation (4-53), we get

$$[M]^t z_6 + [A_0]^t z_5 - [A_1]^t z_3 + [A_2]^t z_1 + [K]^t z_4 = \underline{F}(t) \quad (4-54)$$

where the sign (.) under a letter denotes 'vector.'

Further substitution and rearrangement of terms give the explicit integration scheme

$$\begin{aligned} \left(\frac{1}{\Delta t^2} [M] + \frac{1}{2\Delta t} [A_0] \right)^{t+\Delta t} z_4 = F(t) - \left([K] - \frac{2}{\Delta t^2} [M] \right)^t z_4 \\ - \left(\frac{1}{\Delta t^2} [M] - \frac{1}{2\Delta t} [A_0] \right)^{t-\Delta t} z_4 + [A_1]^t z_3 - [A_2]^t z_1 \end{aligned} \quad (4-55)$$

where:

$$[A_0] = \alpha[K]$$

$$[A_1] = \beta[K] \quad (4-56)$$

$$[A_2] = \gamma[K]$$

Taylor's expansion of z_3 , z_2 , z_1 results in the following expressions

$${}^{t+\Delta t} z_3 = {}^t z_3 + \Delta t {}^t z_4 + \frac{\Delta t^2}{2} {}^t z_5 + \dots \quad (4-57)$$

$${}^{t+\Delta t} z_2 = {}^t z_2 + \Delta t {}^t z_3 + \frac{\Delta t^2}{2} {}^t z_4 + \dots \quad (4-58)$$

$${}^{t+\Delta t} z_1 = {}^t z_1 + \Delta t {}^t z_2 + \frac{\Delta t^2}{2} {}^t z_3 + \dots \quad (4-59)$$

4.3. Implicit Time Integration Scheme Using Newmark Method

From Newmark's Method we assume

$${}^{t+\Delta t}z_5 = {}^t z_5 + \Delta t \left[(1 - \delta) {}^t z_6 + \delta {}^{t+\Delta t} z_6 \right] \quad (4-60)$$

$${}^{t+\Delta t}z_4 = {}^t z_4 + \Delta t {}^t z_5 + \Delta t^2 \left[\left(\frac{1}{2} - \lambda \right) {}^t z_6 + \lambda {}^{t+\Delta t} z_6 \right] \quad (4-61)$$

where δ and λ are integration parameters, and z_4 , z_5 and z_6 are the displacement, velocity, and acceleration vectors, respectively. The dynamic equation of motion at time $t + \Delta t$ is given by

$$[M]\ddot{X}_{t+\Delta t} + [C]\dot{X}_{t+\Delta t} + [K]X_{t+\Delta t} = F(t + \Delta t) \quad (4-62)$$

or

$$[M] {}^{t+\Delta t}z_6 + [A_0] {}^{t+\Delta t}z_5 - [A_1] {}^{t+\Delta t}z_3 + [A_2] {}^{t+\Delta t}z_1 + [K] {}^{t+\Delta t}z_4 = F(t + \Delta t) \quad (4-63)$$

From equation (4-61)

$${}^{t+\Delta t}z_6 = \frac{1}{\lambda \Delta t^2} \left({}^{t+\Delta t}z_4 - {}^t z_4 - \Delta t {}^t z_5 \right) - \left(\frac{1}{2\lambda} - 1 \right) {}^t z_6 \quad (4-64)$$

substitute values ${}^{t+\Delta t}z_6$ from equation (4-64) into equation (4-60) and get

$${}^{t+\Delta t}z_5 = {}^t z_5 + \Delta t (1 - \delta) {}^t z_6 + \delta \Delta t {}^{t+\Delta t} z_6 \quad (4-65)$$

Substitute values of ${}^{t+\Delta t}z_5$ and ${}^{t+\Delta t}z_6$ from equations (4-65) and (4-64) into equation (4-63) and

rearrange terms and get

$$\begin{aligned}
 & \left(\frac{1}{\lambda \Delta t^2} [M] + \frac{\delta}{\lambda \Delta t} [A_o] + [K] \right)^{t+\Delta t} z_4 \\
 &= F(t + \Delta t) + [M] \left(\frac{1}{\lambda \Delta t^2} {}^t z_4 + \frac{1}{\lambda \Delta t} {}^t z_5 + \left(\frac{1}{2\lambda} - 1 \right) {}^t z_6 \right) \\
 &+ [A_o] \left(\frac{\delta}{\lambda \Delta t} {}^t z_4 + \left(\frac{\delta}{\lambda} - 1 \right) {}^t z_5 + \frac{\Delta t}{2} \left(\frac{\delta}{\lambda} - 2 \right) {}^t z_6 \right) \\
 &+ [A_1] {}^{t+\Delta t} z_3 - [A_2] {}^{t+\Delta t} z_1
 \end{aligned} \tag{4-66}$$

where again

$$[A_o] = \alpha[K]$$

$$[A_1] = \beta[K]$$

$$[A_2] = \gamma[K]$$

α, β, γ are constants from Q-model.

4.4 The Numerical Algorithm

1. Determine stress-strain relation of material using any appropriate elastoplastic material model.
2. Obtain elasto-plastic constitutive matrix $[D^{ep}]$.
3. Obtain elasto-plastic stiffness matrix using the relation

$$[K] = \int_v [B]^T [D^{ep}] [B] dv$$

where [B] is the strain displacement matrix.

4. Solve the dynamic equation of motion

$$[M]\ddot{X} + [C]\dot{X} + [K]X = F(t)$$

where:

$$[C]\dot{X} = [A_0]\dot{X}(t) - [A_1] \int_0^t X(\tau) d\tau + [A_2] \int_0^t \int_0^\tau \int_0^\tau X(\tau) d\tau d\tau d\tau$$

$$[A_0] = \alpha[K]$$

$$[A_1] = \beta[K]$$

$$[A_2] = \gamma[K]$$

α, β, γ are constants from Q-model.

4.4.1 Explicit Time Integration Scheme

For explicit Time Integration Scheme using the Central Difference Method, the dynamic equation of motion at time t is solved as follows:

1. Select time step Δt , $\Delta t < \Delta t_{cr}$
2. Compute integration constants

$$\tau_0 = \frac{1}{\Delta t^2}$$

$$\tau_1 = \frac{1}{2\Delta t}$$

$$\tau_2 = 2\tau_0$$

$$\tau_3 = \frac{1}{\tau_2}$$

3. Initialize $z_1, z_2, z_3, z_4, z_5, z_6$ and calculate

$${}^{-\Delta t}z_4 = {}^0z_4 - \Delta t {}^0\dot{z}_4 + \tau_3 {}^0\ddot{z}_4$$

4. Obtain effective mass matrix

$$[\hat{M}] = \tau_0 [M] + \tau_1 [A_0]$$

5. Triangularize $[\hat{M}]$: $\hat{M} = [L] [D] [L]^T$

6. For each time step:

a. Calculate effective force at time t:

$$\begin{aligned} \hat{F}(t) = & F(t) - ([K] - \tau_2 [M]) {}^t z_4 - (\tau_0 [M] - \tau_1 [A_0]) {}^{t-\Delta t} z_4 \\ & + [A_1] {}^t z_3 - [A_2] {}^t z_3 \end{aligned}$$

b. Solve for displacements at time $t + \Delta t$

$$[L] [D] [L]^T {}^{t+\Delta t} z_4 = \hat{F}(t)$$

c. Solve for z_1, z_2, z_3 at time $t + \Delta t$ as follows

$${}^{t+\Delta t}z_1 = {}^t z_1 + \Delta t {}^t z_2 + \frac{1}{\tau_2} {}^t z_3$$

$${}^{t+\Delta t}z_2 = {}^t z_2 + \Delta t {}^t z_3 + \frac{1}{\tau_2} {}^t z_4$$

$${}^{t+\Delta t}z_3 = {}^t z_3 + \Delta t {}^t z_4 + \frac{1}{\tau_2} {}^t z_5$$

also

$${}^t z_5 = \tau_1 \left(-{}^{t-\Delta t}z_4 + {}^{t+\Delta t}z_4 \right)$$

$${}^t z_6 = \tau_0 \left({}^{t-\Delta t}z_4 - 2{}^t z_4 + {}^{t+\Delta t}z_4 \right)$$

4.4.2 Implicit Time Integration Scheme

For Implicit Time Integration scheme using Newmark's method, the dynamic equation of motion at time $t + \Delta t$ is solved as follows:

1. Select time step Δt
2. Compute integration parameters $\delta \geq 0.5$ $\lambda \geq 0.25 (0.5 + \delta)^2$

$$\tau_0 = \frac{1}{\lambda \Delta t^2}$$

$$\tau_1 = \frac{\delta}{\lambda \Delta t}$$

$$\tau_2 = \frac{1}{\lambda \Delta t}$$

$$\tau_3 = \frac{1}{2\lambda} - 1$$

$$\tau_4 = \frac{\delta}{\lambda} - 1$$

$$\tau_5 = \frac{\Delta t}{2} \left(\frac{\delta}{\lambda} - 2 \right)$$

$$\tau_6 = \Delta t(1 - \delta)$$

$$\tau_7 = \delta \Delta t$$

3. Initialize $z_1, z_2, z_3, z_4, z_5, z_6$

4. Form effective stiffness matrix

$$[\hat{K}] = [K] + \tau_0 [M] + \tau_1 [A_0]$$

5. Triangularize $[\hat{K}]$: $[\hat{K}] = [L][D][L]^T$

6. For each time step

a. Compute $z_1(t + \Delta t), z_2(t + \Delta t), z_3(t + \Delta t)$ from Taylor's expansion as follows

$${}^{t+\Delta t}z_1 = {}^tz_1 + \Delta t {}^tz_2 + \frac{\Delta t^2}{2} {}^tz_3$$

$${}^{t+\Delta t}z_2 = {}^tz_2 + \Delta t {}^tz_3 + \frac{\Delta t^2}{2} {}^tz_4$$

$${}^{t+\Delta t}z_3 = {}^tz_3 + \Delta t {}^tz_4 + \frac{\Delta t^2}{2} {}^tz_5$$

b. Compute effective loads at time $t + \Delta t$

$$\hat{\underline{F}}(t + \Delta t) = \underline{F}(t + \Delta t) + [M] \left(\tau_0^t z_4 + \tau_2^t z_5 + \tau_3^t z_6 \right) \\ + [A_0] \left(\tau_1^t z_4 + \tau_4^t z_5 + \tau_5^t z_6 \right) + [A_1]^{t+\Delta t} z_3 - [A_2]^{t+\Delta t} z_1$$

c. Solve for displacements at time $t + \Delta t$:

$$[L] [D] [L]^T {}^{t+\Delta t} z_4 = \hat{\underline{F}}(t + \Delta t)$$

d. Calculate accelerations and velocities at time $t + \Delta t$:

$${}^{t+\Delta t} z_6 = \tau_0 \left({}^{t+\Delta t} z_4 - {}^t z_4 \right) - \tau_2^t z_5 - \tau_3^t z_6 \\ {}^{t+\Delta t} z_5 = {}^t z_5 + \tau_6^t z_6 + \tau_7^{t+\Delta t} z_6$$

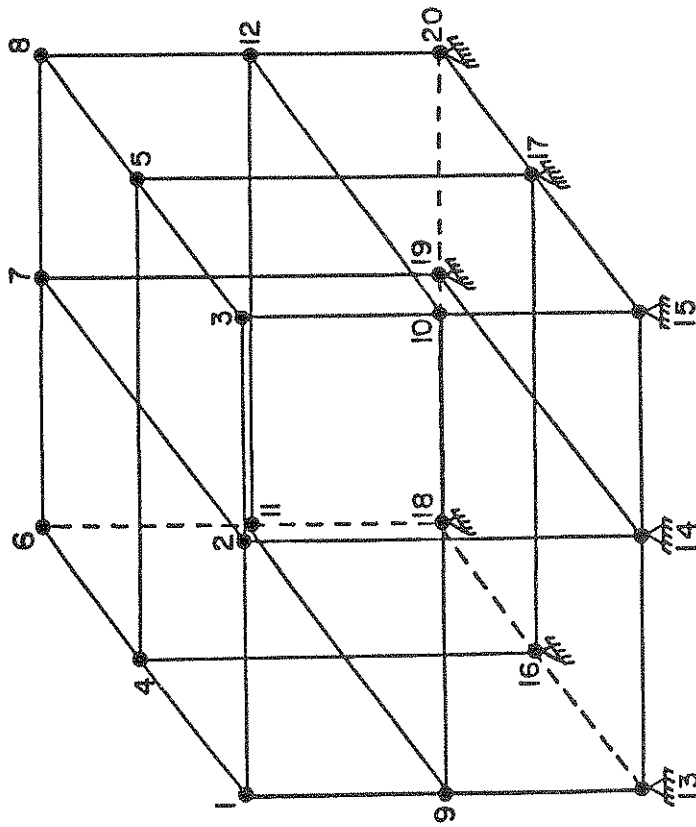
4.5 Numerical Example

The three-dimensional problem shown in Fig. 4-2A is used to test the new proposed model. The system consists of a medium discretized into three-dimensional Lagrangian elements whose interior nodes have been condensed out in the stiffness matrix formulation by the substructure technique (16). The surface nodes were subjected to a sinusoidal excitation as shown in Fig. 4-2A. The support conditions are also indicated in the figure. For clarity, only the displacement history in the direction of excitation (Z direction) for four nodes (nodes 1,2,3,4) on the surface were plotted. The displacement history at these nodes are sufficient to display the essential features of the proposed model. The cap model (1) was utilized to generate stress-strain relations from which an elastoplastic constitutive matrix was derived (see Section 3).

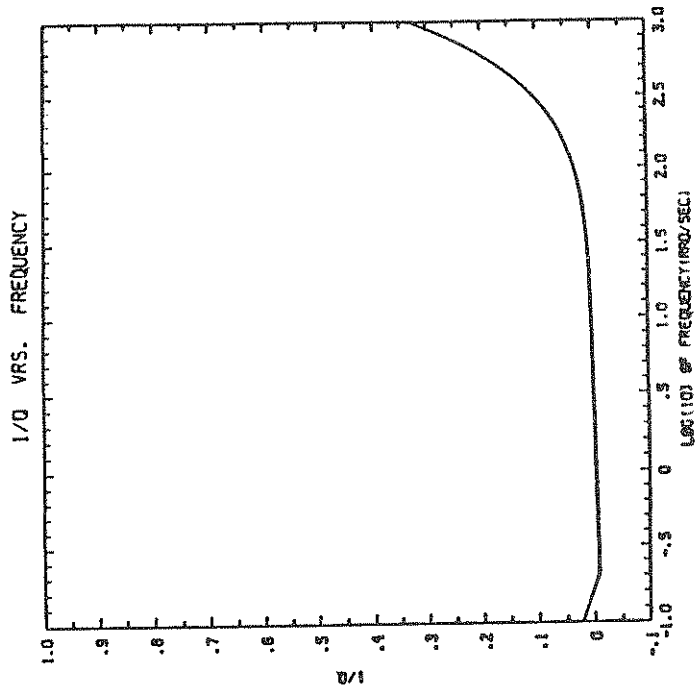
Four different tests were performed:

Test 1. This test was performed to investigate the effect of Q when the proposed viscoplastic model is adapted to predict responses in viscoelastic media. The elastoplastic constitutive matrix

$F = F_0 \sin \omega t$ (at nodes 1 through 8)



- a. NOTES:
1. ALL UNLABELLED NODES ARE CONDENSED OUT.
 2. SUPPORT CONDITIONS:
 - a) NODES 1 THROUGH 12 SUPPORTED IN X & Y DIRECTIONS.
 - b) NODES 13 THROUGH 20 SUPPORTED IN X, Y, Z DIRECTIONS.



b. $Q = 30$

FIGURE 4-2 Test Problem

was replaced by an elastic constitutive matrix in the stiffness matrix formulation. For a constant time step $\Delta t/T$ of 0.2 the value of the constant Q was varied from 10, 20, 30, 40, 50, 100, 500, 1000, using 200 time steps in the implicit time integration algorithm. The displacement history for the various values of Q are shown in Figure 4-3.

Test 2. For a constant time step $\Delta t/T = 0.1$ the effect of Q on displacement was observed for 200 time steps using the outlined implicit time integration algorithm. Here T denotes the fundamental period of the system under investigation. The value of Q was varied from 10, 20, 30, 40, 50, 100, 500, to 1000. The displacement history due to the applied sinusoidal excitation is shown in Figure 4-4.

Test 3. For a constant Q value of 30 the model was tested for convergence by varying the time step $\Delta t/T$ as follows: $\Delta t/T = 0.1, 0.08, 0.05, 0.025$ for 200 time steps using the proposed implicit time integration algorithm. The displacement history due to the applied sinusoidal excitation is shown in Figure 4-5.

Test 4. Test 4 was performed to investigate the stability of the model using the proposed explicit time integration algorithm. Here for a constant Q value of 30, the time step Δt was varied from $\Delta t/T = 0.01$ to 0.005 and the corresponding displacement history observed (Figure 4-6).

4.6 Discussions

From Figure 4-5, it is observed that for the implicit time integration scheme using Newmark's method, the model is stable for $\Delta t/T \leq 0.1$, where T is the fundamental period of the system. It is also noted that no further accuracy of the system response is achieved by using smaller time steps. On the other hand, for the explicit time integration scheme using the Central Difference scheme, a much smaller time step of $\Delta t/T \leq 0.01$ is required for stability (Figure 4-6). Clearly, the implicit time integration scheme is by far superior to the explicit time integration scheme in terms of savings in computer time.

The effect of the value of Q on the displacements history is clearly shown in Figs. 4-4. It is observed that larger values of Q result in less damping, hence larger displacements, and vice

versa. The choice of an appropriate value of Q for a particular material or medium is based on dynamic tests which can readily be performed in the laboratory. A Q value of 35 has been observed on laboratory tests performed on Ottawa sand in Ref. 29.

Fig. 4-2B shows a plot of $1/Q$ versus logarithm to base 10 of frequency for a constant Q value of 30.

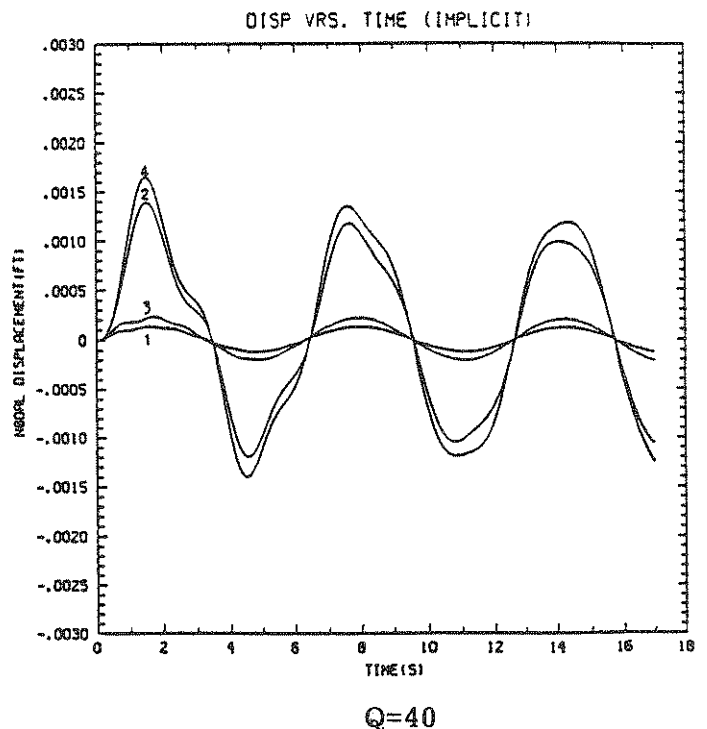
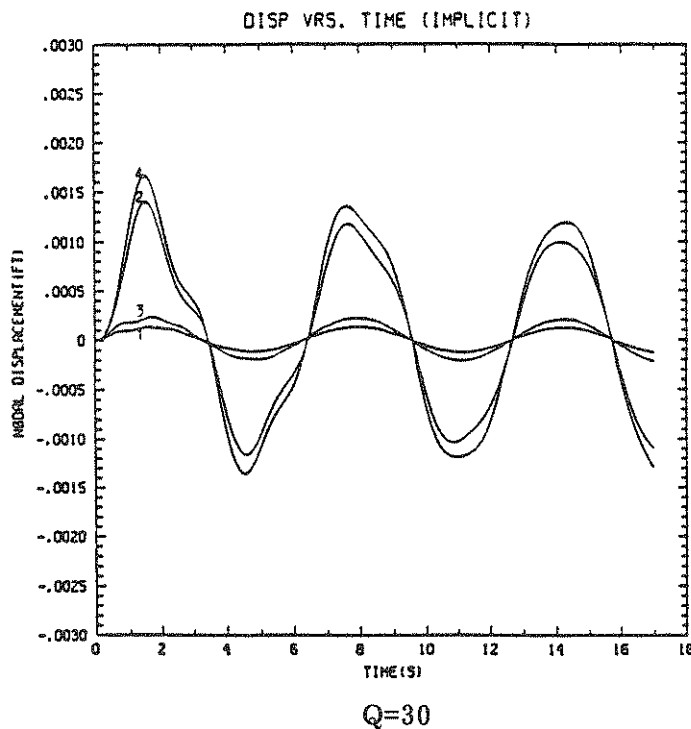
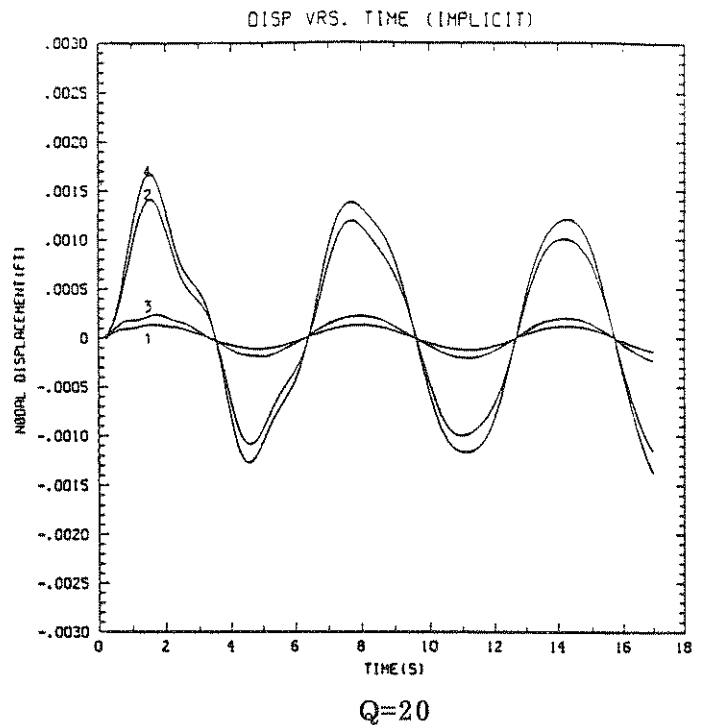
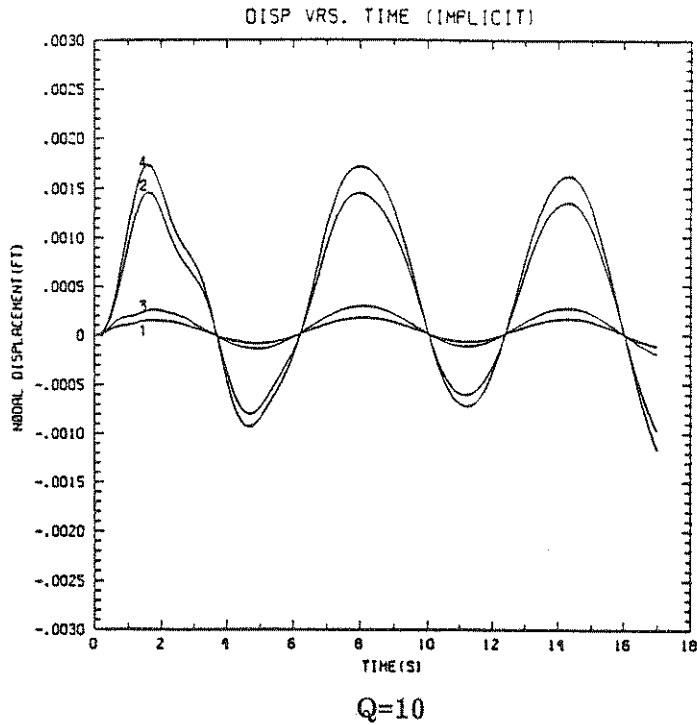
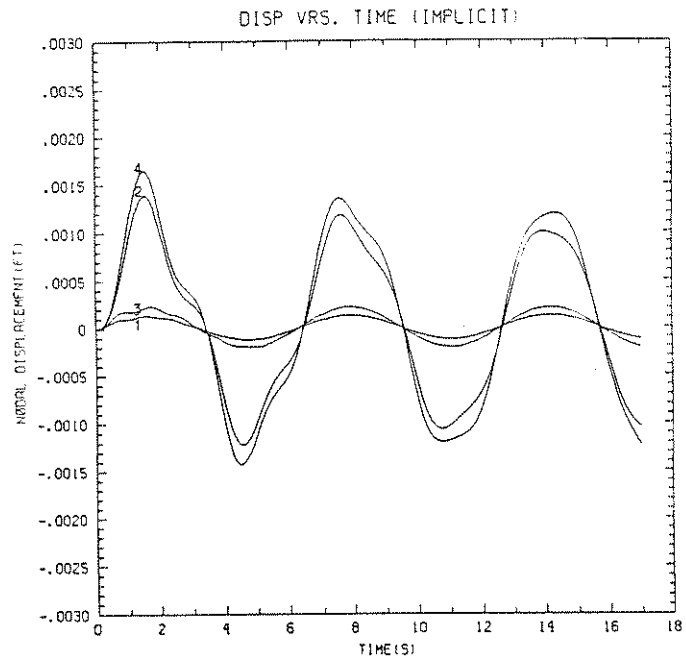
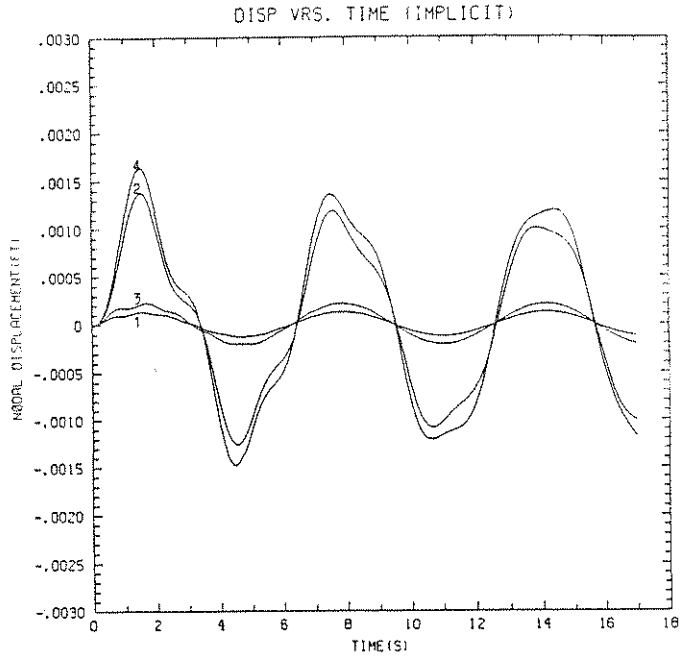


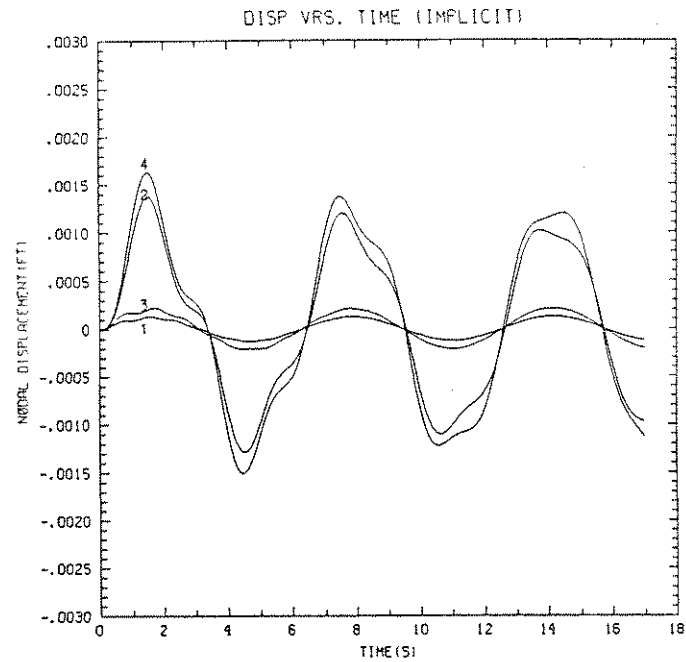
FIGURE 4-3 Test 1: Displacement History for Various Values of Q and a constant time step $\Delta t/T = 0.2$.



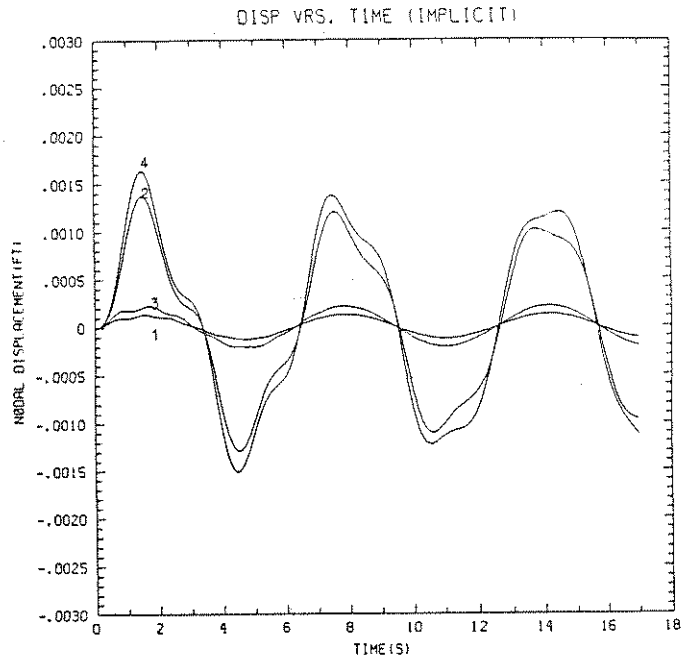
Q=50



Q=100

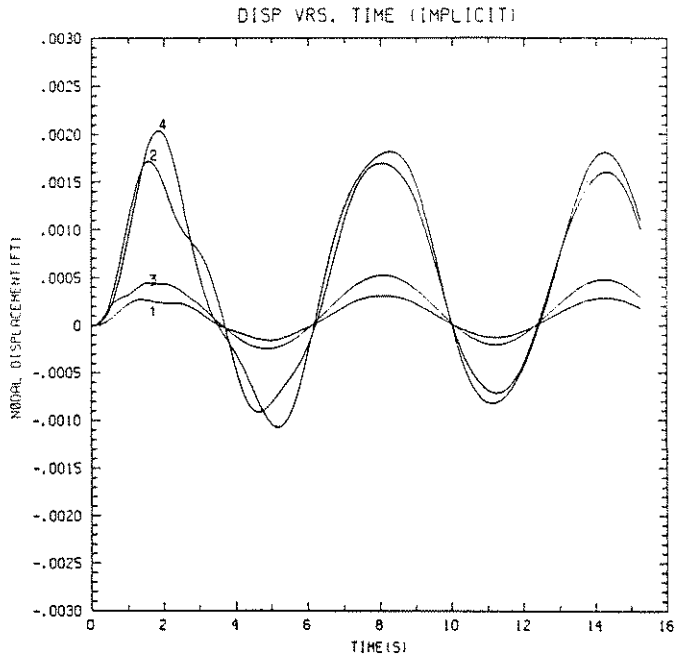


Q=500

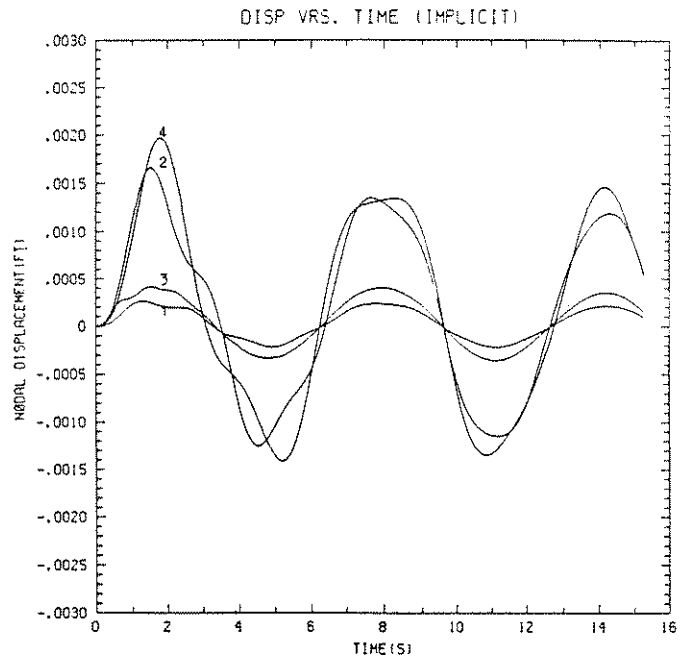


Q=1000

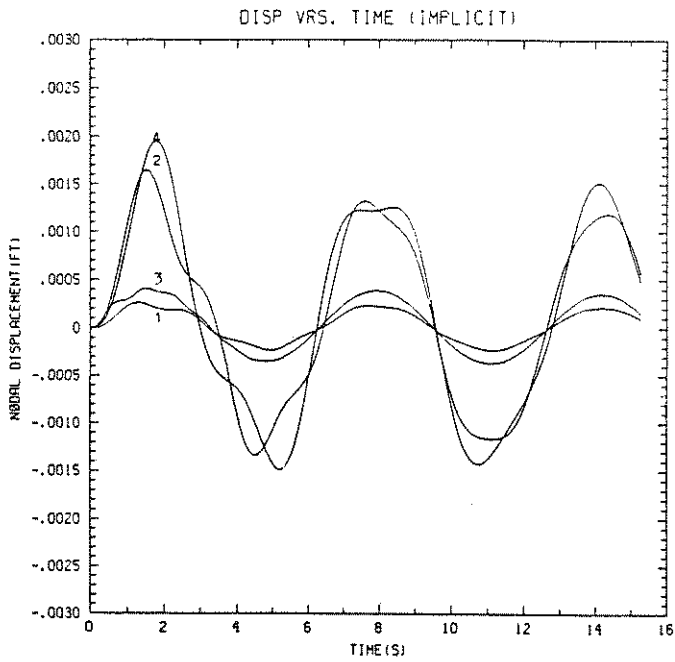
FIGURE 4-3 Test 1: Displacement History for Various Values of Q and a constant time step $\Delta t/T = 0.2$. (Cont'd)



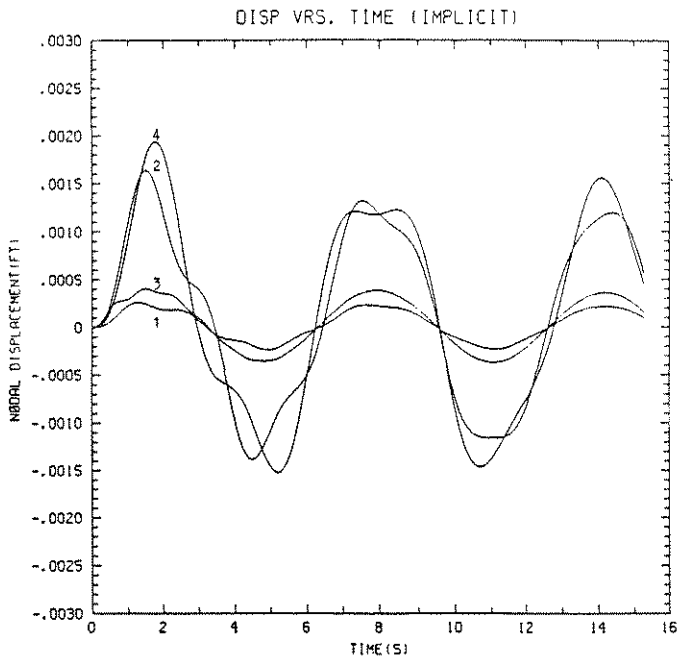
Q=10



Q=20

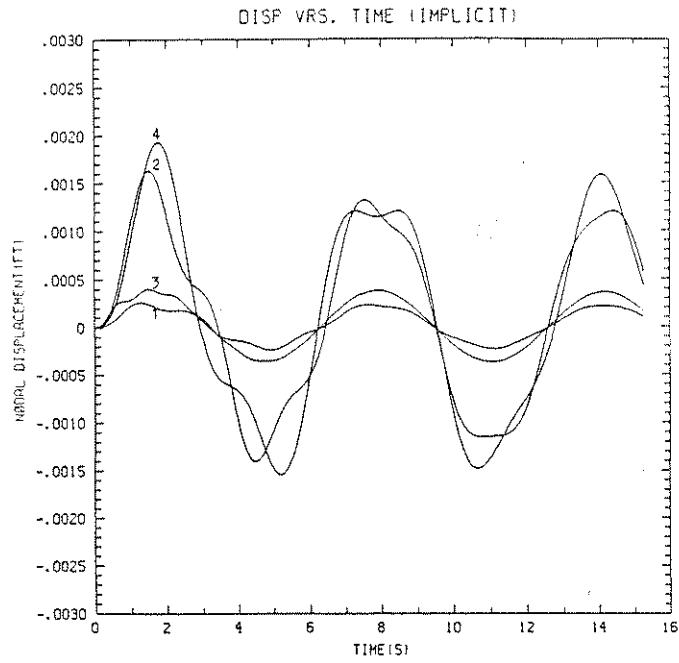


Q=30

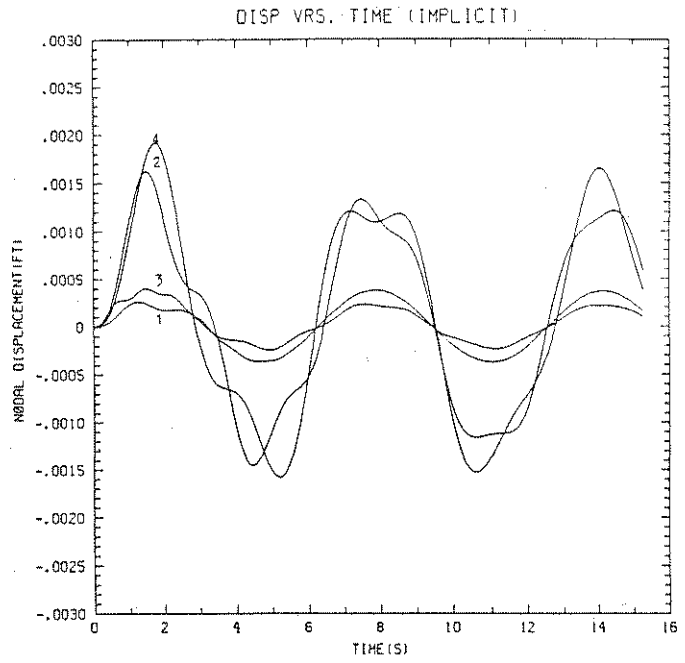


Q=40

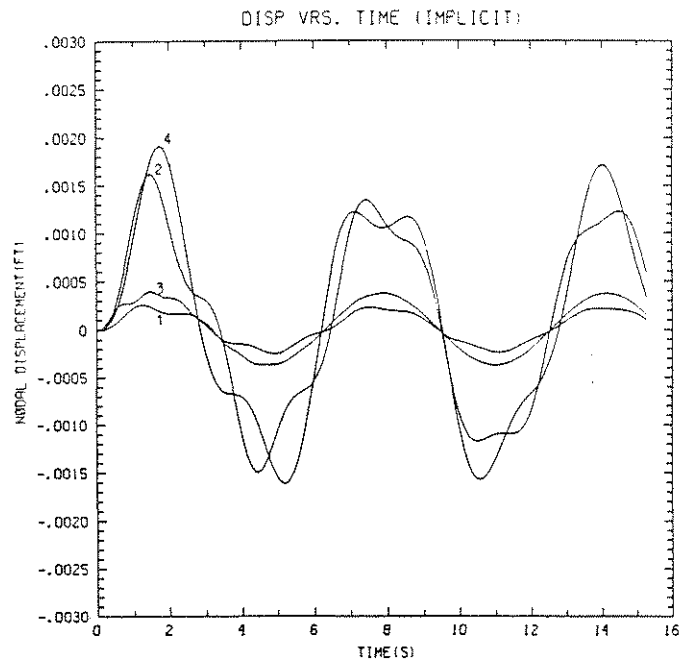
FIGURE 4-4 Test 2: Displacement History Due to Applied Sinusoidal Excitation for a constant time step $\Delta t/T = 0.1$



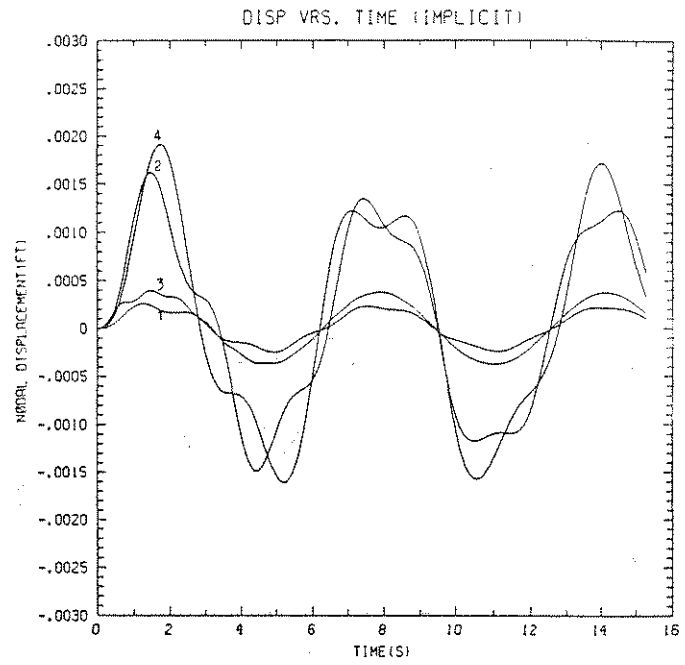
Q=50



Q=100

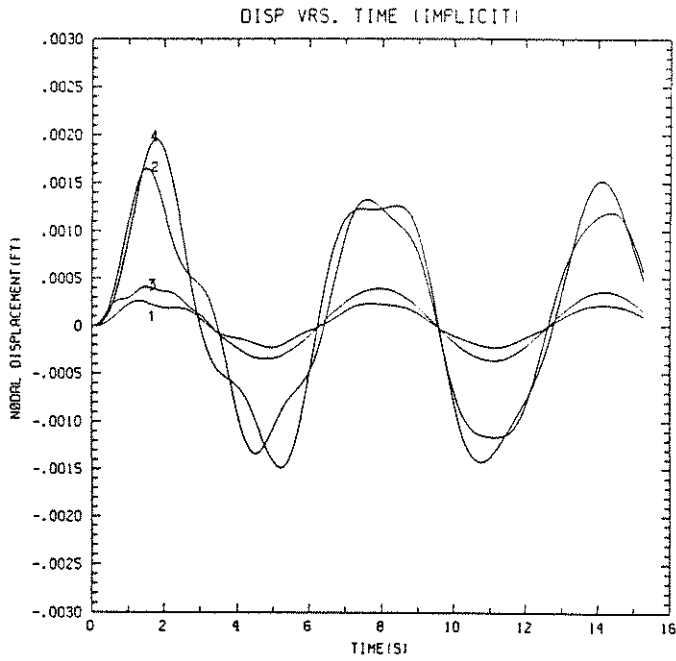


Q=500

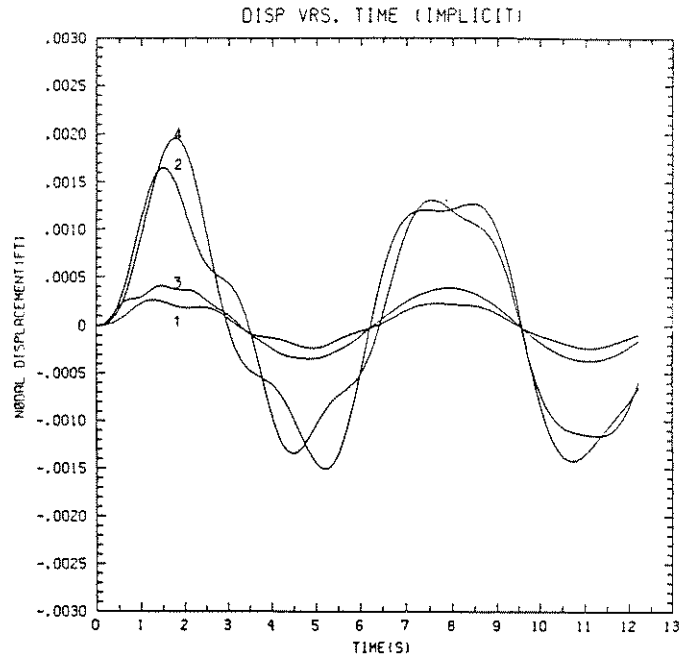


Q=1000

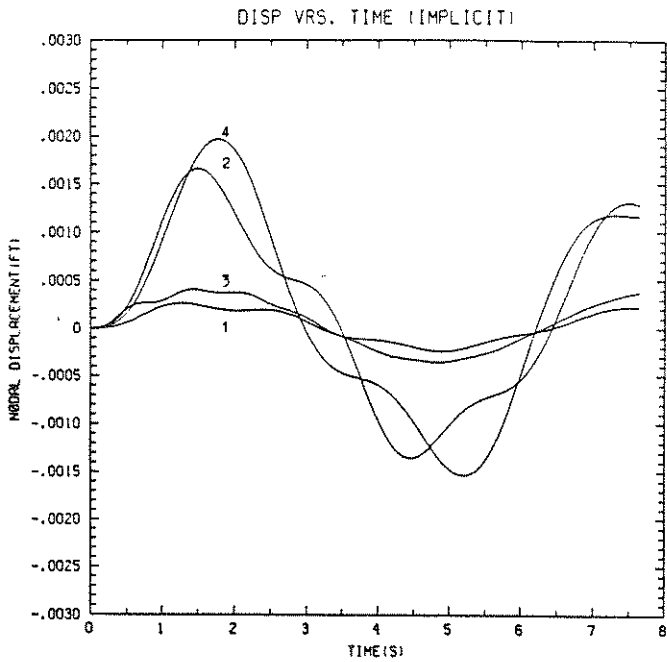
FIGURE 4-4 Test 2: Displacement History Due to Applied Sinusoidal Excitation for a constant time step $\Delta t/T = 0.1$. (Cont'd)



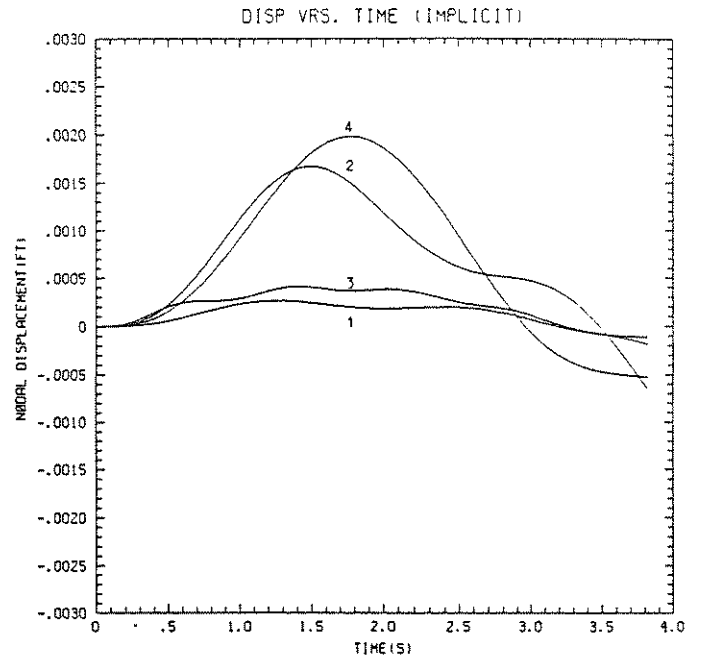
$\Delta t/T=0.1$



$\Delta t/T=0.08$

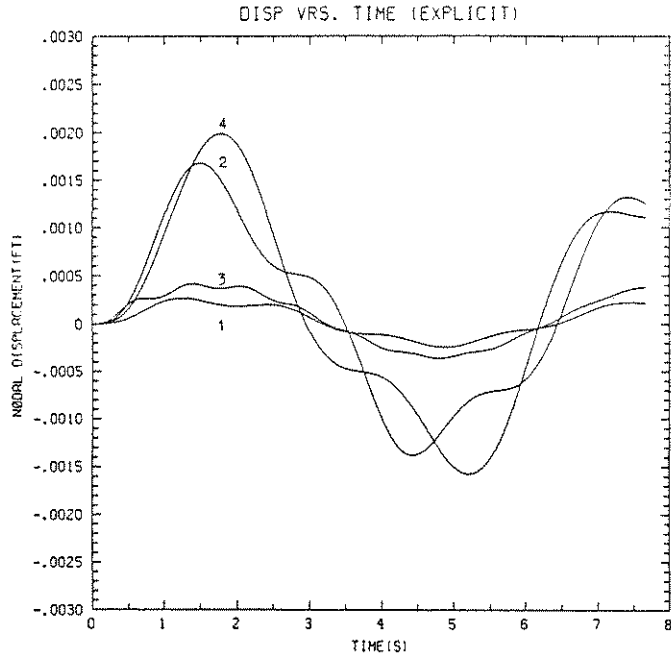


$\Delta t/T=0.05$

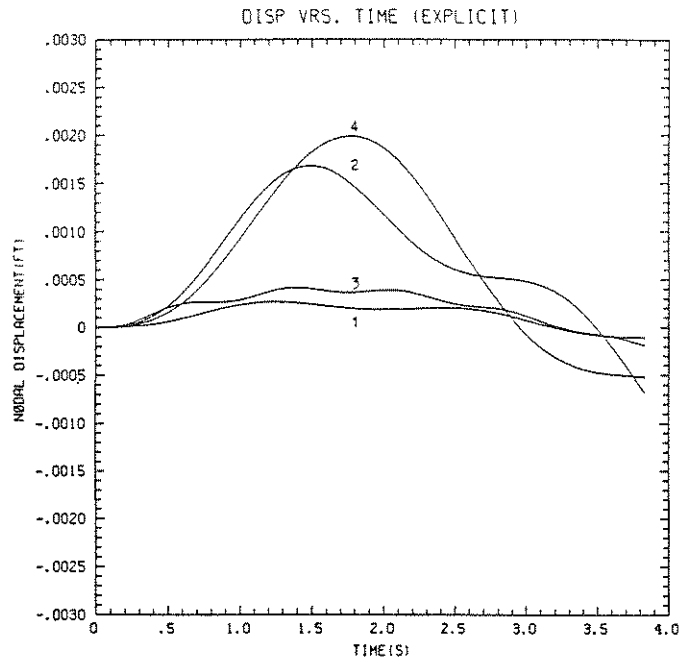


$\Delta t/T=0.025$

FIGURE 4-5 Test 3: Displacement History Due to Applied Sinusoidal Excitation for a constant Q value of 30.



$$\Delta t/T = 0.01$$



$$\Delta t/T = 0.005$$

FIGURE 4-6 Test 4: Displacement History for a constant Q model of 30 with the time step Δt varied from $\Delta t/T = 0.01$ to 0.005 .

SECTION 5 CONCLUSION

The proposed Q-model provides a means to conveniently establish a viscoplastic model which greatly enhances the analyses of earthquake, shock, blast and other vibration excitations often encountered in soil-structure interaction problems.

The frequency dependent $Q^{-1}(\omega)$ is expanded into a Laurent Series which generates a set of damping coefficients. For the special case when Q is constant over a prescribed frequency range, such damping coefficients are readily determined through minimization of the mean square error over the prescribed frequency range.

Numerical tests performed with Q constant over a frequency range of 0.1 to 10 radians/second show the model to be stable and accurate, provided the time step Δt is sufficiently small.

SECTION 6
REFERENCES

1. Sandler, I.S. and Rubin D., "An Algorithm and a Modular Subroutine for the Cap Model." International Journal for Numerical and Analytical Methods in Geomechanics, Vol 3, 1979, pp. 173-186.
2. Liu, H.-P., Anderson, D.L., and Kanamori, H., "Velocity Dispersion due to Anelasticity; Implications for Seismology and Mantle Composition." Geophys. J. Roy. astr. Soc (1976), Vol 47, pp. 41-58.
3. Kjartansson, E., "Constant Q-Wave Propagation and Attenuation." Journal of Geophysical Research, Vol 84, No. B9, Aug. 10, 1979, pp. 4737-4748.
4. Strick, E., "The Determination of Q, Dynamic Viscosity and Transient Creep Curves from Wave Propagation Measurements." Geophys. J. Roy. astr. Soc (1967), Vol. 13, pp. 197-218.
5. Naghdi, P.M. and Murch, S.A. "On the Mechanical Behavior of Viscoelastic/Plastic Solids. Journal Of Applied Mechanics, ASME, pp. 321-328, Sept. 1963.
6. Perzyna, R., "The Constitutive Equations for Work-Hardening and Rate Sensitive Plastic Materials." Proc. of Vibration Problems, Vols. 3,4, pp. 281-290, 1963 (in English).
7. Drucker, D.C., "On Uniqueness in the Theory of Plasticity." Quarterly of Applied Math, Vol 14, 1956, pp. 35-42.
8. Drucker, D.C., Gibson, R.E., and Henkel, D.J., "Soil Mechanics and Work-Hardening Theories of Plasticity." Transactions, ASCE, Vol. 122, 1957, pp. 338-346.
9. Drucker, D.C., and Palgen, L., "On Stress-Strain Relations Suitable for Cyclic and Other Loading." Journal of Applied Mechanics, Vol. 48, pp. 479-485 (1981).

10. Drucker, D.C. and Prager, W., "Soil Mechanics and Plasticity Analysis or Limit Design." Quarterly of Applied Mathematics, Vol. 10, 1952, pp. 157-175.
11. Sandler, I.S., DiMaggio, F.L., and Baron, M.L., "An Extension of the Cap Model- Inclusion of Pore Pressure Effects and Kinematic Hardening to Represent an Anisotropic Wet Clay." Mechanics of Engineering Materials, eds., Desai, C.S. and Gallagher, R.H., Chap. 28, John Wiley & Sons, Ltd. 1984.
12. Rosenblatt M., "A Set of Constitutive Equations for Rocks and Soils." Communication at D.A.S.A. Materials Working Group Meeting, 1970.
13. Reiner, M., "Plastic Yielding in Anelasticity." Journal of the Mechanics and Physics of Solids, Vol. 8, 1960 pp. 255-261.
14. Schofield, A.N., and Wroth, P., Critical State Soil Mechanics. McGraw-Hill, Ltd. 1968.
15. Norwich, A.S., and Berry, B.S., Anelastic Relaxation in Crystalline Solids. Academic Press, N.Y., 1972.
16. Zienkiewicz, O.C., The Finite Element Method, 3rd ed., McGraw - Hill, N.Y., 1977.
17. Bathe, K.J. and Wilson, E.L. Numerical Methods in Finite Element Analysis. Prentice-Hall, Englewood Cliffs, N.J. 1976.
18. Cristescu, N. and Suliciu, I., Viscoplasticity, Martinus Nijhoff Publishers, Boston 1982.
19. Aboim, C.A., and Roth, W.H., "Bounding-Surface-Plasticity Theory Applied to Cyclic Loading of Sand." International Symposium on Numerical Models in Geomechanics, Zurich, Sept. 1982, eds., Dungar, R., Pande, G.N., Studer, J.A.
20. Dafalias, Y.F., "An Elastoplastic -Viscoplastic Constitutive Modelling of Cohesive Soils." International Symposium on Numerical Models in Geomechanics, Zurich, 9/1982, eds. Dungar, R., Pande, G.N., Studer, J.A.

21. Valanis, K.C. and Read, A., "A New Endochronic Plasticity Theory for Soils." Systems, Science and Software Report SSS-R-80-4294 (1979).
22. Kondner, R.L., "Hyperbolic Stress-Strain Response of Soils" Proceedings of Soil Mechanics and Foundation Engineering, Division of ASCE, Vol. 89, 1963.
23. Dasgupta, G., "A Numerical Solution for Viscoelastic Half Planes." Journal of the Engineering Mechanics Division, ASCE, Vol. 102, No. EM4, Aug. 1976, pp. 601-612.
24. Dasgupta, G. and Chopra, A.K., "Dynamic Stiffness Matrix for Viscoelastic Half Planes," Journal of the Engineering Mechanics Division, ASCE, Vol. 105, No. EM5, Oct. 1979, pp. 729-745.
25. Zienkiewicz, O.C. and Corneau, I.C., "Visco-plasticity Plasticity and Creep in Elastic Solids - A Unified Numerical Solution Approach." International Journal for Numerical Methods in Engineering, Vol. 8, pp. 821-84b (1974).
26. Corneau, I., "Numerical Stability in Quasi-static Elasto-Viscoplasticity." International Journal for Numerical Method in Engineering, Vol. 9, pp. 109-127 (1975).
27. DiMaggio, F.L. and Sandler, I.S., "Material Model for Granular Soils." Journal of the Engineering Mechanics Division, ASCE, Vol. 97, No. EM3, Proc. Paper 8212, June 1971, pp. 935-950.
28. Day, S.M. and Minster, B.J., "Numerical Simulation of Attenuated Wavefields using Pade Approximant Method." Geophys. J. Roy. Astr. Soc. (1984), Vol. 78, pp. 105-118.
29. Richart, F.E., Woods, R.D. and Hall, J.R., Jr., "Vibrations of Soils and Foundations." Prentice-Hall, Inc., Englewood Cliffs, N.J., 1970, pp. 162-167.
30. Abramowitz, M. and Stegun, I.A., Handbook of Mathematical Functions. Dover Publications, Inc., N.Y., 1970.

C O M P U T E R

P R O G R A M

The Computer Program follows the following steps:

1. Evaluation of Stress-Strain Relation of Material Using an Appropriate Material Model. In the test problem the cap model(1) was used. The routine STRESS generates the stress-strain relation.
2. Evaluation of the 3-D Material Constitutive Matrix. The derivation of the elastoplastic constitutive matrix is shown in section 3 of the text. The subroutine DMATRX evaluates either elastic or elastoplastic constitutive matrix by setting the counter MMODE to 1 or 3 respectively.
3. Evaluation of 3-D Stiffness Matrix. The subroutine STIFF3D evaluates the 3-D stiffness matrix from a user supplied nodal coordinates of a discretized medium. In the test problem the medium was discretized into 3-D Lagrangian elements with three nodes in X,Y,Z directions - a total of 27 nodes (fig.13a). The interior nodes were then condensed out by the sub structure technique. Integration for stiffness matrix was performed over three Gauss points User has the option to supply the stiffness matrix, in which case steps 1, 2, and 3 can be omitted.
4. Evaluation of Mass Matrix. The subroutine GMASS evaluates the global mass matrix of the discretized medium from user supplied nodal coordinates. User has the option to supply the mass matrix.
5. Evaluation of Damping Coefficients. Damping coefficients are evaluated from the Q-Model by the subroutine QMOD. This routine is called by EXPLCT and IMPLCT for the evaluation of displacement history.
6. Evaluation of Response History. Two integration schemes have been presented. The sub-routines EXPLCT and IMPLCT evaluate the response history of the system by the explicit and implicit direct step by step time integration schemes respectively.

U S E R S U P P L I E D I N P U T

1. STIFFNESS AND MASS MATRIX EVALUATION.

STIFF-----Stiffness Matrix of discretized medium
(Elastic or Elastoplastic).
XMASS-----Mass Matrix of discretized medium
DM-----Material Constitutive Matrix. This is

supplied only when STIFF is not supplied
 by user.
 XXC,YYC,ZZC----Nodal Coordinates of spatially discret-
 ized medium. Omit when STIFF and XMASS
 are supplied by user.
 NCNSTR,JBNDRY--Number of Constraints in the spatially
 discretized medium. Omit when STIFF and
 XMASS are supplied by user.
 NGPS-----Number of Gauss points required for the
 formulation of the stiffness and mass
 matrices. Omit when STIFF and XMASS are
 user supplied.
 NODE-----Total number of nodes in the spatially
 discretized medium.
 NODB3-----3*NODE
 NDOF-----Number of degrees of freedom in the
 discretized medium.
 MX,MY,MZ-----Number of nodes in the X,Y,Z directions
 respectively for the Lagrangian element.

2. EVALUATION OF DAMPING COEFFICIENTS FROM THE Q-MODEL.

Q-----Specific Dissipation Factor of material
 OMEG1,OMEG2----Lower and Upper Frequency Limits for
 which Q is constant.

3. EVALUATION OF DYNAMIC RESPONSE.

CONST1,CONST2--Ratio of time step and fundamental period
 of the system for the explicit and
 implicit time integration schemes resp-
 ectively. These constants help to select
 appropriate time steps for stability and
 accuracy.
 ICOUNT-----A Counter which indicates whether explicit
 or implicit time integration scheme is
 used.
 ICOUNT=0 EXPLICIT
 ICOUNT=1 IMPLICIT
 MTIME1,MTIME2--Number of time steps required for the
 system response for the explicit and
 implicit time integration schemes
 respectively.

S A M P L E I N P U T

| | | | | | | | | | | | | |
|------|-----|-----|-----|-----|------|-------|------|---|----|-------|-----|-----|
| .250 | .65 | .18 | .67 | 2.5 | .066 | 66.67 | 40.0 | 1 | .3 | .4903 | .25 | 001 |
| 0. | 0. | 0. | | | | | | | | | | 002 |
| 1. | 0. | 0. | | | | | | | | | | 003 |
| 2. | 0. | 0. | | | | | | | | | | 004 |


```

C
C   READ(5,*)(PROP(I),I=1,12)
C
C   INPUT NODAL COORDINATES
C
C   DO 1 I=1,NODE
1  READ(5,*)XXC(I),YYC(I),ZXC(I)
C
C   READ(5,*)MAXITR
C
C   READ(5,*)NCNSTR,ICOUNT,CONST1,MTIME1,CONST2,MTIME2,NGP
C
C   READ(5,*)Q,OMEG1,OMEG2
C
C   INPUT BOUNDARY CONSTRAINTS
C
C   READ(5,*)(JBNDRY(I),I=1,NCNSTR)
C   PRINT*,'JBNDRY(I)=',(JBNDRY(I),I=1,NCNSTR)
C
C   MMAT=1
C   VKAPA=0.13304
C
C*****
C
C   COMPUTE STRESS-STRAIN RELATIONS
C   -----
C
C   DECLARE INITIAL STRESSES
C
C   DO 3 I=1,6
3  SIGIJ(I,1)=0.
C
C
C   DO 4 ITER1=1,MAXITR
C
C   CALL STRESS(SIGIJ,MMAT,VKAPA,PROP,SIGMA,DFDSIG,ITER1,
,      MMODE,TTWOG,DFDKAP)
C
C   DO 5 II=1,6
5  SIGIJ(II,1)=SIGMA(II,1)
C
C   SIGMX=SIGMA(1,1)
C   SIGMY=SIGMA(2,1)
C   SIGMZ=SIGMA(3,1)
C   SIGMXY=SIGMA(6,1)
C   SIGMKZ=SIGMA(5,1)
C   SIGMYZ=SIGMA(4,1)

```



```

C          REDUCE MASS MATRIX IN ACCORDANCE WITH
C          PRESCRIBED BOUNDARY CONDITIONS
C
C          CALL REMOVE(ZMASS,NODB3,XMASS,NDOF,JBNDRY,NCNSTR,
C          ,
C          ,      JSIZE)
C
C*****
C
C          WRITE(7,1006)
C          WRITE(7,1007)((STIFF(II,JJ),JJ=1,NDOF),II=1,NDOF)
C
C          WRITE(7,1008)
C          WRITE(7,1009)((XMASS(II,JJ),JJ=1,NDOF),II=1,NDOF)
C
C*****
C
C          COMPUTE THE DYNAMIC RESPONSE OF MATERIAL
C          -----
C
C          TO PERFORM EITHER EXPLICIT OR IMPLICIT TIME
C          INTEGRATION. ICOUNT IS SET TO 0 OR 1 RESP.
C
C          IF(ICOUNT.EQ.0) THEN
C
C          CALL EXPLCT(XMASS,STIFF,NDOF,CONST1,MTIME1,WWRK1,
C          , WWRK2,WWRK3,WWRK4,WWRK5,WWRK8,WWRK10,WWRK11,
C          , WWRK12,WWRK13,ZZ1,ZZ2,ZZ3,ZZ4,ZZ4I,
C          , ZZ5,ZZ6,ZZ66,AA0,AA1,AA2,FFORCE,EFMAS)
C
C          ELSE
C
C          CALL IMPLCT(XMASS,STIFF,NDOF,CONST2,MTIME2,WWRK1,
C          , WWRK2,WWRK3,WWRK4,WWRK5,WWRK6,WWRK7,WWRK8,WWRK9,
C          , WWRK10,WWRK11,WWRK12,ZZ1,ZZ2,ZZ3,ZZ4,ZZ5,ZZ6,
C          , ZZ66,AA0,AA1,AA2,FFORCE,EFK)
C
C          ENDIF
C
C*****

```



```

C  ITERA-----COUNTER FOR CAP ITERATION
C  MAT-----TYPE OF MATERIAL
C           MAT=1  SOIL
C           MAT=2  ROCK
C  MODE-----MODE OF MATERIAL BEHAVIOR
C           MODE=0  TENSION MODE
C           MODE=1  ELASTIC MODE
C           MODE=2  FAILURE MODE
C           MODE=3  CAP MODE
C  NIT-----MAXIMUM NUMBER OF ITERATIONS FOR CAP CONVERGENCE
C  SHEAR-----SHEAR MODULUS OF MATERIAL
C  SIGKK0-----INITIAL VOLUMETRIC STRESS
C  SIGM-----TOTAL STRESS VECTOR
C  STR-----DEVIATORIC STRESS VECTOR
C  TCUT-----TENSION CUT-OFF (MAXIMUM TENSION PERMITTED)
C  VARJ1-----FIRST INVARIANT OF STRESSES
C  VARJ2-----SECOND INVARIANT OF DEVIATORIC STRESSES
C  YA, YB, YC, YD, YR, YW, ----MATERIAL CONSTANTS IN CAP MODEL
C                               (OBTAINED EXPERIMENTALLY)
C  YNU-----POISSON'S RATIO
C  YPROP-----VECTOR CONTAINING THE MATERIAL CONSTANTS IN CAP MODEL
C  X-----FUNCTION DEFINING POSITION OF ELLIPTIC CAP
C
C
C *****
C
C  IMPLICIT REAL*8(A-H,O-Z)
C  DIMENSION STR(6,1),DEPS(6,1),STR0(6,1),DFDS(6,1),SIGM(6,1),
C  +          YPROP(6)
C ***** FUNCTION STATEMENTS *****
C
C  THE FOLLOWING ARE STATEMENT FUNCTIONS DEFINING
C  CAP MODEL EQUATIONS
C
C  EXPS(Z)=DMAX1(DBLE(-500.),DEXP(Z))
C
C  FAILURE ENVELOPE FUNCTION FOR VARJ2  "SQRT J2PRIME"
C
C  F(VARJ1)=YA-YC*EXPS(YB*VARJ1)
C
C  CAP STATEMENT FUNCTIONS (CAPL=BIGL(HKAPA),XX=X(HKAPA))
C
C  BIGL(HKAPA)=DMIN1(DBLE(0.0),HKAPA)
C  R(CAPL)=YR
C  X(HKAPA)=HKAPA-R(BIGL(HKAPA))*F(HKAPA)
C  EVP(XX)=YW*(EXPS(YD*XX)-1.0)
C  FCAP(VARJ1,XX,CAPL)=DSQRT(DABS((XX-CAPL)**2-(VARJ1-CAPL)**2))/
C  /R(CAPL)
C
C  ELASTIC MODULI FUNCTIONS (EV IS CURRENT VALUE OF EVP(XX))-
C
C  BMOD(VARJ1,EV)=BULK
C  SMOD(VARJ2,EV)=SHEAR
C *****
C ***** END OF FUNCTION STATEMENTS *****
C
C

```

```

C
C
YA=YPROP(1)
YB=YPROP(2)
YC=YPROP(3)
YD=YPROP(4)
YR=YPROP(5)
YW=YPROP(6)
BULK=YPROP(7)
SHEAR=YPROP(8)
MAT=YPROP(9)
TCUT=YPROP(10)
FCUT=YPROP(11)
YNU=YPROP(12)
C
CONV=.001
NIT=60
GEOP=0.
ITERA=0
C
C
C
C
DO 1 I=1,6
DEPS(I,1)=0.0
DFDS(I,1)=0.0
STR(I,1)=0.0
1 CONTINUE
C
IF(ITER.LE.10)DEPS(3,1)=-0.005
IF(ITER.GT.10.AND.ITER.LE.15)DEPS(3,1)=0.002
IF(ITER.GT.15)DEPS(3,1)=0.001
C
C
C
THE FOLLOWING ARE EQNS 15,16,17 RESP.
CAPL=BIGL(HKAPA)
XX=X(HKAPA)
EVPI=EVP(XX)
C
C CALCULATE DEVIATORIC STRAIN INCREMENTS
C
DEV=DEPS(1,1)+DEPS(2,1)+DEPS(3,1)
DEVB3=DEV/3.0
DEPX=DEPS(1,1)-DEVB3
DEPY=DEPS(2,1)-DEVB3
DEPZ=DEPS(3,1)-DEVB3
C
C CALCULATE INITIAL DEVIATORIC STRESSES
C
SIGKK0=(STR0(1,1)+STR0(2,1)+STR0(3,1))/3.0
STR0(1,1)=STR0(1,1)-SIGKK0
STR0(2,1)=STR0(2,1)-SIGKK0
STR0(3,1)=STR0(3,1)-SIGKK0
C
C INITIAL STRESS INVARIANTS
C
VARJ1I=3.*(SIGKK0+GEOP)
VARJ2I=DSQRT((STR0(1,1)**2+STR0(2,1)**2+STR0(3,1)**2+
+2.*STR0(4,1)**2+2.*STR0(5,1)**2+2.*STR0(6,1)**2)/2.)

```

```

C
C   ELASTIC MATERIAL PROPERTIES
C
C   THREEK=3.*BMOD(VARJ1I,EVPI)
C   TWOG=2.*SMOD(VARJ2I,EVPI)
C
C*****
C   ELASTIC TRIAL
C*****
C
C   STR(1,1)=STRO(1,1)+TWOG*DEPX
C   STR(2,1)=STRO(2,1)+TWOG*DEPY
C   STR(3,1)=STRO(3,1)+TWOG*DEPZ
C   STR(4,1)=STRO(4,1)+TWOG*DEPS(4,1)
C   STR(5,1)=STRO(5,1)+TWOG*DEPS(5,1)
C   STR(6,1)=STRO(6,1)+TWOG*DEPS(6,1)
C
C   RATIO=1.0
C   MODE=1
C
C   STRESS INVARIANTS
C
C   VARJ1=THREEK*DEV+VARJ1I
C   VARJ2=DSQRT((STR(1,1)**2.+STR(2,1)**2.+STR(3,1)**2.+
C+2.*STR(4,1)**2.+2.*STR(5,1)**2.+2.*STR(6,1)**2.)/2.)
C*****
C   TENSILE CODING
C*****
C
C   TENCUT=DMIN1(FCUT,TCUT+3.*GEOP)
C
C   THE FOLLOWING CONDITION APPLIES
C   WHEN TENSION LIMIT IS NOT EXCEEDED
C
C   IF(VARJ1.LT.TENCUT)GO TO 10
C
C   THE FOLLOWING CONDITION APPLIES
C   WHEN TENSION LIMIT IS EXCEEDED
C   THAT IS SET J1=T
C
C   VARJ1=TENCUT
C   RATIO=0.0
C   MODE=0
C
C   - CONDITION FOR EITHER ROCK MODEL OR
C   KAPPA'(DOT) IS.GE.ZERO
C
C   IF(MAT.EQ.2.OR.HKAPA.GE.0.0)GO TO 200
C
C   TENSION DILATANCY CODING
C
C   HKAPA1=DMAX1(DBLE(0.0),HKAPA+CONV*F(HKAPA))
C
C   EQUATION 16 FOLLOWS
C

```

```

      XXL=X(HKAP1)
      DENOM=EVP(XXL)-EVPI
      IF(DENOM.GT.0.0)GO TO 5
      HKAPA=0.0
      GO TO 200
C
C      EQN. 22 FOLLOWS
C
C      DEVP=DEV-(VARJ1-VARJ1I)/THREEK
C
C      HKAPA=HKAPA+DEVP*(HKAP1-HKAPA)/DENOM
C
C      EQNS (18), & (20) FOLLOW
C
C      HKAPA=DMIN1(DBLE(0.0),HKAPA)
C
C      GO TO 200
C
C*****
C      CHECK FAILURE ENVELOPE      *
C*****
C
C10    CONTINUE
C
C      THE FOLLOWING CONDITION IMPLIES FAILURE
C      MODE DOES NOT APPLY
C
C      IF(VARJ1.LT.CAPL)GO TO 30
C
C      VON MISES TRANSITION
C
C      VONMIS=FCAP(CAPL,XX,CAPL)
C      FJ1=F(VARJ1)
C      FF=VARJ2-DMIN1(FJ1,VONMIS)
C      IF(FF.LE.0.0)GO TO 200
C
C      FAILURE SURFACE CALLCULATION
C
C      MODE=2
C      DFDJ1=0.0
C
C      CALLCULATION OF DFE/DJ1 @ J1E
C
C      IF(FJ1.LT.VONMIS)DFDJ1=(FJ1-F(VARJ1+CONV*VARJ2))/(CONV*VARJ2)
C
C      EQN (33) FOLLOWS
C
C      DEVP=3.*DFDJ1*FF/(3.*THREEK*DFDJ1**2+0.5*TWOG)
C
C      VARJ1=VARJ1-THREEK*DEVP
C
C      DILATANCY AND CORNER CODING
C
C      IF(MAT.EQ.1.AND.HKAPA.LT.0.0.AND.VARJ1.GT.CAPL)GO TO 60
C      VARJ1=DMAX1(VARJ1,CAPL)
C      GO TO 70
60    CONTINUE
C

```

```

C      EQN (37) FOLLOWS
C
      HKAP1=BIGL(VARJ1)
      XXL=X(HKAP1)
      IF(DEVP.LE.0.0)GO TO 70
      DEVPT=DMAX1(DEVP,EVP(XXL)-EVPI)
      HKAPA=HKAPA+(HKAP1-HKAPA)*DEVP/DEVPT
C
70    CONTINUE
      FJ1=F(VARJ1)
      RATIO=DMIN1(FJ1,VONMIS)/VARJ2
C
C*****
C
C          CAP CALCULATION
C
C*****
C
30    CONTINUE
C
C          CONDITION FOR CAP MODE COMPUTATION
C
      IF(VARJ1.LT.XX)GO TO 40
      IF(VARJ2.LE.FCAP(VARJ1,XX,CAPL))GO TO 200
40    CONTINUE
C
      VARJ1E=VARJ1
      VARJ2E=VARJ2
C
C          AN INITIAL VALUE FOR THE HARDENING
C          PARAMETER KAPPA IS ASSUMED AND THEN
C          REFINED USING THE REGULA FALSI
C          ITERATION PROCEDURE
C
      HKAP1=HKAPA
      HKAPA2=VARJ1E
      IF(VARJ1E.LE.XX)FL=(HKAPA-VARJ1E)/(HKAPA-XX)
      IF(VARJ1E.GT.XX)FL=2.*VARJ2E/(VARJ2E+FCAP(VARJ1E,XX,CAPL))-1.0
C
      XR=X(HKAPA2)
      VARJ1R=VARJ1E-THREEK*(EVP(XR)-EVPI)
      FR=(XR-VARJ1R)/(HKAPA2-XR)
      COMP=CONV*F((FL*XR-FR*XX)/(FL-FR))
      IF(DABS(VARJ1)+VARJ2.LT.COMP)GO TO 200
C
      MODE=3
      FOLD=0.0
C
C*****
C
C          ITERATION FOR CAP CONVERGENCE BEGINS
C
C*****
C
      DO 190 ITERA=1,NIT
      HKAPA=(FL*HKAPA2-FR*HKAP1)/(FL-FR)
      XX=X(HKAPA)
      DEVP=EVP(XX)-EVPI

```



```

VARJ1=VARJ1E-THREK*DEVP
CAPL=BIGL(HKAPA)
IF(VARJ1.LE.XX)FC=(HKAPA-VARJ1)/(HKAPA-XX)
IF(VARJ1.GE.CAPL)FC=(XX-VARJ1)/(CAPL-XX)
IF(VARJ1.LE.XX.OR.VARJ1.GE.CAPL)GO TO 300
C
VARJ2=FCAP(VARJ1,XX,CAPL)
DELJ1=CONV*(XX-VARJ1)
DESP=0.0
IF(VARJ1+DELJ1.NE.VARJ1)DESP=(DEVP/6.)*(DELJ1/(VARJ2-
-FCAP(VARJ1+DELJ1,XX,CAPL)))
VARJ2T=VARJ2+TWOG*DESP
FC=(VARJ2E-VARJ2T)/(VARJ2E+VARJ2T)
C
C
C*****
C
C          OUTLET OF CAP ITERATION LOOP FOLLOWS
C
C          IF(DABS(VARJ2E-VARJ2T).LE.COMP)GO TO 195
C          IF(FC.GT.0.0.AND.VARJ1-CAPL.GE.DELJ1)GO TO 195
C*****
C
C 300 IF(FC.GT.0.)GO TO 320
C
C      HKAPA2=HKAPA
C      FR=FC
C      IF(FOLD.LT.0.)FL=0.5*FL
C      GO TO 190
C 320 CONTINUE
C
C      HKAPA1=HKAPA
C      FL=FC
C      IF(FOLD.GT.0.)FR=0.5*FR
C 190 FOLD=FC
C
C
C
C      VARJ1=DMAX1(VARJ1,XX)
C      IF(VARJ1.GT.BIGL(HKAPA2))VARJ1=CAPL
C      VARJ2=DMIN1(VARJ2E,FCAP(VARJ1,XX,CAPL))
C 195 RATIO=0.0
C      IF(VARJ2E.NE.0.)RATIO=VARJ2/VARJ2E
C 200 CONTINUE
C
C
C          COMPUTE FINAL DEVIATORIC STRESSES
C
C      STR(1,1)=STR(1,1)*RATIO
C      STR(2,1)=STR(2,1)*RATIO
C      STR(3,1)=STR(3,1)*RATIO
C      STR(6,1)=STR(6,1)*RATIO
C      STR(5,1)=STR(5,1)*RATIO
C      STR(4,1)=STR(4,1)*RATIO
C
C
C          COMPUTE FINAL TOTAL STRESSES
C
C      SIGKB3=VARJ1/3.-GEOP

```

```

C
SIGM(1,1)=STR(1,1)+SIGKB3
SIGM(2,1)=STR(2,1)+SIGKB3
SIGM(3,1)=STR(3,1)+SIGKB3
SIGM(4,1)=STR(4,1)
SIGM(5,1)=STR(5,1)
SIGM(6,1)=STR(6,1)

C
C
C
XX=X(HKAPA)
CAPL=BIGL(HKAPA)
IF(MODE.LE.1) GO TO 99
IF(MODE.EQ.2) THEN          !FAILURE MODE
DFDJ1=YB*YC*EXPS(YB*VARJ1)
DFDK=0.0
ENDIF

C
IF(MODE.EQ.3) THEN          !CAP MODE
DFDJ1=2.*(VARJ1-CAPL)/(YR*YR*VARJ2)
DFDX=(CAPL-XX)/(YR*YR*VARJ2)
DXDK=1.+(YR*YC*YB*(EXPS(YB*HKAPA)))
DFDL=(XX-VARJ1)/(YR*YR*VARJ2)
IF(HKAPA.GE.0.) THEN
DLDK=0.
ELSE
DLDK=1.
ENDIF
DFDK=DFDX*DXDK+DFDL*DLDK
ENDIF

C
IF(VARJ2.LT.1.E-25) GO TO 99
FCAP2=2.*VARJ2
DFDS(1,1)=STR(1,1)/FCAP2+DFDJ1
DFDS(2,1)=STR(2,1)/FCAP2+DFDJ1
DFDS(3,1)=STR(3,1)/FCAP2+DFDJ1
DFDS(4,1)=STR(4,1)/FCAP2
DFDS(5,1)=STR(5,1)/FCAP2
DFDS(6,1)=STR(6,1)/FCAP2

C
C
C
99 CONTINUE

C
RETURN
END

SUBROUTINE DMATRIX(SIGM,XPROP,MODE,DF,DFDK,TWOG,D)
C*****
C
C   THIS SUBROUTINE COMPUTES THE ELASTO-PLASTIC
C   MATRIX FROM THE CAP MODEL
C
C ARGUMENTS :
C
C   E , XNU ----- YOUNG'S MODULUS AND POISSON'S

```



```

      IF(I.EQ.KBNDRY(L))GO TO 1
10  CONTINUE
C
C
C
      N=1
      DO 2 J=1,NODB3
C
      DO 11 L=1,KONSTR
      IF(J.EQ.KBNDRY(L))GO TO 2
11  CONTINUE
C
C
C
      YM(M,N)=XM(I,J)
C
      N=N+1
C
      2  CONTINUE
C
      M=M+1
C
      1  CONTINUE
C
      M=M-1
      N=N-1
      ISIZE=M
C
      RETURN
      END

```

```

      SUBROUTINE TRANSP(AA,BB,IROW,ICOL)
C
C*****
C
C          THIS SUBROUTINE TRANSPOSES
C          MATRICES
C
C          [BB]=[AA] TRANSPOSE
C
C ARGUMENTS ARE:
C   AA----- (INPUT) MATRIX TO BE TRANSPOSED
C   BB----- (OUTPUT) TRANSPOSED MATRIX
C   IROW,ICOL---NUMBER OF ROWS AND COLUMNS RESP.
C                 OF THE INPUT MATRIX AA
C
C
C THIS SUBROUTINE IS CALLED BY:
C                                     (1) MAIN PROGRAM
C                                     (2) DMATRX
C THIS SUBROUTINE CALLS : NONE
C*****
C
C
C
      IMPLICIT REAL*8(A-H,O-Z)

```

```

C      DIMENSION AA(IROW,ICOL),BB(ICOL,IROW)
C
C      DO 1 M=1,IROW
C      DO 2 J=1,ICOL
C      BB(J,M)=AA(M,J)
C      2 CONTINUE
C      1 CONTINUE
C
C      RETURN
C      END

C      SUBROUTINE MXMULT(AA,BB,CC,IROWA,ICOLA,IROWB,ICOLB)
C      *****
C      THIS SUBROUTINE MULTIPLIES TWO MATRICES
C      ARGUMENTS ARE:
C      [CC] = [AA][BB]
C      IROWA & IROWB ARE NUMBER OF ROWS IN
C      THE MATRICES [AA] , [BB] RESP.
C      ICOLA & ICOLB ARE NUMBER OF COLUMNS
C      IN MATRICES [AA] , [BB] RESP.
C
C      THIS SUBROUTINE IS CALLED BY :
C      (1) MAIN PROGRAM*
C      (2) DMATRX
C
C      THIS SUBROUTINE CALLS :      NONE
C      *****
C
C      IMPLICIT REAL*8(A-H,O-Z)
C      DIMENSION AA(IROWA,ICOLA),BB(IROWB,ICOLB),
C      CC(IROWA,ICOLB)
C
C      DO 1 M=1,IROWA
C      DO 2 J=1,ICOLB
C
C      SUM=0.
C
C      DO 3 JJ=1,ICOLA
C      CC(M,J)=AA(M,JJ)*BB(JJ,J)
C      SUM=SUM+CC(M,J)
C      3 CONTINUE
C
C      CC(M,J)=SUM
C
C      2 CONTINUE
C      1 CONTINUE

```



```

C
RETURN
END

SUBROUTINE PLOT2D(YDISP,XTIME,LDOF,LTIME)
C
C THIS ROUTINE PLOTS RESULTS FROM COMPUTER OUTPUT
C IMPLICIT REAL(A-H,O-Z)
C DIMENSION XTIME(1000),YDISP(1000,100)
C
C SET THE Y (DISPLACEMENT) MINIMUM AND MAXIMUM AS
C -.003 AND +.003 RESPECTIVELY
C
CALL AGSETF('Y/MINIMUM.', -.003)
CALL AGSETF('Y/MAXIMUM.', .003)
C
C SET UP THE LEFT LABEL
C CALL AGSETF('LABEL/NAME.', 'L')
C CALL AGSETI('LINE/NUMBER.', 100)
C CALL AGSETP('LINE/TEXT.', 'NODAL DISPLACEMENTS', 1)
C
C SET UP BOTTOM LABEL
C CALL AGSETF('LABEL/NAME.', 'B')
C CALL AGSETI('LINE/NUMBER.', -100)
C CALL AGSETP('LINE/TEXT.', 'TIMES', 1)
C
C
C SET UP LABELS
C CALL ANOTAT('TIME(S)$', 'NODAL DISPLACEMENT(FT)$',
C 0,0,0,0)
C
C
C DRAW BOUNDARY ARUOND THE EDGE OF THE PLOTTER FRAME
C
CALL BNDARY
C
C DRAW THE GRAPH, USING EZMXY
C
CALL EZMXY(XTIME,YDISP,LTIME,LDOF,LTIME,
'DISP VRS. TIME (EXPLICIT)$')
C
RETURN
END

SUBROUTINE EXPLCT(XMAS,STIF,NDOF,CONST,NTIME,WRK1,
, WRK2,WRK3,WRK4,WRK5,WRK8,WRK10,WRK11,WRK12,
, WRK13,Z1,Z2,Z3,Z4,Z4I,Z5,Z6,Z66,
, A0,A1,A2,FORCE,EFMAS)
C
C*****

```



```

C
  OMSQ=0.
  DO 13 I=1,NDOF
13 OMSQ=DMAX1(EIGV(I),OMSQ)
  OMEG1=DSQRT(OMSQ)
C
  PRIOD=2.*PI/OMEG1
  DELT=PRIOD*CONST
C
C
C*****
C
C      COMPUTE INTEGRATION PARAMETERS
C
  TAU0=1./(DELT**2.)
  TAU1=1./(2.*DELT)
  TAU2=2.*TAU0
  TAU3=1./TAU2
C
C
C*****
C
C      COMPUTE ALPHA, BETA, GAMMA CONSTANTS FROM Q - MODEL
C
  CALL QMODEL(XB)
C
  ALPHA=XB(1,1)
  BETA=XB(2,1)
  GAMMA=XB(3,1)
C
C
C      COMPUTE DAMPING MATRICES [A0], [A1], [A2]
C
  DO 6 II=1,NDOF
  DO 6 JJ=1,NDOF
  A0(II,JJ)=ALPHA*STIF(II,JJ)
  A1(II,JJ)=BETA*STIF(II,JJ)
  A2(II,JJ)=GAMMA*STIF(II,JJ)
  6 CONTINUE
C
C
C*****
C
C      COMPUTE INITIAL VALUES OF VARIABLES
C
  DO 1 I=1,NDOF
  Z1(I)=0.
  Z2(I)=0.
  Z3(I)=0.
  Z4(I)=0.
  Z5(I)=0.
  1 CONTINUE
C
C
C      TRIANGULARIZE MASS MATRIX
C
  CALL DECOMP(XMAS,WRK10,NDOF,NDOF,1)

```



```

      8 EFMAS(I,J)=TAU0*XMAS(I,J)+TAU1*A0(I,J)
C
C
C      TRIANGULARIZE EFFECTIVE MASS MATRIX
C
C      CALL DECOMP(EFMAS,WRK13,NDOF,NDOF,1)
C
C      CALCULATE DISPLACEMENT VECTOR AT TIME T+DELT
C
C      CALL SOLVEQ(WRK13,FORCE,WRK12,NDOF,NDOF,1,1,1,1)
C
C      DO 15 I=1,NDOF
15 UD(I)=WRK12(I,1)
C
C
C      WRITE(7,2000)
C      WRITE(7,2001)(UD(I),I=1,NDOF)
C
C
C      REARRANGE DATA FOR PLOTTING
C
C      ZTIME(JTIME)=TIME
C      DO 11 IJ=1,4
C      JJ=IJ+6
11 ZDISP(JTIME,IJ)=Z4(JJ)
C
C
C      INCREMENT TIME AND UPDATE VARIABLES
C
C      TIME=TIME+DELT
C
C
C      DO 10 I=1,NDOF
C      Z1(I)=Z3(I)/TAU2+Z1(I)+DELT*Z2(I)
C      Z2(I)=Z4(I)/TAU2+Z2(I)+DELT*Z3(I)
C      Z3(I)=Z5(I)/TAU2+Z3(I)+DELT*Z4(I)
C      Z4I(I)=Z4(I)
C      Z4(I)=UD(I)
10 CONTINUE
C
C
C      4 CONTINUE
C
C*****
C
C      PLOT THE RESULTS
C
C      CALL PLOT2D(ZDISP,ZTIME,4,NTIME)
C
C*****
C
C      2000 FORMAT(//,6X,'DISPLACEMENT')

```



```

        TAU3=(1./2.*XLAMD)-1.
        TAU4=DEL/XLAMD-1.
        TAU5=DEL*(DEL/XLAMD-2.)/2.
        TAU6=DEL*(1.-DEL)
        TAU7=DEL*DEL
C
C*****
C
C      COMPUTE ALPHA, BETA, GAMMA CONSTANTS FROM Q - MODEL
C
C      CALL QMODEL(XB)
C
C      ALPHA=XB(1,1)
C      BETA=XB(2,1)
C      GAMMA=XB(3,1)
C
C
C      COMPUTE DAMPING MATRICES [A0],[A1],[A2]
C
C      DO 3 II=1,NDOF
C      DO 3 JJ=1,NDOF
C      A0(II,JJ)=ALPHA*STIF(II,JJ)
C      A1(II,JJ)=BETA*STIF(II,JJ)
C      A2(II,JJ)=GAMMA*STIF(II,JJ)
C 3 CONTINUE
C
C
C*****
C
C      COMPUTE INITIAL VALUES OF VARIABLES
C
C
C      DO 1 I=1,NDOF
C      Z1(I)=0.
C      Z2(I)=0.
C      Z3(I)=0.
C      Z4(I)=0.
C      Z5(I)=0.
C 1 CONTINUE
C
C
C
C      TRIANGULARIZE THE MASS MATRIX
C      CALL DECOMP(XMAS,WRK10,NDOF,NDOF,1)
C
C      COMPUTE INITIAL ACCELERATION AS FOLLOWS
C      -1
C      {Z6(0)}=[M]-1{F(0)-[K]{Z4(0)}-[A0]{Z5(0)}} (SINCE
C      {A1}{Z3(0)},[A2]{Z1(0)} ARE IDENTICALLY EQUAL TO ZERO
C
C      FIRST COMPUTE [K]{Z4(0)}
C      CALL VMULT(STIF,Z4,WRK9,NDOF,NDOF,NDOF)
C
C
C      NEXT COMPUTE [A0]{Z5(0)}
C      CALL VMULT(A0,Z5,Z66,NDOF,NDOF,NDOF)
C
C
C      DO 2 I=1,NDOF
C      WRK8(I,1)=FFCN(TIME,I)-WRK9(I)-Z66(I)
C 2 CONTINUE
C
C      COMPUTE INITIAL VALUE OF Z6

```

```

      CALL SOLVEQ(WRK10,WRK8,WRK11,NDOF,NDOF,1,1,1,1)
      DO 102 I=1,NDOF
102  Z6(I)=WRK11(I,1)
C*****
C
C
C
C      COMPUTE EFFECTIVE STIFFNESS MATRIX
C
C      DO 5 I=1,NDOF
C      DO 5 J=1,NDOF
C      5  EFFK(I,J)=STIF(I,J)+TAU0*XMAS(I,J)+TAU1*A0(I,J)
C
C
C
C      TRIANGULARIZE THE EFFECTIVE STIFFNESS MATRIX
C      AND STORE IN WRK1
C
C      CALL DECOMP(EFFK,WRK1,NDOF,NDOF,1)
C*****
C
C      TIME ITERATION BEGINS
C
C      DO 4 JTIME=1,NTIME
C
C
C      CALCULATE EFFECTIVE LOADS AT TIME T+DELT
C
C      DO 7 II=1,NDOF
C      7  WRK2(II)=TAU0*Z4(II)+TAU2*Z5(II)+TAU3*Z6(II)
C
C
C      COMPUTE {XMAS}{WRK2}
C      DO 104 MM=1,NDOF
C      SUM1=0.
C      DO 105 JJ=1,NDOF
C      WRK3(MM)=XMAS(MM,JJ)*WRK2(JJ)
C      SUM1=SUM1+WRK3(MM)
C      105 CONTINUE
C      WRK3(MM)=SUM1
C      104 CONTINUE
C
C      DO 8 II=1,NDOF
C      8  WRK4(II)=TAU1*Z4(II)+TAU4*Z5(II)+TAU5*Z6(II)
C
C      CALL VMULT(A0,WRK4,WRK5,NDOF,NDOF,NDOF)
C
C
C
C      COMPUTE Z1(T+DELT), Z2(T+DELT), Z3(T+DELT) USING
C      TAYLOR'S EXPANSION
C
C      DO 9 II=1,NDOF
C      Z1(II)=Z1(II)+Z2(II)*DELT+Z3(II)*0.5*DELT**2.
C      Z2(II)=Z2(II)+Z3(II)*DELT+Z4(II)*0.5*DELT**2.
C      Z3(II)=Z3(II)+Z4(II)*DELT+Z5(II)*0.5*DELT**2.
C      9  CONTINUE

```



```

SUBROUTINE STIF3D(DMTRX,NX,NY,NZ,NGS,X,Y,Z,
,   NDIM,NDIM3,DRVTS,FNC,DNRM,BSAVE,WRK,NLST,
,   XX,YY,ZZ,AJAC,STIF1)
C
C
C          *****
C          *
C          *   P R O G R A M   S T I F 3 D   *
C          *
C          *****
C
-----
C
C  A IN-CORE DOUBLE PRECISION PROGRAM TO CALCULATE THE ELEMENT STIFFNESS
C  MATRIX OF A 3-D COMPLEX ELEMENT OF ARBITRARY BRICK SHAPE WITH
C  BOUNDARY NODES ONLY (i.e. SERENDIPITY ELEMENTS); THE NUMBER OF NODES
C  ALONG EACH DIRECTION CAN BE GIVEN AS AN INPUT.
C
C  THE CONCEPT OF OBTAINING THE STIFFNESS MATRIX IS :
C  1. GENERATE INTERNAL NODAL POINTS FROM THE GIVEN BOUNDARY NODES
C     SUCH THAT IT FORMS A LAGRANGIAN ELEMENT
C  2. CALCULATE ALL OF THE SHAPE FUNCTIONS CORRESPONDING L. ELEM.
C  3. FORM STIFFNESS MATRIX OF THE L. ELEMENT BY USING GAUSS-
C     LEGENDRE QUADRATURE INTEGRATION SCHEME.
C  4. CONDENSE OUT ALL OF THE INTERNAL NODES; THE STIFFNESS MATRIX
C     IS OBTAINED.
C
-----
C
C  VARIABLE NAME LISTING :
C  *****
C
C  STIF1 ..... LOCAL STIFFNESS MATRIX WITH DIMENSION 3*NDIM x 3*NDIM
C  DRVTS ..... MATRIX STORAGE WHICH CONTAINS THE DERIVATIVES OF THE
C               NONLINEAR INTERPOLATION FUNCTION AT EACH GAUSS POINT.
C               (i.e. IT HAS DIMENSION NDIM x 3 x 10;
C                 3 --- INDICATES DERIVATIVES AT EACH DIRECTION;
C                 10--- INDICATES THE NUMBER OF GAUSS POINTS.
C  DMTRX ..... A 6 x 6 MATRIX CONTAINS THE MATERIAL CONSTANT MATRIX.
C  FNC ..... MATRIX STORAGE WHICH CONTAINS THE VALUES OF THE
C               NONLINEAR INTERPOLATION FUNCTIONS OBTAINED AT EACH
C               GAUSS POINT. IT HAS SAME DIMENSION AS DRVTS.
C  DNRM ..... MATRIX STORAGE WITH DIMENSION NDIM x 3, WHICH CONTAINS
C               THE VALUES OF THE NORMALIZATION FACTORS.
C  X,Y,Z ..... ARRAYS OF LENGTH NDIM, WHICH CONTAIN THE X, Y, AND Z
C               COORDINATES OF NODES IN GLOBAL COORDINATES SYSTEM;
C               WHERE NDIM(TOTAL NUMBER OF NODES)=NX*NY*NZ.
C  XX,YY,ZZ ..... ARRAYS OF LENGTH NDIM, WHICH CONATINS A SET OF LOCAL
C               COORDINATES.
C  BSAVE ..... STORAGE OF CALCULATED DERIVATIVES OF SHAPF FUNCTIONS
C               WHICH ARE USED TO FORM THE [B]-MATRIX.
C  WRK ..... A 3*NDIM x 3*NDIM WORKING ARRAY FOR WORK'NG PACE.
C  WGHT ..... VECTOR OF LENGTH 10 CONTAINS THE WEIGHTS FOR EACH
C               GAUSS POINT.
C  NGS ..... ACTUAL REQUIRED GAUSS POINTS.
C  NDIM ..... TOTAL NUMBER OF NODES
C  NDIM3 ..... INITIAL DIMENSION OF STIFFNESS MATRIX.(3*NDIM)

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C      NLST ..... A LOCAL NODAL NUMBER LIST FOR REDUCING THE LAGRANGIAN
C      ELEMENT STIFFNESS MATRIX TO SEREDIPITY ONE PROPERLY.
C      IRED ..... INDICATOR IF THE CONDENCING PROCEDURE IS REQUIRED;
C      IRED = 0 ... NO CONDENCE; IRED = 1 ... CONDENCE REQ.
C      NX,NY,NZ ..... NUMBER OF NODES ALONG X, Y, AND Z DIRECTION, RESPECT.
C      ITAPE ..... LOGICAL UNIT NUMBER ON WHICH THE STIFFNESS MATRIX OF
C      LAGRANGIAN ELEMENT IS SAVE IF IT IS NECESSERARY. IF
C      ITAPE = 0, NO SAVE IS DONE.
C-----
C
C
C      IMPLICIT REAL*8 (A-H,O-Z)
C      DIMENSION STIF1(NDIM3,NDIM3),DRVTS(NDIM,3,10),DMTRX(6,6),
+      FNC(NDIM,3,10),DNRM(NDIM,3),BSAVE(NDIM,3),X(NDIM),
+      Y(NDIM),Z(NDIM),WRK(NDIM3,NDIM3),NLST(NDIM),WGHT(10),
+      XX(NDIM),YY(NDIM),ZZ(NDIM)
C
C
C      OBTAIN A LOCAL NODAL NUMBER LISTING --- BOUNDARY NODES IN THE FRONT
C      FOLLOWED BY INTER NODES
C
C      CALL LCNODE(NLST,NX,NY,NZ,NDIM)
C
C      OBTAIN THE SHAPE FUNCTIONS
C
C      CALL SHAPE1(XX,YY,ZZ,DRVTS,FNC,DNRM,NDIM,NX,NY,NZ,NGS,WRK,NDIM3,0)
C
C      OBTAIN THE LOCAL STIFFNESS MATRIX
C
C      CALL LCSTIF(STIF1,DRVTS,DMTRX,FNC,DNRM,BSAVE,X,Y,Z,WRK,NGS,NDIM,
+      NDIM3,NLST,0,NX,NY,NZ,0,AJAC)
C
C
C      RETURN
C      END

C*****
C      SUBROUTINE BMTRX1(X,Y,Z,XX,YY,ZZ,BMTRX,NDIM,NX,NY,NZ,NGS,WRK,
+      VWRK,AJAC,IGEN)
C*****
C
C      CALCULATE THE B-MATRIX BY USING GAUSS-LEGENDRE QUADRATURE FORMULA.
C
C      X ..... ARRAY OF GLOBAL NODAL COORDINATES IN X.
C      Y ..... ARRAY OF GLOBAL NODAL COORDINATES IN Y.
C      Z ..... ARRAY OF GLOBAL NODAL COORDINATES IN Z.
C      XX ..... ARRAY OF NORMAL (natural) NODAL COORDINATES IN XX.
C      YY ..... ARRAY OF NORMAL (natural) NODAL COORDINATES IN YY.
C      ZZ ..... ARRAY OF NORMAL (natural) NODAL COORDINATES IN ZZ.
C      BMTRX ..... ARRAY OF DIMENSION (3*NPE x 3 x NGS), WHICH CONTAINS THE
C      EVALUATED B-MATRIX AT EACH GAUSS POINTS; Where
C      (NPE --- number of nodes per element)
C      NDIM ..... FIRST ROW DIMENSION OF X,Y,Z,XX,YY,ZZ AND BMTRX IN
C      CALLING PROGRAM.
C      NX ..... NUMBER OF NODES ALONG X.
C      NY ..... NUMBER OF NODES ALONG Y.
C      NZ ..... NUMBER OF NODES ALONG Z.
C      NGS ..... NUMBER OF GAUSS POINTS REQUIRED IN NUMERICAL INTEGRATION.

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C   WRK ..... WORKING SPACE OF LENGTH AT LEAST 3*NDIM.
C   VWRK ..... WORKING VECTOR OF LENGTH AT LEAST AS SAME AS X, Y, AND Z.
C   AJAC ..... ARRAY OF LENGTH NGS WITH JACOBIAN VALUES AT EACH GAUSS
C               POINTS. (i.e. determinant of JACOBIAN matrix)
C   IGEN ..... NODAL POINT COORDINATE GENERATIONS INDICATOR
C               IGEN = 0   NORMAL COORDINATES ARE GENERATED;
C               IGEN = 1   NORMAL COORDINATES ARE FROM INPUT.
C .....
C
C   SUBROUTINES CALLED :
C
C               DERVTS ---- CALCULATE DERIVATIVES.
C               JACOBN ---- CALCULATE INVERSE OF JACOBIAN MATRIX, AND
C                           VALUE OF JACOBIAN.
C               NCOORD ---- GENERATE A SET NORMAL NODAL COORDINATES.
C               NUMINT ---- PERFORM THE NUMERICAL INTEGRATIONS.
C -----
C
C   IMPLICIT REAL*8 (A-H,O-Z)
C   DIMENSION X(NDIM),Y(NDIM),Z(NDIM),XX(NDIM),YY(NDIM),ZZ(NDIM),
C   +         BMTRX(NDIM,3,10),VWRK(NDIM),WRK(1),AJAC(10),
C   +         DVX(3,10),DVY(3,10),DVZ(3,10),DFUNC(3),AJACB(3,3),
C   +         XINT1(10),XINT2(10),XINT3(10),DVXX(10),DVYY(10),DVZZ(10)
C
C
C   NPE=8+4*(NX+NY+NZ-6)
C   IWRK1=0
C   IWRK2=IWRK1+NDIM
C   IWRK3=IWRK2+NDIM
C
C   INITIALIZE XX, YY, ZZ, VWRK, WRK AND BMTRX.
C
C   DO 2 I=1,NDIM
C   DO 1 J=1,3
C   DO 1 K=1,NGS
C   BMTRX(I,J,K)=0.0
C 1  CONTINUE
C   IK1=IWRK1+I
C   IK2=IWRK2+I
C   IK3=IWRK3+I
C   IF(IGEN.EQ.0) THEN
C   XX(I)=0.0
C   YY(I)=0.0
C   ZZ(I)=0.0
C   ENDIF
C   WRK(IK1)=X(I)
C   WRK(IK2)=Y(I)
C   WRK(IK3)=Z(I)
C   VWRK(I)=0.0
C 2  CONTINUE
C   DO 3 I=1,3
C   DO 3 J=1,NGS
C   DVX(I,J)=0.0
C   DVY(I,J)=0.0
C   DVZ(I,J)=0.0
C 3  CONTINUE
C
C   GENERATE NODAL COORDINATES IN NORMAL (NATURAL) COORDINATES
C   IF IT IS REQUIRED.

```



```

C
C     IF(IGEN.EQ.0) CALL NCOORD(XX,YY,ZZ,NX,NY,NZ,NDIM)
C
C CALCULATE THE NORMALIZED INTERPOLATION FUNCTIONS (I.F.) AND THE
C CORRESPONDING DERIVATIVES
C     dF/dXX=Sum(dF(i)/dXX), i=1,NPE; etc....
C     where
C         F(i)=Fi(XX)*Fi(YY)*Fi(ZZ),           i.e.
C         dF(i)/dXX=Fi(YY)*Fi(ZZ)*(dFi(XX)/dXX) etc....
C
C     DO 4 IPE=1,NPE                !Evaluations node by node
C         TRAS1=1.0
C         TRAS2=1.0
C         TRAS3=1.0
C         IPTX=0
C         IPTY=0
C         IPTZ=0
C         DO IXYZ=1,NPE
C             X(IXYZ)=0.
C             Y(IXYZ)=0.
C             Z(IXYZ)=0.
C         END DO
C         DO 5 INOD=1,NPE
C             IF(INOD.EQ.IPE) GO TO 5        !Only choose i .ne. j the term
C
C EVALUATE NORMALIZATION FACTORS AND CHOOSE APPROPRIATE NODAL POINTS.
C
C     IF(YY(INOD).EQ.YY(IPE).AND.ZZ(INOD).EQ.ZZ(IPE)) THEN
C         IPTX=IPTX+1                !Interpolation function along XX
C         TERM1=XX(IPE)-XX(INOD)     !Normalization factor in XX
C         TRAS1=TRAS1*TERM1
C         X(IPTX)=XX(INOD)          !XX-points of interpolation func
C     ENDIF
C     IF(XX(INOD).EQ.XX(IPE).AND.ZZ(INOD).EQ.ZZ(IPE)) THEN
C         IPTY=IPTY+1                !Interpolation function along YY
C         TERM2=YY(IPE)-YY(INOD)     !Normalization factor in YY
C         TRAS2=TRAS2*TERM2
C         Y(IPTY)=YY(INOD)          !YY-points of interpolation func
C     ENDIF
C     IF(XX(INOD).EQ.XX(IPE).AND.YY(INOD).EQ.YY(IPE)) THEN
C         IPTZ=IPTZ+1                !Interpolation function along ZZ
C         TERM3=ZZ(IPE)-ZZ(INOD)     !Normalization factor in ZZ
C         TRAS3=TRAS3*TERM3
C         Z(IPTZ)=ZZ(INOD)          !ZZ-points of interpolation func
C     ENDIF
C 5     CONTINUE
C
C PERFORM NUMERICAL INTEGRATIONS WITHOUT WEIGHTS (IWG=0)
C
C     CALL NUMINT(X,NDIM,IPTX,NGS,XINT1,0)
C     CALL NUMINT(Y,NDIM,IPTY,NGS,XINT2,0)
C     CALL NUMINT(Z,NDIM,IPTZ,NGS,XINT3,0)
C     DNRM=TRAS1*TRAS2*TRAS3
C
C CALCULATE DERIVATIVES ABOUT X, Y, AND Z.
C
C     CALL DERVTS(DVXX,X,NDIM,IPTX,VWRK,NGS)
C     CALL DERVTS(DVYY,Y,NDIM,IPTY,VWRK,NGS)
C     CALL DERVTS(DVZZ,Z,NDIM,IPTZ,VWRK,NGS)
C

```

```

C Saving the evaluated values
C
      IK1=IWRK1+IPE
      IK2=IWRK2+IPE
      IK3=IWRK3+IPE
      DO 10 IGS=1,NGS
      XINT23=XINT2(IGS)*XINT3(IGS)
      XINT13=XINT1(IGS)*XINT3(IGS)
      XINT12=XINT1(IGS)*XINT2(IGS)
      XXDV=DVXX(IGS)*XINT23/DNRM
      YYDV=DVYY(IGS)*XINT13/DNRM
      ZZDV=DVZZ(IGS)*XINT12/DNRM
      BMTRX(IPE,1,IGS)=XXDV
      BMTRX(IPE,2,IGS)=YYDV
      BMTRX(IPE,3,IGS)=ZZDV
C Calculate dx(XX,YY,ZZ)/dXX=Sum(X(i)*dF(i)/dXX), i=1,NPE      etc.....
C for each GAUSS point
C
      DVX(1,IGS)=DVX(1,IGS)+WRK(IK1)*XXDV
      DVX(2,IGS)=DVX(2,IGS)+WRK(IK1)*YYDV
      DVX(3,IGS)=DVX(3,IGS)+WRK(IK1)*ZZDV
      DVY(1,IGS)=DVY(1,IGS)+WRK(IK2)*XXDV
      DVY(2,IGS)=DVY(2,IGS)+WRK(IK2)*YYDV
      DVY(3,IGS)=DVY(3,IGS)+WRK(IK2)*ZZDV
      DVZ(1,IGS)=DVZ(1,IGS)+WRK(IK3)*XXDV
      DVZ(2,IGS)=DVZ(2,IGS)+WRK(IK3)*YYDV
      DVZ(3,IGS)=DVZ(3,IGS)+WRK(IK3)*ZZDV
10 CONTINUE
4 CONTINUE
C
C Evaluate the inverse of JACOBIAN-Matrix and determinant of Jacobian
C at each GAUSS points.
C
      NXYZ=6
      DO 11 IGS=1,NGS
      DO 12 IJB=1,3
      DVXX(IJB)=DVX(IJB,IGS)
      DVYY(IJB)=DVY(IJB,IGS)
      DVZZ(IJB)=DVZ(IJB,IGS)
12 CONTINUE
      CALL JACOBN(DVXX,DVYY,DVZZ,AJAC1,AJACB,NXYZ)
      AJAC(IGS)=AJAC1
C
C Form B-matrix at each GAUSS point in terms of X, Y, and Z values,
C i.e. {dF(i)/dX(i)}=[AJACB]*{dF(i)/dXX(i)}, where
C X(i) ..... Global coordinates;
C XX(i) ..... Natural coordinates;
C i ..... the nodal numbers
C
      DO 6 IPE=1,NPE      !Evaluate it node by node.
      DO 7 IDR=1,3
      DFUNC(IDR)=0.0
      DO 8 IDC=1,3
      DFUNC(IDR)=DFUNC(IDR)+AJACB(IDR,IDC)*BMTRX(IPE,IDC,IGS)
8 CONTINUE
7 CONTINUE
      BMTRX(IPE,1,IGS)=DFUNC(1)
      BMTRX(IPE,2,IGS)=DFUNC(2)
      BMTRX(IPE,3,IGS)=DFUNC(3)
6 CONTINUE

```

```

11 CONTINUE
C
C
C RETURN
C END
C
C *****
C SUBROUTINE DECOMP(SAVE,A,NROW,NDIM,ISYM)
C *****
C A DOUBLE PRECISION CODE WHICH PERFORMS THE LU-DECOMPOSITION OF THE
C SQUARE MATRIX [A]; [A]=[L]*[U]. IF [A] IS SYMTRIC MATRIX (i.e.ISYM=1)
C THEN THE SYMTRIC MATRIX DECOMPOSITION IS PERFORMED; [A]=Trans(L)*[L].
C -----
C A ..... DECOMPOSED MATRIX A=LU.
C SAVE.....INPUT MATRIX TO BE DECOMPOSED. THIS MATRIX REMAINS INTACT
C UPON RETURN
C NROW ..... ROW DIMENSION OF MATRIX [A].
C NDIM ..... INITIAL DIMENSION OF MATRIX [A].
C ISYM ..... SYMETRIC INDICATOR; ISYM=0 ... NONSYM. BUT POSITIVE.
C ISYM=1 ... SYMMET. AND POSITIVE.
C ISYM=3 ... NONSYM. AND NONPOSIT.
C -----
C IMPLICIT REAL*8 (A-H,O-Z)
C DIMENSION A(NDIM,NDIM),SAVE(NDIM,NDIM)
C
C SAVE ORIGINAL MATRIX BEFORE DECOMPOSITION
C
C DO 6 I=1,NROW
C DO 6 J=1,NROW
C 6 A(I,J)=SAVE(I,J)
C
C START DECOMPOSITION.
C
C DO 1 IELE=1,NROW
C
C DETERMINE ELEMENT OF [L].
C
C DO 2 IROW=IELE,NROW
C CSUM=0.
C IF(IELE.EQ.1) GO TO 997
C DO 3 ISUM=1,IELE-1
C CSUM=CSUM+A(IROW,ISUM)*A(ISUM,IELE)
C 3 CONTINUE
C 997 A(IROW,IELE)=A(IROW,IELE)-CSUM
C AA=A(IELE,IELE)
C IF(ISYM.LT.3.AND.AA.LT.1.D-10) GO TO 998 !CHECK POSIT. AND SYM.
C IF(DABS(AA).LT.1.D-10) GO TO 998 !CHECK IF PIVOT = 0.
C IF(ISYM.EQ.1) THEN
C IF(IROW.EQ.IELE) THEN
C A(IELE,IELE)=DSQRT(AA)
C ELSE
C A(IROW,IELE)=A(IROW,IELE)/A(IELE,IELE)
C ENDIF
C ENDIF
C 2 CONTINUE

```

```

C
C DETERMINE ELEMENTS OF [U]
C
      IF(IELE.EQ.NROW) GO TO 1
      DO 4 JCOL=IELE+1,NROW
      RSUM=0.0
      IF(ISYM.EQ.1) GO TO 995
      IF(IELE.EQ.1) GO TO 996
      DO 5 ISUM=1,IELE-1
      RSUM=RSUM+A(IELE,ISUM)*A(ISUM,JCOL)
5     CONTINUE
996   A(IELE,JCOL)=(A(IELE,JCOL)-RSUM)/A(IELE,IELE)
      GO TO 4
995   A(IELE,JCOL)=A(JCOL,IELE)
4     CONTINUE
1     CONTINUE
C
      RETURN
998   CONTINUE
      IF(AA.LT.0.) THEN
      PRINT 1000, IROW, AA
      ELSE
      PRINT 1001, IROW
      ENDIF
1000  FORMAT(/,' **** ERROR : MATRIX IS NON-POSITIVE DEFINITE ...',/,
+         '      EQUATION NUMBER ', I10,/,
+         '      PIVOTING VALUE =', D10.4,/)
1001  FORMAT(/,' **** ERROR : MATRIX IS NEARLY SINGULAR ...',/,
+         '      EQUATION NUMBER ', I10,/)
      STOP
      END
C
C*****
C      SUBROUTINE DERVTS(DVXI,XX,NDIM,NPT,WRK,NGS,NWRK,ICOL)
C*****
C      CALCULATE THE DERIVATIVE (Numerically) OF INTERPOLATION FUNCTIONS X
C      BY USING GAUSS-LEGENDRE QUADRATURE FORMULA.
C
C      DVXI ..... ARRAY OF LENGTH NGS, THE DERIVATIVES OF THE
C                  INTERPOLATION FUNCTION AT EACH GAUSS POINTS.
C      XX ..... ARRAY OF INITIAL LENGTH NDIM, CONTAINS THE NODAL
C                COORDINATES IN THE INTERPOLATION FUNCTION (i.e. In
C                normal coordinate system).
C      NDIM ..... INITIAL DIMENSION OF ARRAY X.
C      NPT ..... ACTUAL LENGTH OF ARRAY X.
C      WRK ..... WORKING SPACE WITH DIMENSION NWRK X NWRK.
C      NGS ..... NUMBER OF GAUSS POINTS QUIRED.
C      NWRK ..... INITIAL DIMENSION OF MATRIX WRK.
C      ICOL ..... I-TH COLUMN, WHICH WILL BE USED AS A WORKING SPACE.
C
C      SUBROUTINES CALLED :
C                  NUMINT ----- A NUMERICAL INTERATION ROUTINE.
C
C-----
C
C      IMPLICIT REAL*8 (A-H,O-Z)
C      DIMENSION XX(NDIM),WRK(NWRK,NWRK),DVXI(10),DVXJ(10)
C
C      INITINALIZE WRK

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```

DEC00046
DEC00047
DEC00048
DEC00049
DEC00050
DEC00051
DEC00052
DEC00053
DEC00054
DEC00055
DEC00056
DEC00057
DEC00058
DEC00059
DEC00060
DEC00061
DEC00062
DEC00063
DEC00064
DEC00065
DEC00066
DEC00067
DEC00068
DEC00069
DEC00070
DEC00071
DEC00072
DEC00073
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C
DO 1 I=1,NPT
WRK(I,ICOL)=XX(I)
1 CONTINUE
DO 2 I=1,10
DVXI(I)=0.0
2 CONTINUE
!Initialize dFi(XX)/dXX.
C
Fi(XX)=(XX-XX(1))*(XX-XX(2))*.....*(XX-XX(i-1))*(XX-XX(i+1))*....
hence :
dFi(XX)/dXX=Sum(dFij(XX)/dXX),jj=1,NPT;
there
dFij(XX)/dXX=(XX-XX(1))*(XX-XX(2))*.....*(XX-XX(j-1))*(XX-XX(j+1))*
.....*(XX-XX(i-1))*(XX-XX(i+1))*.....
C
DO 3 ITERM=1,NPT
IPT=0
DO 4 IX=1,NPT
IF(IX.EQ.ITERM) GO TO 4
IPT=IPT+1
XX(IPT)=WRK(IX,ICOL)
4 CONTINUE
CALL NUMINT(XX,NDIM,IPT,NGS,DVXJ,0)
DO 5 IG=1,NGS
DVXI(IG)=DVXI(IG)+DVXJ(IG)
5 CONTINUE
3 CONTINUE
!Calculate dFij(XX)/dXX.
!Choose proper XX-points.
!Skip XX(i) and XX(j).
!Sum of dFij(XX)/dXX in j.
C
RETURN
END
C
C*****
SUBROUTINE FORMVK(VWRK,VWK)
C*****
FORM AN APPROPRIATE B-MATRIX FOR INTERPOLATION FUNCTION # INOD.
VWRK ..... A SUB-MATRIX OF B-MATRIX FOR NODAL NUMBER (INOD) WITH
DIMENSION 6 x 3.
VWK ..... AN ARRAY OF LENGTH 3, WHICH CONTAINS THE DERIVATIVES OF
THE INTERPOLATION FUNCTIONS WITH RESPECT TO X, Y, AND Z
IN GLOBAL COORDINATE SYSTEM.
-----
IMPLICIT REAL*8 (A-H,O-Z)
DIMENSION VWRK(6,3),VWK(3)
INITIALIZE MATRIX VWRK.
DO 1 IVK=1,6
DO 1 JVK=1,3
1 VWRK(IVK,JVK)=0.0
FORM VWRK SUB-MATRIX.
VWRK(1,1)=VWK(1)
VWRK(2,2)=VWK(2)
VWRK(3,3)=VWK(3)
VWRK(4,1)=VWRK(2,2)

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VWRK(4,2)=VWRK(1,1) FOR00030
VWRK(5,1)=VWRK(3,3) FOR00031
VWRK(5,3)=VWRK(1,1) FOR00032
VWRK(6,2)=VWRK(3,3) FOR00033
VWRK(6,3)=VWRK(2,2) FOR00034
C FOR00035
C FINISH FOR00036
C FOR00037
C RETURN FOR00038
C END FOR00039
C GCO00001
C*****GCO00002
C SUBROUTINE GCOORD(X,Y,Z,NX,NY,NZ,NDIM,NLST) GCO00003
C*****GCO00004
C GCO00005
C GENERATE A SET OF NODAL COORDINATES, WHICH ARE NOT ON THE ELEMENT GCO00006
C EDGES, IN GLOBAL COORDINATES SYSTEM IN ORDER TO FORM A LAGRANGIAN GCO00007
C ELEMENTS FROM THE GIVEN SERENDIPITY ELEMENTS. GCO00008
C GCO00009
C-----GCO00010
C X,Y,Z ..... NODAL COORDINATES IN X, Y, AND Z-DIRECTIONS. AS INPUT GCO00011
C THEY MUST CONTAIN THE COORDINATES OF THE NODES WHICH GCO00012
C ARE ON THE EDGES IN PROPER SEQUENCE. GCO00013
C NX,NY,NZ ..... NUMBER OF NODES IN X, Y, AND Z-DIRECTIONS. GCO00014
C NDIM ..... INITIAL DIMENSION OF X, Y, AND Z VECTORS. GCO00015
C NLST ..... LOCAL NODAL NUMBER LISTING. GCO00016
C-----GCO00017
C GCO00018
C IMPLICIT REAL*8 (A-H,O-Z) GCO00019
C DIMENSION X(NDIM),Y(NDIM),Z(NDIM),NLST(NDIM) GCO00020
C GCO00021
C NPE1=8+4*(NX+NY+NZ-6) !NUMBER OF ELEMENTS ON EDGESGCO00022
C IPT1=NPE1+1 GCO00023
C NNX=NX GCO00024
C NNY=NY GCO00025
C NNZ=NZ GCO00026
C IF(NNX.LE.0) NNX=1 GCO00027
C IF(NNY.LE.0) NNY=1 GCO00028
C IF(NNZ.LE.0) NNZ=1 GCO00029
C GCO00030
C GENERATE NODES GCO00031
C GCO00032
C NZ1=2*(NX+NY-2) GCO00033
C NZ2=NZ1+4*(NZ-2) GCO00034
C RATIO=1./DFLOAT(NZ-1) GCO00035
C DO 1 IZ=1,NZ GCO00036
C IZ1=NZ1+4*(IZ-2) !STARTING POINT GCO00037
C IF(IZ.EQ.1) IZ1=NX GCO00038
C IF(IZ.EQ.NZ) IZ1=IZ1+NX GCO00039
C DX0=0. GCO00040
C DY0=0. GCO00041
C DZ0=0. GCO00042
C DX1=0. GCO00043
C DY1=0. GCO00044
C DZ1=0. GCO00045
C IF(IZ.GT.1.AND.IZ.LT.NZ) THEN GCO00046
C DX0=(X(IZ1+3)-X(IZ1+1))*RATIO GCO00047
C DX1=(X(IZ1+4)-X(IZ1+2))*RATIO GCO00048
C DY0=(Y(IZ1+3)-Y(IZ1+1))*RATIO GCO00049
C DY1=(Y(IZ1+4)-Y(IZ1+2))*RATIO GCO00050

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| | | |
|----|--------------------------------------------------------|----------|
| | DZ0=(Z(IZ1+3)-Z(IZ1+1))*RATIO | GCO00051 |
| | DZ1=(Z(IZ1+4)-Z(IZ1+2))*RATIO | GCO00052 |
| | ENDIF | GCO00053 |
| | ITY1=IZ1+1 | GCO00054 |
| | IEY1=ITY1+1 | GCO00055 |
| | DO 2 IY=1,NNY | GCO00056 |
| | IF((IZ.EQ.1.OR.IZ.EQ.NZ).AND.IY.EQ.NNY) GO TO 2 | GCO00057 |
| | IF((IZ.EQ.1.OR.IZ.EQ.NZ).AND.IY.EQ.1) GO TO 2 | GCO00058 |
| | IF(IY.EQ.1.OR.IY.EQ.NY) THEN | GCO00059 |
| | IF(IY.EQ.1) ITY=IZ1+1 | GCO00060 |
| | IF(IY.EQ.NY) ITY=IZ1+3 | GCO00061 |
| | IEY=ITY+1 | GCO00062 |
| | X0=X(ITY) | GCO00063 |
| | Y0=Y(ITY) | GCO00064 |
| | Z0=Z(ITY) | GCO00065 |
| | X1=X(IEY) | GCO00066 |
| | Y1=Y(IEY) | GCO00067 |
| | Z1=Z(IEY) | GCO00068 |
| | ELSE | GCO00069 |
| | X0=X(ITY1)+DFLOAT(IY-1)*DX0 | GCO00070 |
| | Y0=Y(ITY1)+DFLOAT(IY-1)*DY0 | GCO00071 |
| | Z0=Z(ITY1)+DFLOAT(IY-1)*DZ0 | GCO00072 |
| | X1=X(IEY1)+DFLOAT(IY-1)*DX1 | GCO00073 |
| | Y1=Y(IEY1)+DFLOAT(IY-1)*DY1 | GCO00074 |
| | Z1=Z(IEY1)+DFLOAT(IY-1)*DZ1 | GCO00075 |
| | IF(IZ.EQ.1.OR.IZ.EQ.NZ) THEN | GCO00076 |
| | ITY1=IEY1+1 | GCO00077 |
| | IEY1=ITY1+1 | GCO00078 |
| | ENDIF | GCO00079 |
| | ENDIF | GCO00080 |
| | X2=X0 | GCO00081 |
| | X3=X1 | GCO00082 |
| | DX=(X1-X0)/DFLOAT(NNX-1) | GCO00083 |
| | DY=(Y1-Y0)/DFLOAT(NNY-1) | GCO00084 |
| | DZ=(Z1-Z0)/DFLOAT(NNZ-1) | GCO00085 |
| | IF(IZ.EQ.1.OR.IZ.EQ.NZ) GO TO 89 | GCO00086 |
| | IF(IY.GT.1.AND.IY.LT.NY) THEN | GCO00087 |
| | X2=X2-DX | GCO00088 |
| | X3=X3+DX | GCO00089 |
| | ENDIF | GCO00090 |
| 89 | CONTINUE | GCO00091 |
| | DO 3 IX=1,NNX | GCO00092 |
| | DDX1=X0-X2 | GCO00093 |
| | DDX2=X0-X3 | GCO00094 |
| | IF(DABS(DDX1).LE.1.D-8.OR.DABS(DDX2).LE.1.D-8) GO TO 9 | GCO00095 |
| | X(IPT1)=X0 | GCO00096 |
| | Y(IPT1)=Y0 | GCO00097 |
| | Z(IPT1)=Z0 | GCO00098 |
| | IPT1=IPT1+1 | GCO00099 |
| 9 | X0=X0+DX | GCO00100 |
| | Y0=Y0+DY | GCO00101 |
| | Z0=Z0+DZ | GCO00102 |
| 3 | CONTINUE | GCO00103 |
| 2 | CONTINUE | GCO00104 |
| 1 | CONTINUE | GCO00105 |
| C | | GCO00106 |
| | RETURN | GCO00107 |
| | END | GCO00108 |
| C | | JAC0000 |
| C | ***** | JAC0000 |

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SUBROUTINE JACOBN(DX,DY,DZ,AJAC,AJACB,NXYZ) JAC00003
C***** JAC00004
C JAC00005
C PROGRAM TO EVALUATE THE JACOBIAN - MATRIX AND ITS DETERMINANT. JAC00006
C JAC00007
C DX ..... DERIVATIVES WITH RESPECT TO X, DIMENSION 3. JAC00008
C DY ..... DERIVATIVES WITH RESPECT TO Y, DIMENSION 3. JAC00009
C DZ ..... DERIVATIVES WITH RESPECT TO Z, DIMENSION 3. JAC00010
C AJAC ..... DETERMINANT OF JACOBIAN MATRIX JAC00011
C AJACB ..... THE INVERSE JACOBIAN MATRIX. JAC00012
C NXYZ ..... THE INITIAL DIMENSION OF ARRAYS DX, DY, AND DZ. JAC00013
C JAC00014
C-----JAC00015
C JAC00016
C IMPLICIT REAL*8 (A-H,O-Z) JAC00017
C DIMENSION DX(NXYZ),DY(NXYZ),DZ(NXYZ),AJACB(3,3) JAC00018
C JAC00019
C Evaluate the determinant of JACOBIAN - matrix. JAC00020
C JAC00021
C AJAC=DX(3)*(DY(1)*DZ(2)-DY(2)*DZ(1))-(DY(3)*(DX(1)*DZ(2)-
+ DX(2)*DZ(1))+DZ(3)*(DX(1)*DY(2)-DX(2)*DY(1))
C JAC00022
C JAC00023
C Evaluate the inverse of JACOBIAN - matrix. JAC00024
C JAC00025
C JAC00026
C AJACB(1,1)=(DY(2)*DZ(3)-DY(3)*DZ(2))/AJAC JAC00027
C AJACB(1,2)=(DY(3)*DZ(1)-DY(1)*DZ(3))/AJAC JAC00028
C AJACB(1,3)=(DY(1)*DZ(2)-DY(2)*DZ(1))/AJAC JAC00029
C AJACB(2,1)=(DX(3)*DZ(2)-DX(2)*DZ(3))/AJAC JAC00030
C AJACB(2,2)=(DX(1)*DZ(3)-DX(3)*DZ(1))/AJAC JAC00031
C AJACB(2,3)=(DX(2)*DZ(1)-DX(1)*DZ(2))/AJAC JAC00032
C AJACB(3,1)=(DX(2)*DY(3)-DX(3)*DY(2))/AJAC JAC00033
C AJACB(3,2)=(DX(3)*DY(1)-DX(1)*DY(3))/AJAC JAC00034
C AJACB(3,3)=(DX(1)*DY(2)-DX(2)*DY(1))/AJAC JAC00035
C JAC00036
C RETURN JAC00037
C END JAC00038
C LCN00001
C***** LCN00002
C SUBROUTINE LCNODE(NLST,NX,NY,NZ,NDIM) LCN00003
C***** LCN00004
C LCN00005
C THIS SUBROUTINE GENERATES A NODAL NUMBER LIST OF A LAGURANGIAN LCN00006
C ELEMENT SUCH THAT THE FIRST "NSE" NUMBERS ARE IN THE SEQUENCE OF A LCN00007
C SERENDIPITY ELEMENT AND FOLLOWED BY THE "OFF EDGE" NODAL NUMBERS. LCN00008
C LCN00009
C-----LCN00010
C NLST ..... ARRAY; AFTER RETURN IT CONTAINS THE RESULT. LCN00011
C NX,NY,NZ ..... NUMBER OF NODES ALONG X, Y, AND Z DIRECTION. LCN00012
C NDIM ..... INITIAL DIMENSION OF ARRAY NLST. LCN00013
C-----LCN00014
C LCN00015
C IMPLICIT REAL*8 (A-H,O-Z) LCN00016
C DIMENSION NLST(NDIM) LCN00017
C LCN00018
C NSE=8+4*(NX+NY+NZ-6) LCN00019
C NLE=NX*NY*NZ LCN00020
C LCN00021
C INITIALIZE NLST LCN00022
C LCN00023
C DO 1 I=1,NDIM LCN00024

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1      NLST(I)=0
C
C NUMBERING START
C
C 1. FACE AT IZ=1, AND IZ=NZ
C
      IZ1=0
      INOD1=1
      IZN=NSE-2*(NX+NY-2)
      INODZ=NLE-NX*NY+1
      ISE1=NSE
      ISEZ=INODZ+2*(NX+NY-2)-1
      DO 2 IY=1,NY
      DO 3 IX=1,NX
      IF((IY.GT.1.AND.IY.LT.NY).AND.(IX.GT.1.AND.IX.LT.NX)) GO TO 900
      IZ1=IZ1+1
      IZN=IZN+1
      NLST(IZ1)=INOD1
      NLST(IZN)=INODZ
      GO TO 901
900   ISE1=ISE1+1
      ISEZ=ISEZ+1
      NLST(ISE1)=INOD1
      NLST(ISEZ)=INODZ
901   INOD1=INOD1+1
      INODZ=INODZ+1
      3   CONTINUE
      2   CONTINUE
      IF(NZ.LE.2) RETURN
      DO 4 IZ=2,NZ-1
      DO 5 IY=1,NY
      DO 6 IX=1,NX
      IF(IY.GT.1.AND.IY.LT.NY) GO TO 902
      IF(IX.GT.1.AND.IX.LT.NX) GO TO 902
      IZ1=IZ1+1
      NLST(IZ1)=INOD1
      GO TO 903
902   ISE1=ISE1+1
      NLST(ISE1)=INOD1
903   INOD1=INOD1+1
      6   CONTINUE
      5   CONTINUE
      4   CONTINUE
C
      RETURN
      END
C
C*****
C      SUBROUTINE LCSTIF(STIF1,DRVTS,DMTRX,FNC,DNRM,BSAVE,X,Y,Z,WRK,NGS,LCS0000
+      NDIM,NDIM3,NLST,IRED,NX,NY,NZ,ITAPE,AJAC)
C*****
C      LCS0000
C      LCS0000
C A DOUBLE PRECISION IN-CORE PROGRAM TO FORM THE LOCAL STIFFNESS MATRIX
C OF COMPLEX CUBIC ELEMENTS. AT THE FIRST, THE STIFFNESS MATRIX IS
C FORMED FOR LAGRANGE ELEMENT, AND THEN A REDUCTION IS PERFORMED TO
C REDUCE THE L.G. STIFFNESS MATRIX TO A SERENDIPITY ELEMENT STIFFNESS
C MATRIX, (i.e. NODAL POINTS ON THE EDGES OF THE CUBIC ONLY.).
C
C-----
C      STIFF ..... LOCAL STIFFNESS MATRIX WITH DIMENSION 3*NOD x 3*NOD

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C      DRVTS ..... MATRIX STORAGE WHICH CONTAINS THE DERIVATIVES OF THE LCS00015
C      NONLINEAR INTERPOLATION FUNCTION AT EACH GAUSS POINT. LCS00016
C      (i.e. IT HAS DIMENSION NOD x 3 x 6; LCS00017
C      3 ---- INDICATES DERIVATIVES AT EACH DIRECTION; LCS00018
C      6 ---- INDICATES THE NUMBER OF GAUSS POINTS. LCS00019
C      DMTRX ..... A 6 x 6 MATRIX CONTAINS THE MATERIAL CONSTANT MATRIX. LCS00020
C      FNC ..... MATRIX STORAGE WHICH CONTAINS THE VALUES OF THE LCS00021
C      NONLINEAR INTERPOLATION FUNCTIONS OBTAINED AT EACH LCS00022
C      GAUSS POINT. IT HAS SAME DIMENSION AS DRVTS. LCS00023
C      DNRM ..... MATRIX STORAGE WITH DIMENSION NOD x 3, WHICH CONTAINS LCS00024
C      THE VALUES OF THE NORMALIZATION FACTORS. LCS00025
C      X,Y,Z ..... ARRAYS OF LENGTH NOD, WHICH CONTAIN THE X, Y, AND Z LCS00026
C      COORDINATES OF NODES IN GLOBAL COORDINATES SYSTEM; LCS00027
C      WHERE NOD=NX*NY*NZ. LCS00028
C      BSAVE ..... STORAGE OF CALCULATED DERIVATIVES OF SHAPE FUNCTIONS LCS00029
C      WHICH ARE USED TO FORM THE [B]-MATRIX. LCS00030
C      WRK ..... A 3*NOD x 3*NOD WORKING ARRAY FOR WORKING PACE. LCS00031
C      WGHT ..... VECTOR WITH LENGTH 6 CONTAINS THE WEIGHTS FOR EACH LCS00032
C      GAUSS POINT. LCS00033
C      NGS ..... ACTUAL REQUIRED GAUSS POINTS. LCS00034
C      NDIM ..... INITIAL DIMENSION OF X, Y, AND Z. LCS00035
C      NDIM3 ..... INITIAL DIMENSION OF STIFFNESS MATRIX. LCS00036
C      NLST ..... A LOCAL NODAL NUMBER LIST FOR REDUCING THE LAGRANGIAN LCS00037
C      ELEMENT STIFFNESS MATRIX TO SEREDIPITY ONE PROPERLY. LCS00038
C      IRED ..... INDICATOR IF THE CONDENCING PROCEDURE IS REQUIRED; LCS00039
C      IRED = 0 ... NO CONDENCE; IRED = 1 ... CONDENCE REQ. LCS00040
C      NX,NY,NZ ..... NUMBER OF NODES ALONG X, Y, AND Z DIRECTION, RESPECT. LCS00041
C      ITAPE ..... LOGICAL UNIT NUMBER ON WHICH THE STIFFNESS MATRIX OF LCS00042
C      LAGRANGIAN ELEMENT IS SAVE IF IT IS NECESSERARY. IF LCS00043
C      ITAPE = 0, NO SAVE IS DONE. LCS00044
C      ----- LCS00045
C      LCS00046
C      LCS00047
C      IMPLICIT REAL*8 (A-H,O-Z) LCS00048
C      DIMENSION STIF1(NDIM3,NDIM3),DRVTS(NDIM,3,10),DMTRX(6,6) LCS00049
C      + FNC(NDIM,3,10),DNRM(NDIM,3),BSAVE(NDIM,3),X(NDIM), LCS00050
C      + Y(NDIM),Z(NDIM),WRK(NDIM3,NDIM3),NLST(NDIM),WGHT(10) LCS00051
C      LCS00052
C      NOD=NX*NY*NZ !FIND THE TOTAL NUMBER OF NODES. LCS00053
C      NOD3=3*NOD LCS00054
C      NSE=8+4*(NX+NY+NZ-6) !FIND # NODES OF SEREN. ELMEMENT. LCS00055
C      NSE3=3*NSE LCS00056
C      LCS00057
C      OBTAIN THE WEIGHTS CORRESPONDING TO THE REQUIRED NUMBER OF GAUSS LCS00058
C      POINTS. LCS00059
C      LCS00060
C      CALL NUMINT(X,NDIM,NX,NGS,WGHT,1) LCS00061
C      LCS00062
C      OBTAIN THE GLOBAL COORDINATES OF A LAGRANGIAN ELEMENT FROM THE GIVEN LCS00063
C      COORDINATES OF THE BOUNDARY NODES. LCS00064
C      LCS00065
C      CALL GCOORD(X,Y,Z,NX,NY,NZ,NDIM,NLST) LCS00066
C      LCS00067
C      RENUMBERING THE GLOBAL COORDINATES CALCULATED FROM ABOVE. LCS00068
C      LCS00069
C      CALL RENUMB(X,Y,Z,NX,NY,NZ,NDIM,0,NLST) LCS00070
C      LCS00071
C      EVALUATE THE LOCAL STIFFNESS MATRIX IN LAGRANGIAN ELEMENT FORM. LCS00072
C      LCS00073
C      CALL STIFF1(WRK,DRVTS,DMTRX,FNC,DNRM,STIF1,WGHT,NGS,NOD,NDIM3, LCS00074

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+          NDIM,X,Y,Z,BSAVE,AJAC)
C
  IF(ITAPE.NE.0) WRITE(ITAPE) WRK
  IF(IREQ.EQ.0) THEN
  DO 1 I=1,NOD3
  DO 1 J=1,NOD3
  STIF1(I,J)=WRK(I,J)
  STIF1(I,I)=DABS(WRK(I,I))
  IF(I.NE.J) STIF1(J,I)=STIF1(I,J)
1  CONTINUE
C
  ELSE
C
  CONDENSE OUT THE INTERIOR NODES FROM [STIF1] SUCH THAT AT RETURN
  THE STIFFNESS MATRIX REPRESENTS THE BOUNDARY NODES ONLY.
C
  [K] =  $\begin{bmatrix} [K11] & [K12] \\ [K21] & [K22] \end{bmatrix}$ ;   [K'] = [K11] - [K12]*[K22]-1*[K21]
C
  REFORM THE STIFFNESS MATRIX SO THAT THE CONDENSING PROCEDURE CAN BE
  DONE CORRECTLY.
C
  CALL REFORM(WRK,STIF1,NDIM3,NLST,NDIM,NOD)
  CALL SWITCH(WRK,STIF1,NDIM3,NOD,NSE,0)
C
  FORM A R-H-S MATRIX IN ORDER TO CALCULATE [K22]-1*[K21].
C
  NK22=NOD3-NSE3
  DO 3 I=1,NK22
  DO 3 J=1,NSE3
  JK21=NK22+J
  STIF1(I,J)=WRK(I,JK21)
3  CONTINUE
C
  CALCULATE [K22]-1*[K21].
C
  CALL SOLVEQ(WRK,STIF1,NK22,NDIM3,NSE3,NDIM3,1,0)
C
  CALCULATE [K'] = [K11] - [K12]*[K22]-1*[K21].
C
  DO 4 I=1,NSE3
  IK12=NK22+I
  IK11=IK12
  DO 5 J=1,NSE3
  SUM=0.0
  JK11=NK22+J
  DO 6 K=1,NK22
  SUM=SUM+WRK(IK12,K)*STIF1(K,J)
6  CONTINUE
  WRK(IK11,JK11)=WRK(IK11,JK11)-SUM
  IF(DABS(WRK(IK11,JK11)).LT.1.D-10) WRK(IK11,JK11)=0.0
5  CONTINUE
4  CONTINUE
C
  PUT [STIF1] BACK IN CORRECT ORDER.
C
  DO 7 I=1,NSE3

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C
C CALCULATE Z,X,AND Y VALUES IN THE NATURAL (NORMAL) COORDINATES
C
      IPT=0
      NNZ=NZ
      NNY=NY
      NNX=NX
      IF(NZ.EQ.0) NNZ=1
      IF(NY.EQ.0) NNY=1
      IF(NX.EQ.0) NNX=1
C START NODAL POINT NUMBERING.
      IF(NZ.GE.1) Z0=-1.-DZ1
      DO 1 IZ=1,NNZ
      Z0=Z0+DZ1
      IF(NY.GE.1) Y0=-1.-DY1
      DO 2 IY=1,NNY
      Y0=Y0+DY1
      IF(NX.GE.1) X0=-1.-DX1
      DO 3 IX=1,NNX
      IPT=IPT+1
      X0=X0+DX1
      X1(IPT)=X0
      Y1(IPT)=Y0
      Z1(IPT)=Z0
  3 CONTINUE
  2 CONTINUE
  1 CONTINUE
C
      RETURN
      END

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NCO00025
NCO00026
NCO00027
NCO00028
NCO00029
NCO00030
NCO00031
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NCO00051
NCO00052
NCO00053
NCO00054

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C
C*****
C SUBROUTINE NUMINT(X,NDIM,NX,NGS,XR,IWG)
C*****

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```

C
C NUMERICAL INTEGRATION BY GAUSS-LEGENDRE QUADRATURE FORMULA FOR
C INTERPOLATION FUNCTIONS IN NATURAL (NORMAL) COORDINATES ONLY.
C
C X ..... ARRAY OF LENGTH NX, WHICH CONTAINS THE NODAL
C NDIM ..... INITIAL DIMENSION OF ARRAY X.
C NX ..... ACTUAL LENGTH OF ARRAY X.(TOT # OF NODES)
C NGS ..... NUMBER OF GAUSS POINTS REQUIRED FOR THE INTEGRATIONS.
C IT HAS TO BE ONE OF THE FOLLOWING :
C NGS = 2, 3, 4, 5, 6,7,8,9,10.
C XR ..... ARRAY OF LENGTH NGS, WHICH CONTAIN THE INTEGRATED
C VALUES AT EACH GAUSS POINTS.
C IWG ..... EVALUATION INDICATOR
C IWG = 0 ----- EVALUATE FUNCTION AT GAUSS POINT ONLY
C (i.e. without multiplying the WEIGHTS.
C IWG = 1 ----- TRANSFER THE VALUES OF WEIGHTS INTO
C XR.
C IWG = 2 ----- PERFORM THE NUMERICAL INTEGRATION
C COMPLETELY - i.e. G=SUM(wt(i)F(gs(i)))

```

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C-----
C
      IMPLICIT REAL*8 (A-H,O-Z)
      DIMENSION GS2(2),WT2(2),GS3(3),WT3(3),GS4(4),WT4(4),
+             GS5(5),WT5(5),GS6(6),WT6(6),GS7(7),WT7(7),
+             GS8(8),WT8(8),GS9(9),WT9(9),GS10(10),

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      +          WT10(10),X(NDIM),XR(10)
C
C DATA BLOCKS OF GAUSS POINTS AND CORRESPONDING WEIGHTS.
C
C 1. DATA BLOCK FOR TWO-POINT GAUSS QUADRATURE
C
C   DATA GS2/-0.5773502692, 0.5773502692/
C   DATA WT2/ 1.0000000000, 1.0000000000/
C
C 2. DATA BLOCK FOR THREE-POINT GAUSS QUADRATURE
C
C   DATA GS3/-0.7745966692, 0.0000000000, 0.7745966692/
C   DATA WT3/ 0.5555555555, 0.8888888889, 0.5555555555/
C
C 3. DATA BLOCK FOR FOUR-POINT GAUSS QUADRATURE
C
C   DATA GS4/-0.8611363116,-0.3399810435, 0.3399810435,
C   +          0.8611363116/
C   DATA WT4/ 0.3478548451, 0.6521451548, 0.6521451548,
C   +          0.3478548451/
C
C 4. DATA BLOCK FOR FIVE-POINT GAUSS QUADRATURE
C
C   DATA GS5/-0.9061798459,-0.5384693101, 0.0000000000,
C   +          0.5384693101, 0.9061798459/
C   DATA WT5/ 0.2369268850, 0.4786286705, 0.5688888889,
C   +          0.4786286705, 0.2369268850/
C
C 5. DATA BLOCK FOR SIX-POINT GAUSS QUADRATURE
C
C   DATA GS6/-0.9324695142,-0.6612093865,-0.2386191861,
C   +          0.2386191861, 0.6612093865, 0.9324695142/
C   DATA WT6/ 0.1713244924, 0.3607615730, 0.4679139346,
C   +          0.4679139346, 0.3607615730, 0.1713244924/
C
C 6. DATA BLOCK FOR SEVEN-POINT GAUSS QUADRATURE
C
C   DATA GS7/-0.9491079123,-0.7415311856,-0.4058451514,
C   +          0.0000000000, 0.4058451514, 0.7415311856,
C   +          0.9491079123/
C   DATA WT7/ 0.1294849662, 0.2797053915, 0.3818300505,
C   +          0.4179591837, 0.3818300505, 0.2797053915,
C   +          0.1294849662/
C
C 7. DATA BLOCK FOR EIGHT-POINT GAUSS QUADRATURE
C
C   DATA GS8/-0.9602898565,-0.7966664774,-0.5255324099,
C   +          -0.1834346425, 0.1834346425, 0.5255324099,
C   +          0.7966664774, 0.9602898565/
C   DATA WT8/ 0.1012285363, 0.2223810345, 0.3137066459,
C   +          0.3626837834, 0.3626837834, 0.3137066459,
C   +          0.2223810345, 0.1012285363/
C
C 8. DATA BLOCK FOR NINE-POINT GAUSS QUADRATURE
C
C   DATA GS9/-0.9681602395,-0.8360311073,-0.6133714327,
C   +          -0.3242534234, 0.0000000000, 0.3242534234,
C   +          0.6133714327, 0.8360311073, 0.9681602395/
C   DATA WT9/ 0.0812743884, 0.1806481607, 0.2606106964,
C   +          0.3123470770, 0.3302393550, 0.3123470770,

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+          0.2606106964, 0.1806481607, 0.0812743884/
C
C 9. DATA BLOCK FOR TEN-POINT GAUSS QUADRATURE
C
DATA GS10/-0.9739065285,-0.8650633667,-0.6794095683,
+          -0.4333953941,-0.1488743390, 0.1488743390,
+          0.4333953941, 0.6794095683, 0.8650633667,
+          0.9739065285/
DATA WT10/ 0.0666713443, 0.1494513492, 0.2190863625,
+          0.2692667193, 0.2955242247, 0.2955242247,
+          0.2692667193, 0.2190863625, 0.1494513492,
+          0.0666713443/
C
----- END OF DATA BLOCKS -----
C
IF(NX.LT.0) GO TO 99
IDN=IWG-1
C
C START EVALUATION PROCEDURE (POINT BY POINT) :
C
DO 1 IG=1,NGS
  XX=1.0
  IF(NGS.EQ.2) THEN
    WT=WT2(IG)
    GS=GS2(IG)
  ENDIF
  IF(NGS.EQ.3) THEN
    WT=WT3(IG)
    GS=GS3(IG)
  ENDIF
  IF(NGS.EQ.4) THEN
    WT=WT4(IG)
    GS=GS4(IG)
  ENDIF
  IF(NGS.EQ.5) THEN
    WT=WT5(IG)
    GS=GS5(IG)
  ENDIF
  IF(NGS.EQ.6) THEN
    WT=WT6(IG)
    GS=GS6(IG)
  ENDIF
  IF(NGS.EQ.7) THEN
    WT=WT7(IG)
    GS=GS7(IG)
  ENDIF
  IF(NGS.EQ.8) THEN
    WT=WT8(IG)
    GS=GS8(IG)
  ENDIF
  IF(NGS.EQ.9) THEN
    WT=WT9(IG)
    GS=GS9(IG)
  ENDIF
  IF(NGS.EQ.10) THEN
    WT=WT10(IG)
    GS=GS10(IG)
  ENDIF
C EVALUATE EACH TERM OF THE INTERPOLATION FUNCTIONS AT THE GAUSS POINT.
C

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          IF(IDN) 96,97,96
96       IF(NX.NE.0) THEN
          DO 2 IX=1,NX
          XX=XX*(GS-X(IX))
          2   CONTINUE
          ELSE
          XX=1.
          ENDIF
          XR(IG)=XX
          IF(IWG.EQ.2) XR(IG)=WT*XX
          GO TO 1
97       XR(IG)=WT
          1   CONTINUE
C
C   END OF GAUSS-LEGENDRE QUADRATURE EVALUATIONS
C
          RETURN
99       WRITE(6,1001) NX
1001    FORMAT(/,
          + ' ***** ERROR IN NUMERICAL INTEGRATION, NO FUNCTIONS EXIST *****',
          +/, ' ***** ORDER OF INTERPOLATION FUNCTION NITP =',I8)
          RETURN
          END
C
C ***** REF00001
C ***** REF00002
          SUBROUTINE REFORM(STIF1,STIF2,NDIM1,NLST,NDIM2,NPE) REF00003
C ***** REF00004
C ***** REF00005
C THIS SUBROUTINE RE-ORDERS THE STIFFNESS MATRIX SUCH THAT IT HAS THE REF00006
C "EDGE-NODES" IN THE FRONT AND ALL OF THE "OFF-EDGE" NODES ON THE REF00007
C BACK. REF00008
C REF00009
C ----- REF00010
C STIF1 ..... ARRAY OF DIMENSION 3*NPE X 3*NPE. IT CONTAINS THE INPUT REF00011
C STIFFNESS MATRIX OF A LAGURANGIAN ELEMENT. REF00012
C STIF2 ..... ARRAY OF DIMENSION 3*NPE X 3*NPE. AFTER RETURN IT HAS REF00013
C THE RE-ORDERED STIFFNESS MATRIX. REF00014
C NDIM1 ..... INITIAL DIMENSION OF MATRICES "STIF1" AND "STIF2". REF00015
C NLST ..... LOCAL NODAL NUMBER LISTING. REF00016
C NDIM2 ..... INITIAL DIMENSION OF VECTOR "NLST". REF00017
C NPE ..... TOTAL NUMBER OF NODES OF THE GIVEN LAGURANGIAN ELEMENT. REF00018
C ----- REF00019
C
C   IMPLICIT REAL*8 (A-H,O-Z) REF00020
C   DIMENSION STIF1(NDIM1,NDIM1),STIF2(NDIM1,NDIM1),NLST(NDIM2) REF00021
C
C INITIALIZE STIF2 ... REF00022
C REF00023
C   DO 1 I=1,NDIM3 REF00024
C   DO 1 J=1,NDIM3 REF00025
          1   STIF2(I,J)=0.0 REF00026
C REF00027
C REF00028
C PERFORMING THE RE-ORDERING ..... REF00029
C REF00030
C   DO 2 I=1,NPE REF00031
C   IROW1=3*(I-1) REF00032
C   JROW1=3*(NLST(I)-1) REF00033
C   DO 3 J=1,NPE REF00034

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ICOL1=3*J-3
JCOL1=3*(NLST(J)-1)
DO 4 KI=1,3
IROW=IROW1+KI
JROW=JROW1+KI
DO 5 KJ=1,3
ICOL=ICOL1+KJ
JCOL=JCOL1+KJ
STIF2(IROW,ICOL)=STIF1(JROW,JCOL)
5 CONTINUE
4 CONTINUE
3 CONTINUE
2 CONTINUE
C
RETURN
END
C
C*****
SUBROUTINE RENUMB(X,Y,Z,NX,NY,NZ,NDIM,NDIR,NLST)
C*****
C THIS PROGRAM RENUMBERS THE NODAL NUMBER SEQUENCE SUCH THAT IT FORMS
C A LAGUANGIAN ELEMENT FROM SERENDIPITY ELEMENT, OR V.V.S..
C-----
C X,Y,Z ..... X, Y, AND Z COORDINATES WHICH ARE IN THE SEQUENCE THAT
C FIRST "NSE" VALUES ARE NODES ON EDGES AND REST OF THEM
C (i.e. NX*NY*NZ-NSE) ARE THE NODES OFF THE EDGES.
C NODE: ALL OF THE VALUES MUST BE IN THE CORRECT SEQUENCE.
C NX,NY,NZ ..... NUMBER OF NODES ALONG X, Y, AND Z DIRECTION, RESPECT.
C NDIM ..... INITIAL DIMENSION OF THE CALLING PROGRAM.
C NDIR ..... OPERATING INDICATOR :
C NDIR = 0 ... RENUMBER TO L. ELEMENT;
C NDIR = 1 ... RENUMBER TO S. ELEMENT.
C NLST ..... LOCAL NODAL NUMBER LIST.
C-----
C
IMPLICIT REAL*8 (A-H,O-Z)
DIMENSION X(NDIM),Y(NDIM),Z(NDIM),NLST(NDIM)
C
C FIND THE TOTAL NUMBERS OF NQDES IN SERENDIPITY ELEMENT AND
C CORRESPONDING LAGURANGIAN ELEMENT.
C
NSE=8+4*(NX+NY+NZ-6)      !# OF NODES OF S. ELEMENT
NLE=NX*NY*NZ              !# OF NODES OF L. ELEMENT
NDF=NLE-NSE                !DIFFERENCE BTW. NSE & NLE
C
IF(NDIR.EQ.1) GO TO 9999
C
DO 1 I=1,NDF
IST=NSE+I
ILC=NLST(IST)
NCH=IST-ILC
XSAVE=X(ILC)
YSAVE=Y(ILC)
ZSAVE=Z(ILC)
X(ILC)=X(IST)
Y(ILC)=Y(IST)
Z(ILC)=Z(IST)
IF(NCH.LT.1) GO TO 1

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DO 2 J=1,NCH                                REN00045
JLC=ILC+J                                    REN00046
XSAV1=X(JLC)                                REN00047
YSAV1=Y(JLC)                                REN00048
ZSAV1=Z(JLC)                                REN00049
X(JLC)=XSAVE                                 REN00050
Y(JLC)=YSAVE                                 REN00051
Z(JLC)=ZSAVE                                 REN00052
XSAVE=XSAV1                                  REN00053
YSAVE=YSAV1                                  REN00054
ZSAVE=ZSAV1                                  REN00055
2 CONTINUE                                    REN00056
1 CONTINUE                                    REN00057
C                                              REN00058
RETURN                                        REN00059
9999 CONTINUE                                REN00060
RETURN                                        REN00061
END                                            REN00062
C                                              SHA00001
C*****SHA00002
SUBROUTINE SHAPE1 (XX,YY,ZZ,DRVTS,FNC,DNRM,NDIM,NX,NY,NZ,NGS,WRK,
+                NWRK,IGEN)                  SHA00003
C*****SHA00005
C CALCULATE THE INTERPOLATION FUNCTIONS ABOUT EVERY NODAL POINT BY
C USING GAUSS-LEGENDRE QUADRATURE FORMULA.    SHA00008
C                                              SHA00009
C XX ..... ARRAY OF NORMAL (natural) NODAL COORDINATES IN XX.    SHA00010
C YY ..... ARRAY OF NORMAL (natural) NODAL COORDINATES IN YY.    SHA00011
C ZZ ..... ARRAY OF NORMAL (natural) NODAL COORDINATES IN ZZ.    SHA00012
C DRVTS ..... ARRAY OF DIMENSION (NPE x 3 x NGS), WHICH CONTAINS THE
C EVALUATED DERIVATIVES ABOUT NORMAL COORDINATES SYSTEM          SHA00014
C AT EACH GAUSS POINTS; Where                                     SHA00015
C (NPE --- number of nodes per element)                          SHA00016
C FNC ..... ARRAY OF NPE x 3 x 10, WHICH CONTAINS THE VALUES OF THE
C INTERPOLATION FUNCTION AT EACH GAUSS POINT.                     SHA00018
C DNRM ..... ARRAY OF NPE x 3, WHICH CONTAINS THE VALUES OF THE
C NORMALIZATION FACTORS IN X, Y, AND Z AT EACH NODE.              SHA00020
C NDIM ..... FIRST ROW DIMENSION OF X,Y,Z,XX,YY,ZZ AND DRVTS IN
C CALLING PROGRAM.(EQ # OF NODES)                                 SHA00022
C NX ..... NUMBER OF NODES ALONG X.                                SHA00023
C NY ..... NUMBER OF NODES ALONG Y.                                SHA00024
C NZ ..... NUMBER OF NODES ALONG Z.                                SHA00025
C NGS ..... NUMBER OF GAUSS POINTS REQUIRED IN NUMERICAL INTEGRATION.
C WRK ..... WORKING SPACE OF DIMENSION NWRK X NWRK.              SHA00027
C NWRK ..... INITIAL DIMENSION OF MATRIX WRK.                    SHA00028
C IGEN ..... NORMAL COORDINATES GENERATION INDICATOR :          SHA00029
C IGEN = 0   NORMAL COORDINATES ARE GENERATED;                   SHA00030
C IGEN = 1   NORMAL COORDINATES ARE FROM INPUT.                   SHA00031
C .....SHA00032
C SUBROUTINES CALLED :                                           SHA00033
C DERVTS ---- CALCULATE DERIVATIVES.                              SHA00036
C NCOORD ---- GENERATE A SET NORMAL NODAL COORDINATES.          SHA00037
C NUMINT ---- PERFORM THE NUMERICAL INTEGRATIONS.                SHA00038
C .....SHA00039
C .....SHA00040
C .....SHA00041
C .....SHA00042
IMPLICIT REAL*8 (A-H,O-Z)

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DIMENSION XX(NDIM),YY(NDIM),ZZ(NDIM),DRVTS(NDIM,3,10),
+   WRK(NWRK,NWRK),DNRM(NDIM,3),XINT1(10),XINT2(10),XINT3(10),
+   DVXX(10),DVYY(10),DVZZ(10),FNC(NDIM,3,10)
C
NPE=NX*NY*NZ
IWRK1=0
IWRK2=IWRK1+NDIM
IWRK3=IWRK2+NDIM
C
C INITIALIZE XX, YY, ZZ, VWRK, WRK, DRVTS AND FNC.
C
DO 2 I=1,NDIM
DO 1 J=1,3
DO 1 K=1,NGS
DRVTS(I,J,K)=0.0
FNC(I,J,K)=0.0
1 CONTINUE
IK1=IWRK1+I
IK2=IWRK2+I
IK3=IWRK3+I
IF(IGEN.EQ.0) THEN
XX(I)=0.0
YY(I)=0.0
ZZ(I)=0.0
ELSE
WRK(IK1,1)=XX(I)
WRK(IK2,1)=YY(I)
WRK(IK3,1)=ZZ(I)
ENDIF
WRK(I,2)=0.0
2 CONTINUE
C
C GENERATE NODAL COORDINATES IN NORMAL (NATURAL) COORDINATES
C IF IT IS REQUIRED.
C
IF(IGEN.EQ.0) THEN
CALL NCOORD(XX,YY,ZZ,NX,NY,NZ,NDIM)
DO IKK=1,NPE
IK1=IWRK1+IKK
IK2=IWRK2+IKK
IK3=IWRK3+IKK
WRK(IK1,1)=XX(IKK)
WRK(IK2,1)=YY(IKK)
WRK(IK3,1)=ZZ(IKK)
END DO
END IF
C
C CALCULATE THE NORMALIZED INTERPOLATION FUNCTIONS (I.F.) AND THE
C CORRESPONDING DERIVATIVES IN X, Y, AND Z, RESPECTIVELY.
C df/dXX=Sum(df(i)/dXX), i=1,NPE; etc....
C where
C   F(i)=Fi(XX)*Fi(YY)*Fi(ZZ),           i.e.
C   df(i)/dXX=Fi(YY)*Fi(ZZ)*(dFi(XX)/dXX) etc.....
C
DO 4 IPE=1,NPE           !Evaluations node by node
TRAS1=1.0
TRAS2=1.0
TRAS3=1.0
IPTX=0
IPTY=0

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IPTZ=0
IK1=IWRK1+IPE
IK2=IWRK2+IPE
IK3=IWRK3+IPE
DO IXYZ=1,NPE
XX(IXYZ)=0.
YY(IXYZ)=0.
ZZ(IXYZ)=0.
END DO
DO 5 INOD=1,NPE
IF(INOD.EQ.IPE) GO TO 5 !Only choose i .ne. j the term
IKN1=IWRK1+INOD
IKN2=IWRK2+INOD
IKN3=IWRK3+INOD
C
C EVALUATE NORMALIZATION FACTORS AND CHOOSE APPROPRIATE NODAL POINTS.
C
IF(WRK(IKN2,1).EQ.WRK(IK2,1).AND.WRK(IKN3,1).EQ.WRK(IK3,1))
+ THEN
IPTX=IPTX+1 !Interpolation function along XX
TERM1=WRK(IK1,1)-WRK(IKN1,1) !Normalization factor in XX
TRAS1=TRAS1*TERM1
XX(IPTX)=WRK(IKN1,1) !XX-points of interpolation func
ENDIF
IF(WRK(IKN1,1).EQ.WRK(IK1,1).AND.WRK(IKN3,1).EQ.WRK(IK3,1))
+ THEN
IPTY=IPTY+1 !Interpolation function along YY
TERM2=WRK(IK2,1)-WRK(IKN2,1) !Normalization factor in YY
TRAS2=TRAS2*TERM2
YY(IPTY)=WRK(IKN2,1) !YY-points of interpolation func
ENDIF
IF(WRK(IKN1,1).EQ.WRK(IK1,1).AND.WRK(IKN2,1).EQ.WRK(IK2,1))
+ THEN
IPTZ=IPTZ+1 !Interpolation function along ZZ
TERM3=WRK(IK3,1)-WRK(IKN3,1) !Normalization factor in ZZ
TRAS3=TRAS3*TERM3
ZZ(IPTZ)=WRK(IKN3,1) !ZZ-points of interpolation func
ENDIF
S CONTINUE
DNRM(IPE,1)=TRAS1
DNRM(IPE,2)=TRAS2
DNRM(IPE,3)=TRAS3
C
C PERFORM NUMERICAL INTEGRATIONS WITHOUT WEIGHTS (IWG=0)
C
CALL NUMINT(XX,NDIM,IPTX,NGS,XINT1,0)
CALL NUMINT(YY,NDIM,IPTY,NGS,XINT2,0)
CALL NUMINT(ZZ,NDIM,IPTZ,NGS,XINT3,0)
C
C CALCULATE DERIVATIVES ABOUT X, Y, AND Z.
C
CALL DERVTS(DVXX,XX,NDIM,IPTX,WRK,NGS,NWRK,2)
CALL DERVTS(DVYY,YY,NDIM,IPTY,WRK,NGS,NWRK,2)
CALL DERVTS(DVZZ,ZZ,NDIM,IPTZ,WRK,NGS,NWRK,2)
C
C Saving the evaluated values
C
DO 10 IGS=1,NGS
FNC(IPE,1,IGS)=XINT1(IGS)
FNC(IPE,2,IGS)=XINT2(IGS)

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| | | |
|----|---------------------------------------------------------------------|-----------|
| | FNC(IPE,3,IGS)=XINT3(IGS) | SHA00163 |
| | DRVTS(IPE,1,IGS)=DVXX(IGS) | SHA00164 |
| | DRVTS(IPE,2,IGS)=DVYY(IGS) | SHA00165 |
| | DRVTS(IPE,3,IGS)=DVZZ(IGS) | SHA00166 |
| 10 | CONTINUE | SHA00167 |
| 4 | CONTINUE | SHA00168 |
| C | | SHA00169 |
| C | RETURN BACK XX, YY, AND ZZ VALUE. | SHA00170 |
| C | | SHA00171 |
| | DO 11 I=1,NPE | SHA00172 |
| | IK1=IWRK1+I | SHA00173 |
| | IK2=IWRK2+I | SHA00174 |
| | IK3=IWRK3+I | SHA00175 |
| | XX(I)=WRK(IK1,1) | SHA00176 |
| | YY(I)=WRK(IK2,1) | SHA00177 |
| | ZZ(I)=WRK(IK3,1) | SHA00178 |
| 11 | CONTINUE | SHA00179 |
| C | | SHA00180 |
| C | FINISH | SHA00181 |
| C | | SHA00182 |
| | RETURN | SHA00183 |
| | END | SHA00184 |
| C | | SOL00001 |
| C | ***** | SOL00002 |
| | SUBROUTINE SOLVEQ(A,SAV1,B,NEQ,NRD,NB,NBD,ISYM,JUMP) | SOL00003 |
| C | ***** | SOL00004 |
| C | | SOL00005 |
| C | A IN CORE LINEAR EQUATION SOLVER SUBROUTINE IN DOUBLE PRECISION. | SOL00006 |
| C | IT CALLS A LU-DECOMPOSITION SUBROUTINE "DECOMP" | SOL00007 |
| C | | SOL00008 |
| C | ----- | SOL00009 |
| C | A DECOMPOSED COEFFECIENT MATRIX IN LU FORM | SOL00010 |
| C | B SOLUTION AFTER SOLVING SYSTEM OF EQUATIONS | SOL00011 |
| C | SAV1 INPUT R.H.S VECTORS | SOL00011A |
| C | NEQ NUMBER OF EQUATIONS | SOL00012 |
| C | NB NUMBER OF RIGHT HAND SIDE VECTORS | SOL00013 |
| C | NBD INITIAL COLUMN DIMENSION OF MATRIX [B] | SOL00014 |
| C | ISYM SYMMETRIC INDICATOR; ISYM=0 ... NONSYM.; ISYM=1 ... SYM. | SOL00015 |
| C | JUMP OPERATING INDICATOR; JUMP=1 ... LU-DECOMP. IS SKIPPED | SOL00016 |
| C | ----- | SOL00017 |
| C | | SOL00018 |
| | IMPLICIT REAL*8 (A-H,O-Z) | SOL00019 |
| | DIMENSION A(NRD,NRD),B(NRD,NBD),SAV1(NRD,NBD) | SOL00020 |
| C | | SOL00021 |
| C | SAVE R.H.S VECTORS BEFORE SYSTEM OF EQUATIONS ARE SOLVED | SOL00021A |
| C | | SOL00021B |
| | DO 7 I=1,NEQ | SOL00021C |
| | DO 7 J=1,NBD | SOL00021D |
| 7 | B(I,J)=SAV1(I,J) | SOL00021E |
| C | | SOL00021F |
| C | FORM LU DECOMPOSITION FORM | SOL00022 |
| C | | SOL00023 |
| | IF(JUMP.NE.1) THEN | SOL00023A |
| | DO 8 I=1,NEQ | SOL00023B |
| | DO 8 J=1,NEQ | SOL00023C |
| 8 | SAV2(I,J)=A(I,J) | SOL00023D |
| | CALL DECOMP(SAV2,A,NEQ,NRD,ISYM) | SOL00023E |
| | ENDIF | SOL00024 |
| C | | SOL00025 |
| C | CALCULATE MATRIX [Y] FROM EQUATION [L][Y]=[B] | SOL00026 |

```

C SOL00027
DO 1 ICOL=1,NB SOL00028
B(1,ICOL)=B(1,ICOL)/A(1,1) SOL00029
DO 2 IROW=2,NEQ SOL00030
SUM=0. SOL00031
DO 3 ISUM=1,IROW-1 SOL00032
SUM=SUM+A(IROW,ISUM)*B(ISUM,ICOL) SOL00033
3 CONTINUE SOL00034
B(IROW,ICOL)=(B(IROW,ICOL)-SUM)/A(IROW,IROW) SOL00035
2 CONTINUE SOL00036
1 CONTINUE SOL00037
C SOL00038
C OBTAIN THE FINAL SOLUTION FROM [U][X]=[Y] SOL00039
C SOL00040
DO 4 ICOL=1,NB SOL00041
IBRW=NEQ SOL00042
IF(ISYM.EQ.1) B(IBRW,ICOL)=B(IBRW,ICOL)/A(NEQ,NEQ) SOL00043
DO 5 IROW=2,NEQ SOL00044
ISTR=IBRW SOL00045
IBRW=IBRW-1 SOL00046
SUM=0. SOL00047
DO 6 ISUM=ISTR,NEQ SOL00048
SUM=SUM+A(IBRW,ISUM)*B(ISUM,ICOL) SOL00049
6 CONTINUE SOL00050
B(IBRW,ICOL)=B(IBRW,ICOL)-SUM SOL00051
IF(ISYM.EQ.1) B(IBRW,ICOL)=B(IBRW,ICOL)/A(IBRW,IBRW) SOL00052
5 CONTINUE SOL00053
4 CONTINUE SOL00054
C SOL00055
C EQUATION HAS BEEN SOLVED SOL00056
C SOL00057
RETURN SOL00058
END SOL00059
C STI00001
C*****STI00001
SUBROUTINE STIFF1(STIFF,DRVTS,DMTRX,FNC,DNRM,WRK,NGS,NOD, STI00001
+ NDIM3,NDIM,X,Y,Z,BSAVE,AJAC) STI00004
C*****STI00005
C STI00006
C THIS SUBROUTINE CALCULATES THE LOCAL STIFFNESS MATRIX BY USING GAUSS STI00007
C QUADRATURE FORMULA FOR ELEMENTS WITH NONLINEAR INTERPOLATION STI00008
C FUNCTIONS. STI00009
C STI00010
C STIFF ..... LOCAL STIFFNESS MATRIX WITH DIMENSION 3*NOD x 3*NOD STI00011
C DRVTS ..... MATRIX STORAGE WHICH CONTAINS THE DERIVATIVES OF THE STI00012
C NONLINEAR INTERPOLATION FUNCTION AT EACH GAUSS POINT. STI00013
C (i.e. IT HAS DIMENSION NOD x 3 x 10; STI00014
C 3 --- INDICATES DERIVATIVES AT EACH DIRECTION; STI00015
C 10 --- INDICATES THE NUMBER OF GAUSS POINTS. STI00016
C DMTRX ..... A 6 x 6 MATRIX CONTAINS THE MATERIAL CONSTANT MATRIX. STI00017
C FNC ..... MATRIX STORAGE WHICH CONTAINS THE VALUES OF THE STI00018
C NONLINEAR INTERPOLATION FUNCTIONS OBTAINED AT EACH STI00019
C GAUSS POINT. IT HAS SAME DIMENSION AS DRVTS. STI00020
C DNRM ..... MATRIX STORAGE WITH DIMENSION NOD x 3, WHICH CONTAINS STI00021
C THE VALUES OF THE NORMALIZATION FACTORS. STI00022
C WRK ..... A NDIM3 X NDIM3 WORKING SPACE FOR WORKING PACE. S.I.00022
C WGHT ..... VECTOR OF LENGTH 10 CONTAINS THE WEIGHTS FOR EACH STI00023
C GAUSS POINT. STI00024
C NGS ..... ACTUAL REQUIRED GAUSS POINTS. STI00025
C NOD ..... NUMBER OF NODES FOR THE ELEMENT. STI00026

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C   NDIM3 ..... INITIAL DIMENSION OF STIFFNESS MATRIX.          STI00028
C   NDIM ..... INITIAL DIMENSION OF X, Y, AND Z.                STI00029
C   X,Y,Z ..... ARRAYS OF LENGTH NOD, WHICH CONTAIN THE X, Y, AND Z STI00030
C                                     COORDINATES OF NODES IN GLOBAL COORDINATES SYSTEM. STI00031
C                                     STI00032
C-----STI00033
C                                     STI00034
C                                     STI00035
C   IMPLICIT REAL*8 (A-H,O-Z)
C   DIMENSION STIFF(NDIM3,NDIM3),DRVTS(NDIM,3,10),DMTRX(6,6),    STI00036
C   +          BSAVE(NDIM,3),FNC(NDIM,3,10),DNRM(NDIM,3),        STI00037
C   +          WRK(NDIM3,NDIM3),WGHT(10),VWRK(6,3),DVX(3),DVY(3), STI00038
C   +          DVZ(3),AJACM(3,3),VWK(3),X(NDIM),Y(NDIM),Z(NDIM) STI00039
C   NGSX=NGS STI00040
C   NGSY=NGS STI00041
C   NGSZ=NGS STI00042
C                                     STI00043
C   INITIALIZE STIFF STI00044
C                                     STI00045
C   NDF=3*NOD STI00046
C   DO 1 IST=1,NDIM3 STI00047
C   DO 1 JST=1,NDIM3 STI00048
C 1 STIFF(IST,JST)=0.0 STI00049
C                                     STI00050
C   START EVALUATING GAUSS POINT BY GAUSS POINT STI00051
C   [STIFF] = Sum{WGHT[IGS].[STIFF(IGS)]}; IGS=1,NGS. STI00052
C   where the sumation has to be Tripled. STI00053
C   i.e. [STIFF] = Sum{wght(i)*Sum{wght(j)*Sum{wght(k)*STIFF[I,J,K]}} STI00054
C                                     STI00055
C   DO 2 IGSX=1,NGSX !Sum in XX-direction STI00056
C   DO 12 IGSY=1,NGSY !Sum in YY-direction STI00057
C   DO 22 IGSZ=1,NGSZ !Sum in ZZ-direction STI00058
C   DO IZR=1,3 STI00059
C     DVX(IZR)=0.0 STI00060
C     DVY(IZR)=0.0 STI00061
C     DVZ(IZR)=0.0 STI00062
C   END DO STI00063
C                                     STI00064
C   FORM THE SHAPE FUNCTIONS AND ITS DERIVATIVES STI00065
C   N(XXi,YYj,ZZk) = f(XXi)*g(YYj)*h(ZZk) STI00066
C   dN(XXi,YYj,ZZk)/dXX = (df(XXi)/dXX)*g(YYj)*h(ZZk) STI00067
C   dN(XXi,YYj,ZZk)/dYY = f(XXi)*(dg(YYj)/dYY)*h(ZZk) STI00068
C   dN(XXi,YYj,ZZk)/dZZ = f(XXi)*g(YYj)*(dh(ZZk)/dZZ) STI00069
C                                     STI00070
C   DO 14 INZ=1,NOD STI00071
C   AINT12=FNC(INZ,1,IGSX)*FNC(INZ,2,IGSY) !f(XXi)*g(YYj) STI00072
C   AINT13=FNC(INZ,1,IGSX)*FNC(INZ,3,IGSZ) !f(XXi)*h(ZZk) STI00073
C   AINT23=FNC(INZ,2,IGSY)*FNC(INZ,3,IGSZ) !g(YYj)*h(ZZk) STI00074
C   DNMT=DNRM(INZ,1)*DNRM(INZ,2)*DNRM(INZ,3) STI00075
C   DX=DRVTS(INZ,1,IGSX) !df(XXi)/dXX STI00076
C   DY=DRVTS(INZ,2,IGSY) !dg(YYj)/dYY STI00077
C   DZ=DRVTS(INZ,3,IGSZ) !dh(ZZk)/dZZ STI00078
C   WRK(INZ,7)=DX*AINT23/DNMT !dN(XXi,YYj,ZZk)d/XX STI00079
C   WRK(INZ,8)=DY*AINT13/DNMT !dN(XXi,YYj,ZZk)d/YY STI00080
C   WRK(INZ,9)=DZ*AINT12/DNMT !dN(XXi,YYj,ZZk)d/ZZ STI00081
C                                     STI00082
C   FORM JACOBIAN MATRIX STI00083
C                                     STI00084
C   dx(XXi,YYj,ZZk)/dXX= Sum{X(ip)*dN(XXi,YYj,ZZk)/dXX} STI00085
C   dx(XXi,YYj,ZZk)/dYY= Sum{X(ip)*dN(XXi,YYj,ZZk)/dYY} STI00086
C   dx(XXi,YYj,ZZk)/dZZ= Sum{X(ip)*dN(XXi,YYj,ZZk)/dZZ} STI00087

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| | | |
|----|----------------------------------------------------------------|----------|
| C | | STI00148 |
| C | STEP 1 : CALCULATE [WRK(IGS)] = Trans[B(IGS)].[D(IGS)] | STI00149 |
| C | | STI00150 |
| | IB1=3*(INOD-1)+1 | STI00151 |
| | IB2=IB1+1 | STI00152 |
| | IB3=IB2+1 | STI00153 |
| | DO 4 IWK=1,6 | STI00154 |
| | WRK(IB1,IWK)=0.0 | STI00155 |
| | WRK(IB2,IWK)=0.0 | STI00156 |
| | WRK(IB3,IWK)=0.0 | STI00157 |
| | DO 5 ISUM=1,6 | STI00158 |
| | WRK(IB1,IWK)=WRK(IB1,IWK)+VWRK(ISUM,1)*DMTRX(ISUM,IWK) | STI00159 |
| | WRK(IB2,IWK)=WRK(IB2,IWK)+VWRK(ISUM,2)*DMTRX(ISUM,IWK) | STI00160 |
| | WRK(IB3,IWK)=WRK(IB3,IWK)+VWRK(ISUM,3)*DMTRX(ISUM,IWK) | STI00161 |
| 5 | CONTINUE | STI00162 |
| | IF(DABS(WRK(IB1,IWK)).LT.1.D-10) WRK(IB1,IWK)=0.0 | STI00163 |
| | IF(DABS(WRK(IB2,IWK)).LT.1.D-10) WRK(IB2,IWK)=0.0 | STI00164 |
| | IF(DABS(WRK(IB3,IWK)).LT.1.D-10) WRK(IB3,IWK)=0.0 | STI00165 |
| 4 | CONTINUE | STI00166 |
| 3 | CONTINUE | STI00167 |
| C | | STI00168 |
| C | STEP 2 : EVALUATE LOCAL STIFFNESS MATRIX AT GAUSS POINT "IGS". | STI00169 |
| C | [STIFF(IGS)] = [WRK(IGS)].[B(IGS)] ; AT G.P. "IGS". | STI00170 |
| C | | STI00171 |
| | WGHTS=WGHT(IGSX)*WGHT(IGSY)*WGHT(IGSZ) | STI00172 |
| | DO 6 INOD=1,NOD | STI00173 |
| | IB1=3*(INOD-1)+1 | STI00174 |
| | IB2=IB1+1 | STI00175 |
| | IB3=IB2+1 | STI00176 |
| | VWR(1)=BSAVE(INOD,1) | STI00177 |
| | VWR(2)=BSAVE(INOD,2) | STI00178 |
| | VWR(3)=BSAVE(INOD,3) | STI00179 |
| | CALL FORMVK(VWRK,VWR) | STI00180 |
| | DO 7 IST=1,NDF | STI00181 |
| | SUM1=0.0 | STI00182 |
| | SUM2=0.0 | STI00183 |
| | SUM3=0.0 | STI00184 |
| | IF(IST.LT.IB1) GO TO 97 !AVOID REPEAT CALCULATIONS (DUE SYM.) | STI00185 |
| | DO 8 ISUM=1,6 | STI00186 |
| | SUM1=SUM1+WRK(IST,ISUM)*VWRK(ISUM,1) | STI00187 |
| | IF(IST.LT.IB2) GO TO 8 !AVOID REPEAT CALCULATIONS (DUE SYM.) | STI00188 |
| | SUM2=SUM2+WRK(IST,ISUM)*VWRK(ISUM,2) | STI00189 |
| | IF(IST.LT.IB3) GO TO 8 !AVOID REPEAT CALCULATIONS (DUE SYM.) | STI00190 |
| | SUM3=SUM3+WRK(IST,ISUM)*VWRK(ISUM,3) | STI00191 |
| 8 | CONTINUE | STI00192 |
| | STIFF(IST,IB1)=STIFF(IST,IB1)+AJAC*WGHTS*SUM1 | STI00193 |
| | STIFF(IST,IB2)=STIFF(IST,IB2)+AJAC*WGHTS*SUM2 | STI00194 |
| | STIFF(IST,IB3)=STIFF(IST,IB3)+AJAC*WGHTS*SUM3 | STI00195 |
| | IF(DABS(STIFF(IST,IB1)).LT.1.D-10) STIFF(IST,IB1)=0.0 | STI00196 |
| | IF(DABS(STIFF(IST,IB2)).LT.1.D-10) STIFF(IST,IB2)=0.0 | STI00197 |
| | IF(DABS(STIFF(IST,IB3)).LT.1.D-10) STIFF(IST,IB3)=0.0 | STI00198 |
| 97 | STIFF(IST,IB1)=STIFF(IST,IB1) | STI00199 |
| | STIFF(IST,IB2)=STIFF(IST,IB2) | STI00200 |
| | STIFF(IST,IB3)=STIFF(IST,IB3) | STI00201 |
| 7 | CONTINUE | STI00202 |
| 6 | CONTINUE | STI00203 |
| C | | STI00204 |
| 22 | CONTINUE | STI00205 |
| 12 | CONTINUE | STI00206 |
| 2 | CONTINUE | STI00207 |

```

C STI00208
C FINISH STI00209
C STI00210
C RETURN STI00211
C END STI00212
C SWI00001
C*****SWI00002
C SUBROUTINE SWITCH(STIF1,WRK,NDIM,NOD,NSE,IDR) SWI00003
C*****SWI00004
C SWI00005
C A IN-CORE SUBROUTINE WHICH SWITCHES BLOCK SUBMATRICES OF STIFFNESS SWI00006
C MATRIX FROM THE GIVEN FORM. SWI00007
C SWI00008
C MAKE [STIF1] =  $\begin{bmatrix} [K22] & [K21] \\ [K12] & [K11] \end{bmatrix}$  FROM [K] =  $\begin{bmatrix} [K11] & K[12] \\ [K21] & K[22] \end{bmatrix}$ ; SWI00009
C SWI00010
C OR VICE VER. SWI00011
C SWI00012
C SWI00013
C SWI00014
C-----SWI00015
C STIF1 ..... OUT PUT SWITCHED MATRIX. SWI00016
C WRK ..... INPUT MATRIX TO BE SWITCHED. SWI00017
C NDIM ..... INITIAL DIMENSION OF MATRICES [STIF1] AND [WRK]. SWI00018
C NOD ..... TOTAL NUMBER OF NODES OF A LAGRANGIAN ELEMENT. SWI00019
C NSE ..... NUMBER OF NODES ON THE BOUNDARY. SWI00020
C IDR ..... SWITCHING DIRECTION INDICATOR. SWI00021
C-----SWI00022
C SWI00023
C IMPLICIT REAL*8 (A-H,O-Z) SWI00024
C DIMENSION STIF1(NDIM,NDIM),WRK(NDIM,NDIM) SWI00025
C SWI00026
C DECIDE SWITCHING DIRECTION. SWI00027
C SWI00028
C NSE3=3*NSE SWI00029
C NOD3=3*NOD SWI00030
C NK22=NOD3-NSE3 SWI00031
C IF(IDR.EQ.1) THEN SWI00032
C NSAV=NSE3 SWI00033
C NSE3=NK22 SWI00034
C NK22=NSAV SWI00035
C ENDIF SWI00036
C SWI00037
C NSTP=NK22 SWI00038
C IF(NSTP.LT.NSE3) NSTP=NSE3 SWI00039
C DO 1 I=1,NSTP SWI00040
C IK22=NSE3+I SWI00041
C IS22=NK22+I SWI00042
C DO 2 J=1,NSTP SWI00043
C JK22=NSE3+J SWI00044
C JS22=NK22+J SWI00045
C SWI00046
C--- SWITCH [K21] ..... SWI00047
C IF(I.LE.NK22.AND.J.LE.NSE3) STIF1(I,J)=WRK(IK22,J) SWI00048
C--- SWITCH [K21] ..... SWI00049
C IF(I.LE.NSE3.AND.J.LE.NK22) STIF1(IS22,J)=WRK(I,JK22) SWI00050
C--- SWITCH [K22] ..... SWI00051
C IF(I.LE.NK22.AND.J.LE.NK22) STIF1(I,J)=WRK(IK22,JK22) SWI00052
C--- SWITCH [K11] ..... SWI00053
C IF(I.LE.NSE3.AND.J.LE.NSE3) STIF1(IS22,JS22)=WRK(I,J) SWI00054
C SWI00055

```

```

2 CONTINUE
1 CONTINUE

RETURN
END

```

```

SWI00056
SWI00057
SWI00058
SWI00059
SWI00060

```

```

SUBROUTINE GMASS(FNC,AJAC,XX,YY,ZZ,NODE,NODB3,NGS,SHAPE,
BIGN,BIGNT,FN,XMAS)

```

```

*****
THIS SUBROUTINE COMPUTES GLOBAL MASS MATRIX BY THE
CONSISTENT MASS APPROACH
*****

```

```

IMPLICIT REAL*8(A-H,O-Z)
DIMENSION SHAPE(NODE),FNC(NODE,3,10),BIGN(3,NODB3),
BIGNT(NODB3,3),FN(NODB3,NODB3),WGT1(10),
WGT2(10),WGT3(10),XMAS(NODB3,NODB3)

```

```

*****
DEFINITION OF VARIABLES
VARIABLE I/O DEFINITION
-----
FNC-----I-----MATRIX STORAGE CONTAINING THE VALUES
NONLINEAR INTERPOLATION FUNCTIONS EVALUATED
AT EACH GAUSS POINT. ITS DIMENSION IS
(NODE,3,10); WHERE
3-----INDICATES EVALUATION IN X,Y,Z DIRECTIONS
10----INDICATES NUMBER OF GAUSS POINTS
SHAPE-----ARRAY OF SHAPE FUNCTIONS AT NODE
BIGN-----MATRIX PRESENTATION OF SHAPE FUNCTIONS IN
ACCORDANCE WITH 3-D ELASTICITY.
i.e [BIGN]= N1 0 0 N2 0 0 N3 0 0 --- N(NODE) 0 0
0 N1 0 0 N2 0 0 N3 0 ----0 N(NODE) 0
0 0 N1 0 0 N2 0 0 N3 ----0 0 N(NODE)
BIGNT-----TRANSPOSE OF [BIGN]
FN-----THE FUNCTION [BIGN]T*[BIGN]
WGT1,WGT2,WGT3-----VECTOR CONTAINING THE WEIGHTS OF THE GAUSS
LEGENDRE QUADRATURE FOR INTEGRATION IN THE
X,Y,Z DIRECTIONS RESPECTIVELY
XMAS-----GLOBAL MASS MATRIX

```

```

RHO=120.

```

```

OBTAIN WEIGHTS FOR THE GAUSS-LEGENDRE QUADRATURE

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```

CALL NUMINT(XX,NODE,NODE,NGS,WGT1,1)
CALL NUMINT(YY,NODE,NODE,NGS,WGT2,1)
CALL NUMINT(ZZ,NODE,NODE,NGS,WGT3,1)

```

```

C      INITIALIZE MASS MATRIX
DO 1 I=1,NODB3
DO 1 J=1,NODB3
1 XMAS(I,J)=0.
C
DO 2 IGSX=1,NGS
DO 2 IGSY=1,NGS
DO 2 IGSZ=1,NGS
C
C      OBTAIN SHAPE FUNCTIONS FOR EACH NODE (EVALUATED AT EACH GAUSS POINT
C      N(XXi,YYi,ZZi)=f(XXi).g(YYi).h(ZZi)
C
DO 3 IPE=1,NODE
3 SHAPE(IPE)=FNC(IPE,1,IGSX)*FNC(IPE,2,IGSY)*FNC(IPE,3,IGSZ)
C
C      PLACE SHAPE FUNCTIONS IN MATRIX FORM FOR 3-D ELASTICITY
C
DO 4 I=1,3
DO 4 J=1,NODB3
4 BIGN(I,J)=0.
C
DO 5 IPE=1,NODE
J=3*IPE-2
JJ=J+1
JJJ=J+2
BIGN(1,J)=SHAPE(IPE)
BIGN(2,JJ)=SHAPE(IPE)
BIGN(3,JJJ)=SHAPE(IPE)
5 CONTINUE
C
C      T
C      COMPUTE [N] [N]
C
CALL TRANSP(BIGN,BIGNT,3,NODB3)
CALL MXMULT(BIGNT,BIGN,FN,NODB3,3,3,NODB3)
C
C      MULTIPLY THE FUNCTION FN BY THE WEIGHTS FOR THE GAUSS-LEGENDRE
C      QUADRATURE
C
WGHTS=WGT1(IGSX)*WGT2(IGSY)*WGT3(IGSZ)
DO 6 I=1,NODB3
DO 6 J=1,NODB3
6 XMAS(I,J)=XMAS(I,J)+WGHTS*FN(I,J)
C
2 CONTINUE
C
C      COMPUTE THE CONSISTENT GLOBAL MASS MATRIX
C
DO 7 I=1,NODB3
DO 7 J=1,NODB3
7 XMAS(I,J)=RHO*AJAC*XMAS(I,J)/32.2
C
PRINT*,'AJAC=',AJAC,'WGHTS=',WGHTS
PRINT*,'WGT1=',(WGT1(I),I=1,3)
PRINT*,'WGT2=',(WGT2(I),I=1,3)
PRINT*,'WGT3=',(WGT3(I),I=1,3)
C

```



```

WRITE(7,2020)
RETURN
C
  4 D(I)=A(I,I)/B(I,I)
 10 EIGV(I)=D(I)
    DO 30 I=1,NDOF
      DO 20 J=1,NDOF
        20 X(I,J)=0.
        30 X(I,I)=1.
          IF(NDOF.EQ.1)RETURN
C
C      INITIALIZE SWEEP COUNTER AND BEGIN ITERATION
C
      NSWEEP=0.
      NR=NDOF-1
    40 NSWEEP=NSWEEP+1
      IF(IFPR.EQ.1)WRITE(7,2000)NSWEEP
C
C CHECK IF PRESENT OFF-DIAGONAL ELEMENT IS LARGE ENOUGH TO REQUIRE ZEROING
C
      EPS=(.01**NSWEEP)**2
      DO 210 J=1,NR
        JJ=J+1
        DO 210 K=JJ,NDOF
          EPTOLA=(A(J,K)*A(J,K))/(A(J,J)*A(K,K))
          EPTOLB=(B(J,K)*B(J,K))/(B(J,J)*B(K,K))
          IF(EPTOLA.LT.EPS.AND.EPTOLB.LT.EPS)GO TO 210
C
C IF ZEROING IS REQUIRED, CALCULATE THE ROTATION MATRIX ELEMENT CA AND CG
C
      ARK=A(K,K)*B(J,K)-B(K,K)*A(J,K)
      AJJ=A(J,J)*B(J,K)-B(J,J)*A(J,K)
      AB=A(J,J)*B(K,K)-A(K,K)*B(J,J)
      CHECK=(AB*AB+4.*ARK*AJJ)/4.
      IF(CHECK)50,60,60
    50 WRITE(7,2020)
      RETURN
C
    60 SQCH=DSQRT(CHECK)
      D1=AB/2.+SQCH
      D2=AB/2.-SQCH
      DEN=D1
      IF(DABS(D2).GT.DABS(D1))DEN=D2
      IF(DEN)80,70,80
    70 CA=0.
      CG=-A(J,K)/A(K,K)
      GO TO 90
    80 CA=ARK/DEN
      CG=-AJJ/DEN
C
C PERFORM THE GENERALIZED ROTATION TO ZERO THE PRESENT OFF-DIAGONAL ELEMENT
C
    90 IF(NDOF-2)100,190,100
    100 J1=J+1
      JM1=J-1
      K1=K+1
      KM1=K-1
      IF(JM1-1)130,110,110
    110 DO 120 I=1,JM1
      AJ=A(I,J)

```

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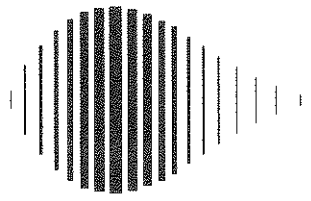
      BJ=B(I,J)
      AK=A(I,K)
      BK=B(I,K)
      A(I,J)=ZJ+CG*AK
      B(I,J)=BJ+CG*BK
      A(I,K)=AK+CA*AJ
120  B(I,K)=BK+CA*BJ
130  IF(KP1-NDOF)140,140,160
140  DO 150 I=KP1,NDOF
      AJ=A(J,I)
      BJ=B(J,I)
      AK=A(K,I)
      BK=B(K,I)
      A(J,I)=AJ+CG*AK
      B(J,I)=BJ+CG*BK
      A(K,I)=AK+CA*AJ
150  B(K,I)=BK+CA*BJ
160  IF(JP1-KM1)170,170,190
170  DO 180 I=JP1,KM1
      AJ=A(J,I)
      BJ=B(J,I)
      AK=A(I,K)
      BK=B(I,K)
      A(J,I)=AJ+CG*AK
      B(J,I)=BJ+CG*BK
      A(I,K)=AK+CA*AJ
180  B(I,K)=BK+CA*BJ
190  AK=A(K,K)
      BK=B(K,K)
      A(K,K)=AK+2.*CA*A(J,K)+CA*CA*A(J,J)
      B(K,K)=BK+2.*CA*B(J,K)+CA*CA*B(J,J)
      A(J,J)=A(J,J)+2.*CG*A(J,K)+CG*CG*AK
      B(J,J)=B(J,J)+2.*CG*B(J,K)+CG*CG*BK
      A(J,K)=0.
      B(J,K)=0.
C
C  UPDATE EIGENVECTOR MATRIX AFTER EACH ROTATION
C
      DO 200 I=1,NDOF
      XJ=X(I,J)
      XK=X(I,K)
      X(I,J)=XJ+CG*XK
200  X(I,K)=XK+CA*XJ
210  CONTINUE
C
C  UPDATE THE EIGENVALUES AFTER EACH SWEEP
C
      DO 220 I=1,NDOF
      IF(A(I,I).GT.0..AND.B(I,I).GT.0.)GO TO 220
      WRITE(7,2020)
      RETURN
C
220  EIGV(I)=A(I,I)/B(I,I)
      IF(IPPR.EQ.0)GO TO 230
      WRITE(7,2030)
      WRITE(7,2010)(EIGV(I),I=1,NDOF)
C
C  CHECK FOR CONVERGENCE
C
230  DO 240 I=1,NDOF

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TOL=RTOL*D(I)
DIF=DABS(EIGV(I)-D(I))
IF(DIF.GT.TOL)GO TO 280
240 CONTINUE
C
C CHECK ALL OFF-DIAGONAL ELEMENTS TO SEE IF ANOTHER SWEEP IS REQUIRED
C
EPS=RTOL**2
DO 250 J=1, NR
JJ=J+1
DO 250 K=JJ, NDOF
EPSA=(A(J,K)*A(J,K))/(A(J,J)*A(K,K))
EPSB=(B(J,K)*B(J,K))/(B(J,J)*B(K,K))
IF((EPSA.LT.EPS).AND.(EPSB.LT.EPS))GO TO 250
GO TO 280
250 CONTINUE
C
C FILL OUT BOTTOM TRIANGLE OF RESULTANT MATRICES AND SCALE EIGENVECTORS
C
255 DO 260 I=1, NDOF
DO 260 J=1, NDOF
A(J,I)=A(I,J)
260 B(J,I)=B(I,J)
DO 270 J=1, NDOF
BB=DSQRT(B(J,J))
DO 270 K=1, NDOF
270 X(K,J)=X(K,J)/BB
RETURN
C
C UPDATE [D] MATRIX AND START NEW SWEEP, IF ALLOWED
C
280 DO 290 I=1, NDOF
290 D(I)=EIGV(I)
IF(NSWEEP.LT.NSMAX)GO TO 40
GO TO 255
C
C
C
2000 FORMAT(27H0SWEEP NUMBER IN * EIGN * =,I4)
2010 FORMAT(1H0,6E20.12)
2020 FORMAT(25H0*** ERROR SOLUTION STOP /
30H MATRICES NOT POSITIVE DEFINITE)
2030 FORMAT(36H0CURRENT EIGENVALUES IN * EIGN * ARE,/)
END

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